B Physics: New Physics and The Next Generation

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<u>Complex phases in the weak</u> <u>interaction</u>: V_{td} and V_{ts} and associated CPV asymmetries

Excitement in Flavor Physics:

- -Connections to the *charged Higgs*
- -Rare B Decays + NP

Flavor Physics, The Next Generation: Belle II and the LHCb upgrade

Apologies: In the limited time, I cannot cover all the recent results from BaBar, Belle, LHCb, CMS, ATLAS, Tevatron ... I have borrowed slides from many excellent physicists and will aim for the "big picture" but skip most details.

Amplitudes and Phases in the Weak Interaction

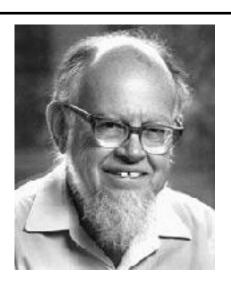






$$V_{\rm CKM} \equiv V_L^u V_L^{d\dagger} =$$

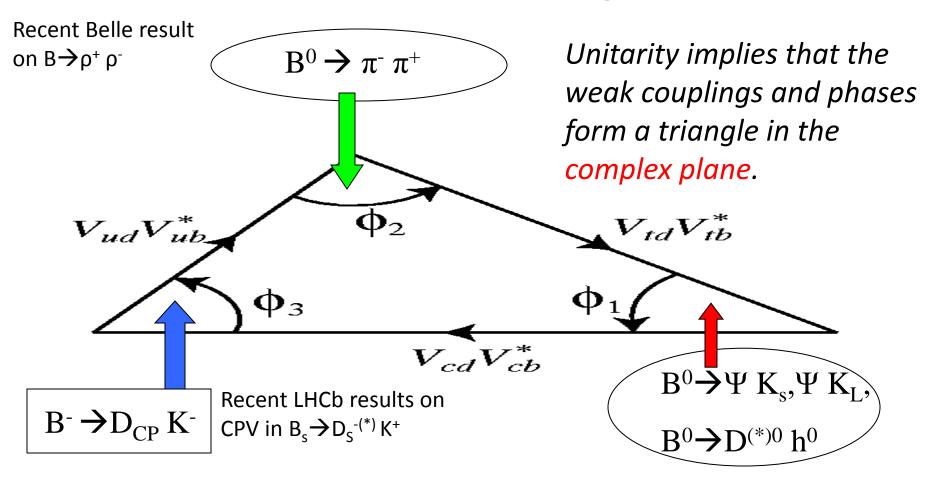
$$egin{pmatrix} V_{ud} & V_{us} & V_{ub} \ V_{cd} & V_{cs} & \overline{V_{cb}} \ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



$$V = \begin{pmatrix} \begin{cases} 1 - \frac{1}{2}\lambda^2 \\ -\lambda \\ A\lambda^3 (1 - \rho - i\eta) \end{cases}$$

L. Wolfenstein (1923-2015)

Three Angles: $(\varphi_1, \varphi_2, \varphi_3)$ or (β, α, γ)

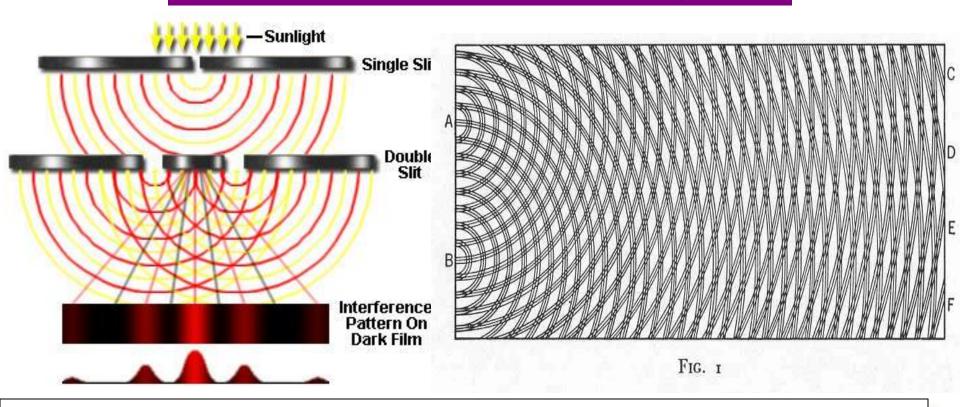


Big Questions: Are determinations of angles consistent with determinations of the sides of the triangle? Are angle determinations from loop and tree decays consistent?

Time-dependent CP violation is

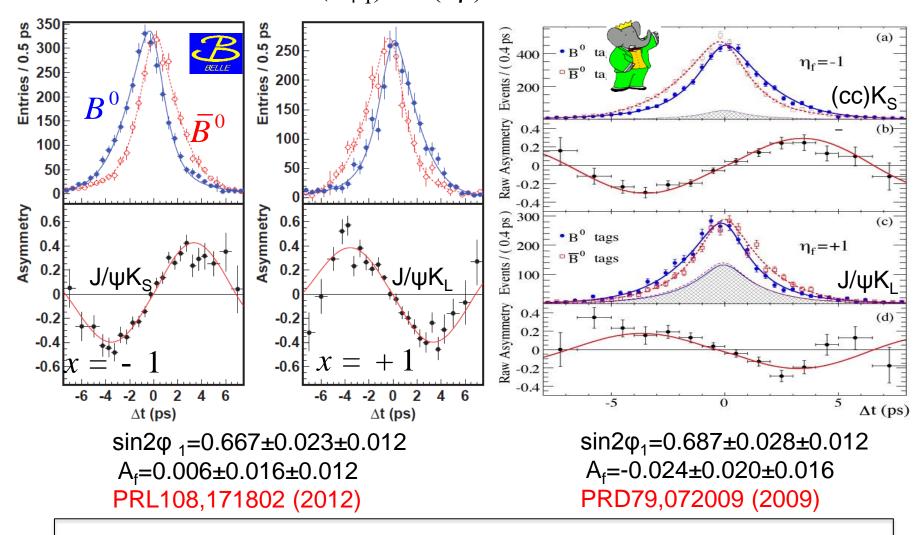
"A Double-Slit experiment" with particles and antiparticles

QM interference between two diagrams



Measures the <u>phase</u> of V_{td} or equivalently the <u>phase</u> of B_d —anti B_d mixing.

Measurement of $\sin(2\varphi_1)/\sin(2\beta)$ in B \rightarrow Charmonium K⁰ modes



Overpowering evidence for CP violation (matter-antimatter asymmetries). >>>> The phase of V_{td} is in good agreement with Standard Model expectations. This is the phase of B_d mixing.

Breaking news: 2016 APS Panofsky Prize for Experimental Particle Physics Awarded to

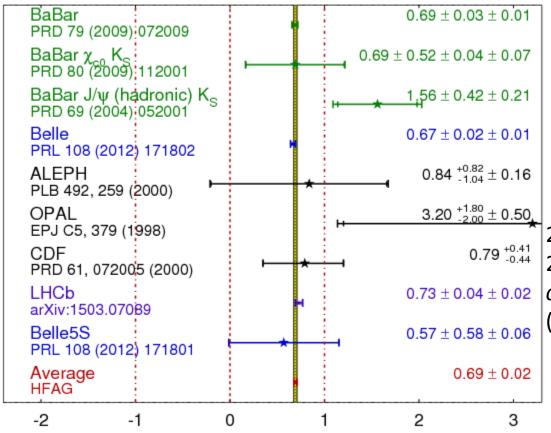
Steve Olsen, Dave Hitlin, Jonathan Dorfan, and Fumihiko Takasaki



A scene in Stockholm, Sweden 2008

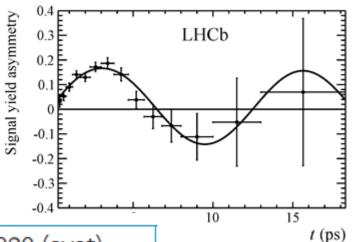
Front row 2008 Physics Nobelists: T. Maskawa, M. Kobayashi

$\sin(2\beta) \equiv \sin(2\phi_1) \frac{\text{HFAG}}{\text{Moriond 2015}}$



B factories: High precision CPV measurement and a calibration for NP.

2013: LHCb joins the game [±0.07(stat)] 2015: Latest LHCb measurement is comparable in precision to B factories (now uses same-side B tagging)

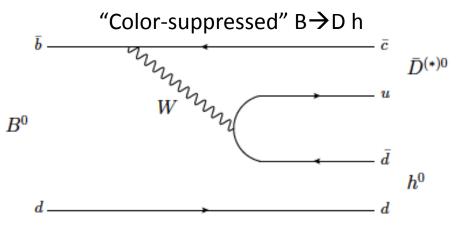


$$S(B^0 \to J/\psi K_s^0) = 0.731 \pm 0.035 \text{ (stat)} \pm 0.020 \text{ (syst)},$$

 $C(B^0 \to J/\psi K_s^0) = -0.038 \pm 0.032 \text{ (stat)} \pm 0.005 \text{ (syst)}.$

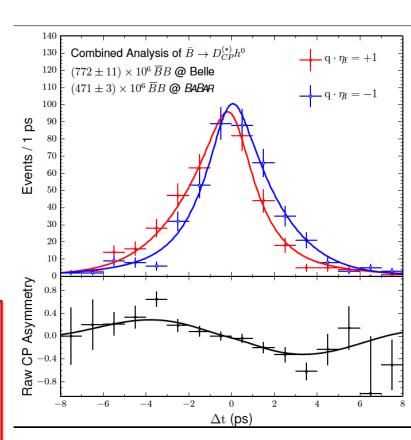
B factories: Check CP violation in $b \rightarrow c$ [ubar d] processes

2015: First joint BaBar-Belle data analysis м. Rohrken et al



where D^0 is a CP eigenstate and $h^0=\pi^0$, η , ω

Combining Belle and BaBar datasets, ~1260 signal events, obtain a 5.4 σ CP violation signal \rightarrow First observation sin(2 β_{eff})=0.66±0.10(stat)±0.06(sys)



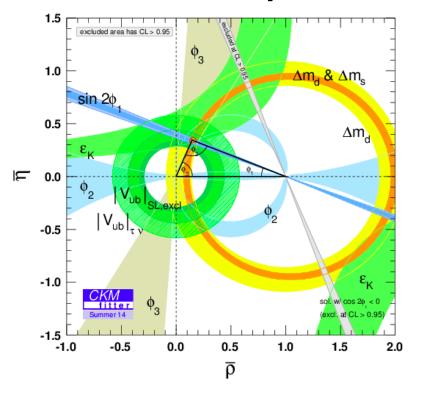
Phase of V_{td} again

Conclusion: CP violation in $b \rightarrow c$ ubar d modes is the same as in $b \rightarrow c$ cbar s modes (e.g. $B \rightarrow J/\psi K_s$)

Results from Global Fits to Data (CKMFitter Group)

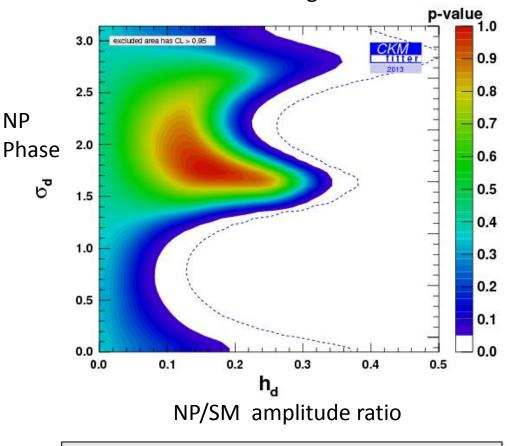
Great progress on ϕ_3 or γ (first from B factories and now in the last two years from LHCb). *These measure* the phase of V_{ub} [CKM2014, K.

Trabelsi's review: ±70]



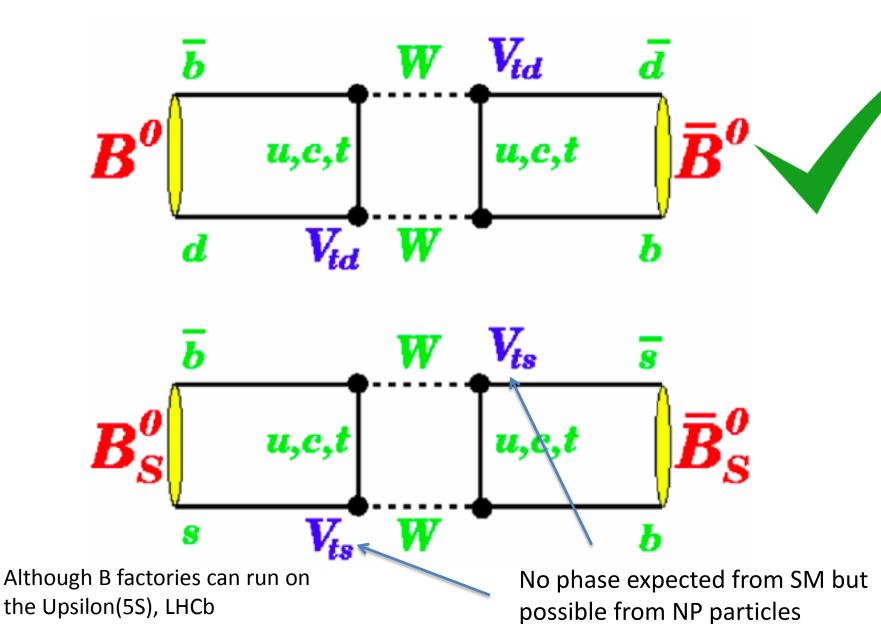
Looks good (except for an issue with $|V_{ijh}|$)

Similar results from UTFIT as well from G. Eigen et al.



But a 10-20% NP amplitude in B_d mixing is perfectly compatible with all current data.

Boxes

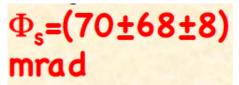


dominates here

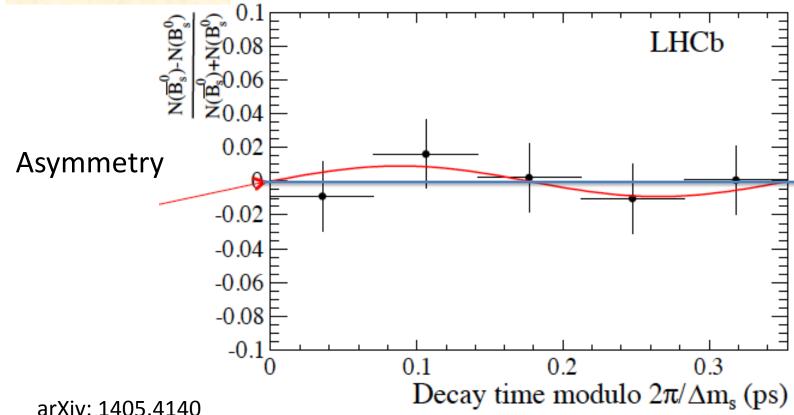
10

B_s \rightarrow J/ψφ, a pseudoscalar to vector-vector mode, is usually used. However, B_s \rightarrow J/ψ f₀(980) is a pure CP eigenstate since the f₀ (980) is a scalar.

Stone & Zhang pointed out that this mode provides more statistics and a more straightforward analysis. Phys. Rev. D79 (2009) 074024.

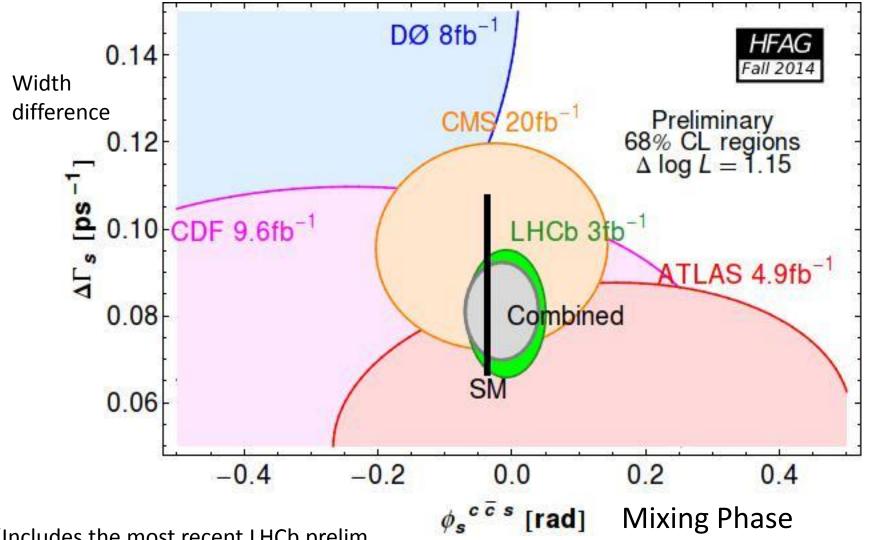


Red curve: expectation for Φ_s = 70 mrad



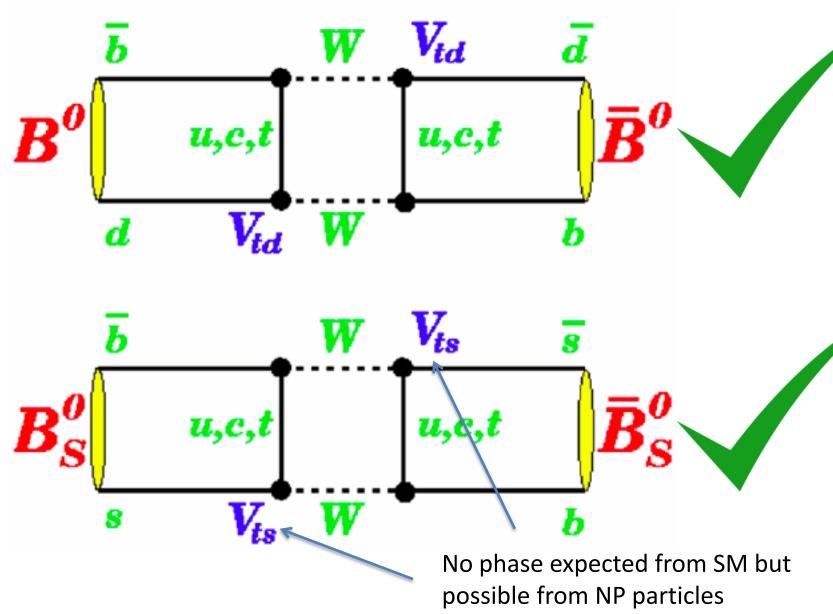
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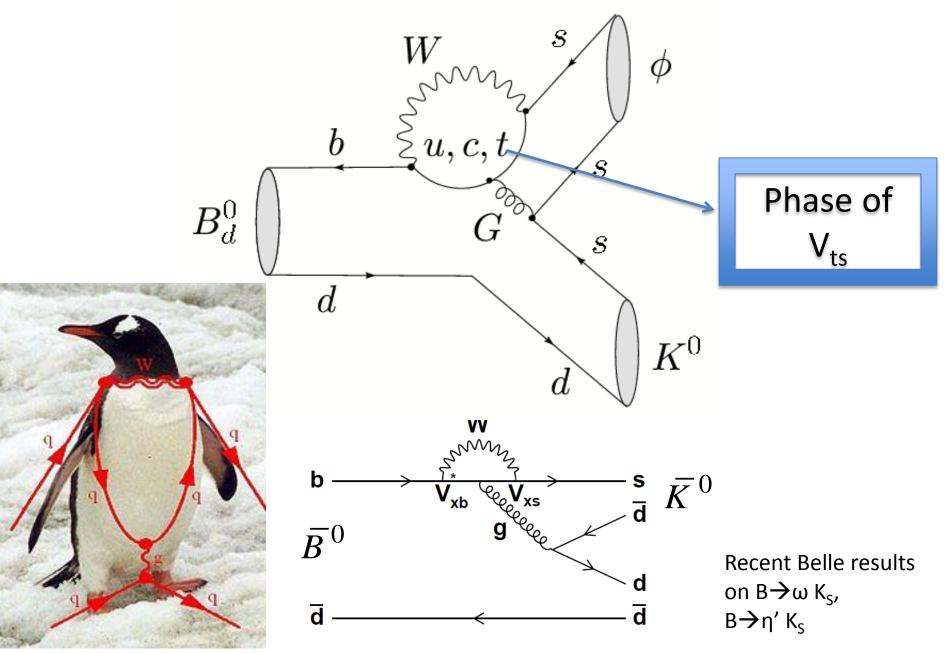
Results on the <u>phase of B_s-anti B_s mixing</u> (i.e. phase of V_{ts}) [use B_s \rightarrow J/ $\psi \phi$; J/ $\psi \pi \pi$ modes]



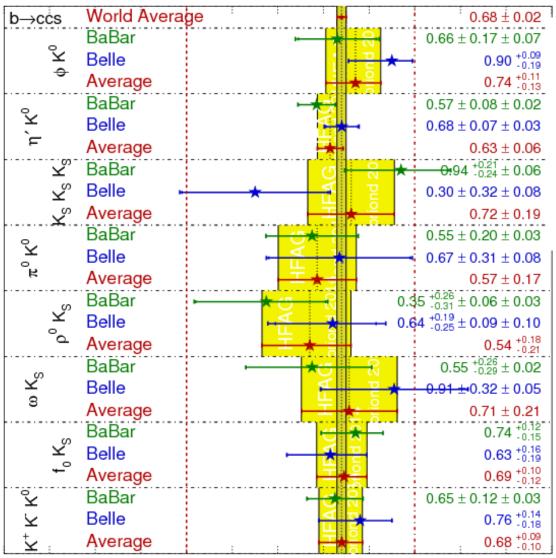
(Includes the most recent LHCb prelim result, gives WA of -36 ±13 mrad)

Boxes

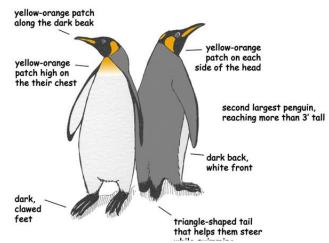




$sin(2\beta^{eff}) \equiv sin(2\phi_1^{eff}) \underset{\text{PRELIMINARY}}{\text{HFAG}}$



New Physics Phases in Penguin b→s decays



No evidence for NP at current level of sensitivity

LHCb is absent from this game (lower K_S eff and flavor tagging eff) but contributes in B_S modes.

"Missing Energy" Decays

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2 Accelerators Find Particles That May Break Known Laws of Physics

The LHC and the Belle experiment have found particle decay patterns that violate the Standard Model of particle physics, confirming earlier observations at the BaBar facility

By Clara Moskowitz | September 9, 2015 | Véalo en español

News & Features Topics



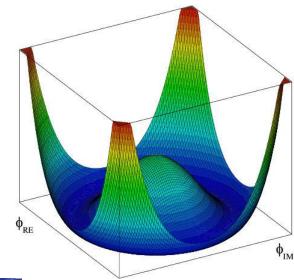
Democracy suffers a blow-in particle physics

Three independent B-meson experiments suggest that the charged leptons may not be so equal after all.

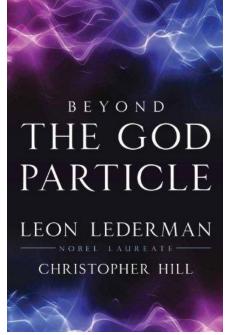
Steven K. Blau 17 September 2015

Subscribe

The BEH boson is now firmly established by experimental results from ATLAS and CMS. Now planning for Higgs flavor factory facilities (e.g ILC, FCC, CEPC, H-LHC).



Does the GP (Brout-Englert-Higgs particle) have a "brother" i.e. the charged Higgs?





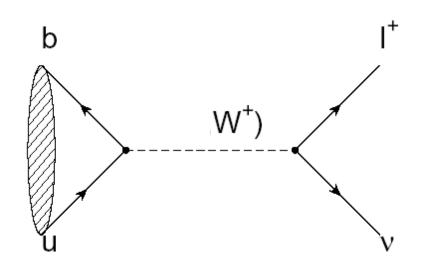
Y. Nambu, 1921-2015

Measurements at Belle II and direct searches at hadron colliders take complementary approaches to this important question.

$B \rightarrow \tau \nu$

(Decay with Large Missing Energy)

Sensitivity to new physics from a charged Higgs



$$\mathcal{B}(B^{+} \to \tau^{+} \nu_{\tau}) = \frac{G_{F}^{2} m_{B}}{8\pi} m_{\tau}^{2} \left(1 - \frac{m_{\tau}^{2}}{m_{B}^{2}} \right)^{2} f_{B}^{2} |V_{ub}|^{2} \tau_{B}$$

$$\mathcal{B}_{(B \to \tau \nu)} = \mathcal{B}_{SM} \times \left(1 - \tan^{2} \beta \frac{m_{B^{\pm}}^{2}}{m_{H^{\pm}}^{2}} \right)$$

W.S.Hou,. PRD 48, 2342 (1993)

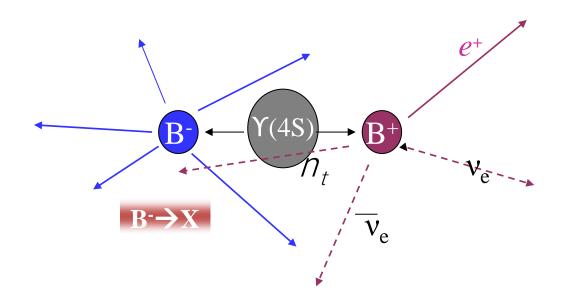
The B meson decay constant, determined by the B wavefunction at the origin

 $(|V_{ub}|)$ taken from indep. measurements.)

Consumer's guide to charged Higgs

- <u>Higgs doublet of type I</u> (ϕ_1 couples to upper (u-type) and lower (d-type) generations. No fermions couple to ϕ_2)
- <u>Higgs doublet of type II</u> (φ_u couples to u type quarks, φ_d couples to d-type quarks, u and d couplings are different; $tan(\beta) = v_u/v_d$) [<u>favored NP scenario</u> e.g. MSSM, generic SUSY]
- Higgs doublet of type III (not type I or type II; anything goes. "FCNC hell"→many FCNC signatures)

Why measuring $B^+ \rightarrow \tau^+ v$ is non-trivial



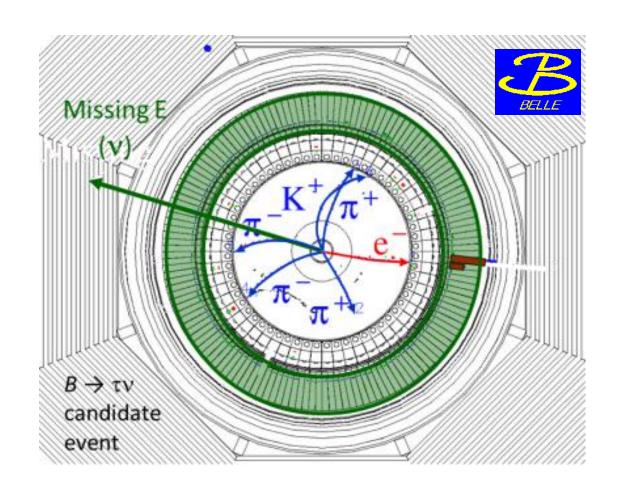
Most of the sensitivity is from tau modes with 1-prongs.

The experimental signature is rather difficult: B decays to a single charged track + nothing

(This may be hard at a hadron collider)

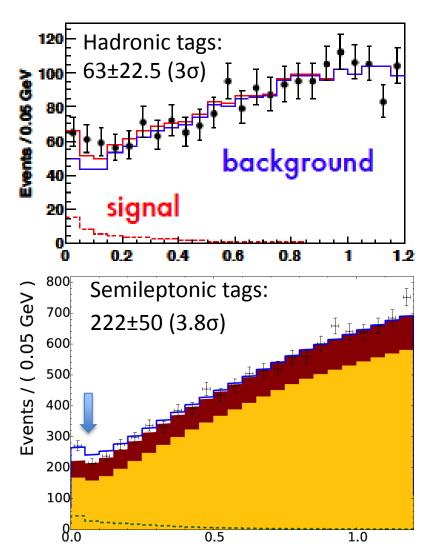
Example of a Missing Energy Decay ($B \rightarrow \tau \nu$) in <u>Data</u>

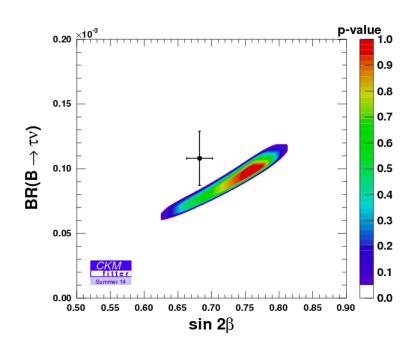
$$B^+ o D^0\pi^+ \ (o K\pi^-\pi^+\pi^-) \ B^- o au(o e
uar
u)
u$$



The clean e+e- environment makes this possible

Example: Belle B \rightarrow tv results with full *reprocessed* data sample and either hadronic or semileptonic tags (arXiv: 1409.5269 \rightarrow PRD)





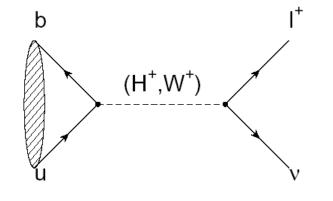
With the full B factory statistics only "evidence". No single observation from either Belle or BaBar.

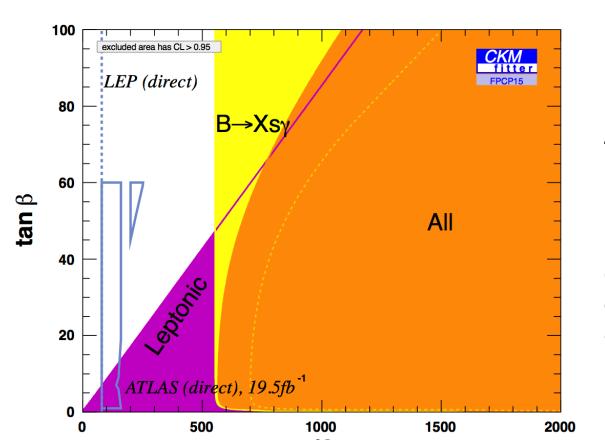
The horizontal axis is the "Extra Calorimeter Energy"

Complementarity of e+ e- factories and LHC

(Slide adapted from A. Bevan)

The current combined B→τυ limit places a stronger constraint than direct searches from LHC exps. for the next few years.



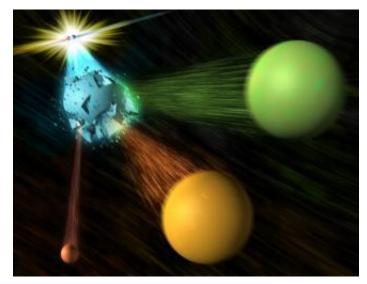


$$r_H = \left(1 - \frac{m_B^2}{m_H^2} \tan^2 \beta\right)^2$$

Currently inclusive b to sy rules out m_{H+} below ~480 GeV/c² range at 95% CL (independent of tan β), M. Misiak et al.

http://arxiv.org/abs/1503.01789

$$B\{rac{b}{ar{q}} - C^{ar{ au}}\}D^{(*)}$$



$$\mathcal{R}(D^{(*)})_{\mathrm{2HDM}} = \mathcal{R}(D^{(*)})_{\mathrm{SM}} + A_{D^{(*)}} \frac{\tan^2\beta}{m_{H^+}^2} + B_{D^{(*)}} \frac{\tan^4\beta}{m_{H^+}^4}$$

$$D\tau \mathbf{v} \qquad D^*\tau \mathbf{v}$$

$$A_{D^{(*)}} (\mathrm{GeV^2}) \qquad -3.25 \pm 0.32 \quad -0.230 \pm 0.029$$

$$B_{D^{(*)}} (\mathrm{GeV^4}) \qquad 16.9 \pm 2.0 \qquad 0.643 \pm 0.085$$

$$R(D^{(*)}) = \frac{\mathcal{B}(\overline{B} \to D^{(*)}\tau^{-}\overline{\nu}_{\tau})}{\mathcal{B}(\overline{B} \to D^{(*)}\ell^{-}\overline{\nu}_{\ell})} \stackrel{\text{Signal}}{\longrightarrow} \text{Normalization } (l = e \text{ or } \mu)$$

Example from a BaBar paper

Signals in B \rightarrow D^(*) τv (489±63, 888±63)

Missing mass variable:

$$m_{miss}^2 = p_{miss}^2 = (p[e^+e^-] - p_{tag}^2 - p_{D(*)}^2 - p_I)^2$$

P_I* = momentum of lepton in B rest frame

But wait !!! Now possible at LHCb.

Production of B meson pairs at threshold is critical to the separation of backgrounds from the missing energy/momentum signal.

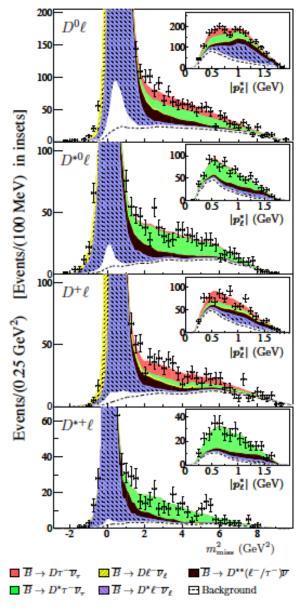
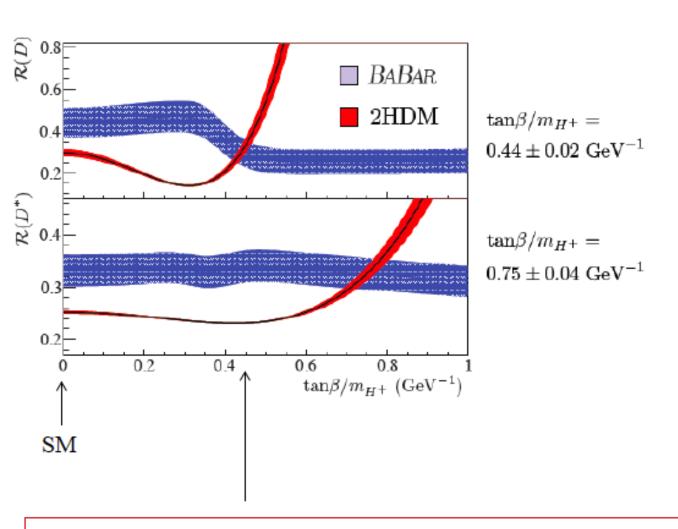
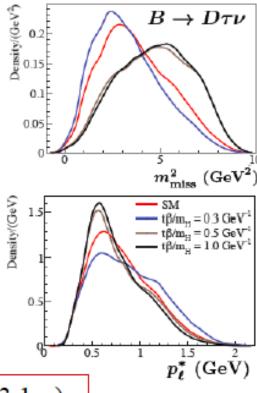


FIG. 1. (Color online) Comparison of the data and the fit projections for the four $D^{(*)}\ell$ samples. The insets show the $|p_{\ell}^*|$ projections for $m_{\text{miss}}^2 > 1 \text{ GeV}^2$, which excludes most of the normalization modes. In the background component, the region above the dashed line corresponds to charge cross-feed, and the region below corresponds to continuum and $B\overline{B}$.

Limits on type-II 2HDM



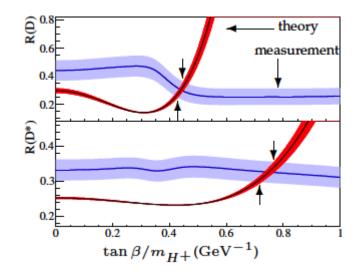
2HDM modifies fitvariable distribution and hence the efficiency



Best point is $\tan \beta / m_{H^+} = 0.45 \text{ GeV}^{-1}$, excluded at 99.8% CL (3.1 σ). All other values (with $m_{H^+} > 15 \text{ GeV}$) are worse.

BaBar collaboration, Phys. Rev. Lett. 109, 101802 (2012)

"However, the combination of R(D) and R(D*) excludes the type II 2HDM charged Higgs boson with a 99.8% confidence level for any value of $\tan(\beta)/m_{H+}$ "



In other words, found NP but killed the 2HDM NP model.

T. Kuhr

http://xxx.lanl.gov/abs/1507.03233; to appear in Phys Rev D

Warning: color-coding different from BaBar

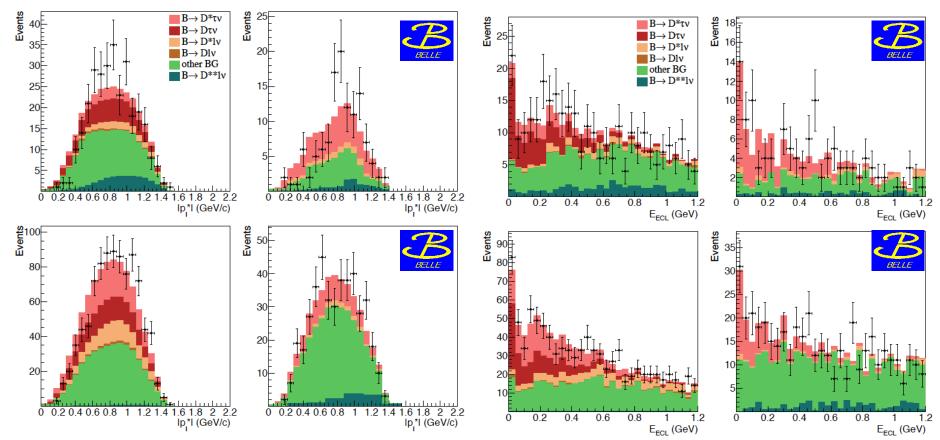


FIG. 6. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of $M_{\text{miss}}^2 > \text{FIG. 5}$.

Signal enhanced projections of <u>lepton momenta</u> in the high M²_{miss} region

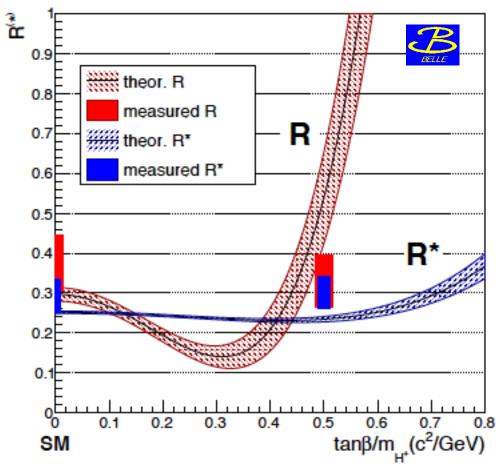
Signal enhanced projections of extra calorimeter energy in the high M²_{miss} region

New Belle result with hadronic tags

May 25 2015, Nagoya FPCP

http://xxx.lanl.gov/abs/1507.03233; to appear in Phys Rev D

Compatible with both BaBar and the 2HDM model (and SM!).



Need more data and more Belle analyses to resolve the issue.

FIG. 8. Theoretical predictions with 1σ error ranges for R (red) and R^* (blue) for different values of $\tan \beta/m_{H^+}$ in the 2HDM of type II. This analysis' fit results for $\tan \beta/m_{H^+} = 0.5 c^2/\text{GeV}$ and SM are shown with their 1σ ranges as red and blue bars with arbitrary width for better visibility.

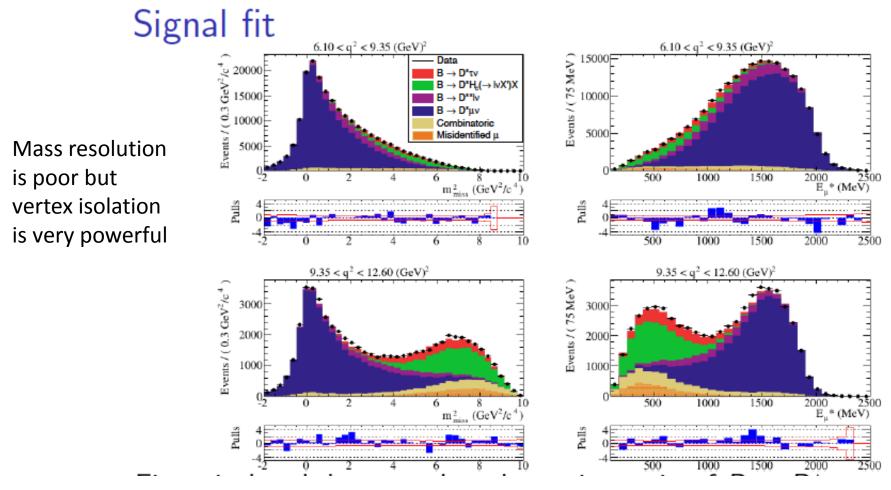
New LHCb result

Compatible with BaBar $D^*\tau U$ BF (B \rightarrow D τU in the pipeline)

May 25 2015, Nagoya FPCP

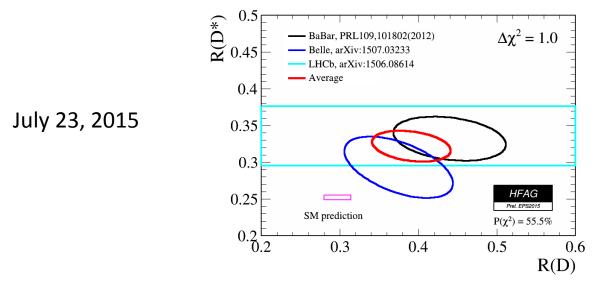
Published in **Phys. Rev. Lett. 115, 111803** (2015)

3. $B \rightarrow D^* \tau \nu$ (LHCB-PAPER-2015-025)



Oui, c'est possible!

Apres Nagoya: New World Averages for R(D) and R(D*)

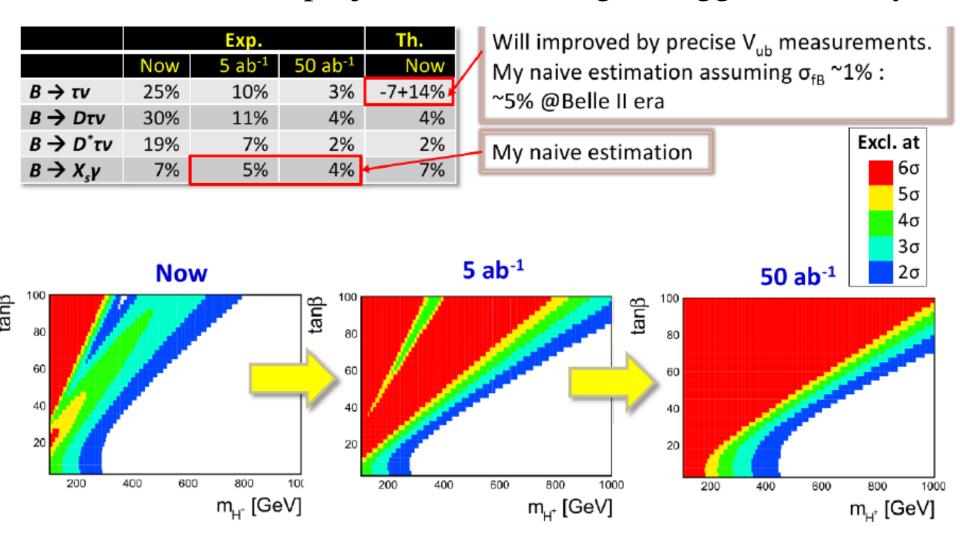


Now 3.9σ from SM

	R(D)	$R(D^*)$
BaBar	$0.440 \pm 0.058 \pm 0.042$	$0.332 \pm 0.024 \pm 0.018$
Belle	$0.375^{+0.064}_{-0.063} \pm 0.026$	$0.293^{+0.039}_{-0.037} \pm 0.015$
LHCb		$0.336 \pm 0.027 \pm 0.030$
Average	0.388 ± 0.047	0.321 ± 0.021
SM expectation	0.300 ± 0.010	0.252 ± 0.005
Belle II, 50/ab	± 0.010	± 0.005

It is *obvious* that we <u>need two orders of magnitude of data</u> to solve these issues related to the *charged Higgs*.

Initial Belle II projections for charged Higgs sensitivity

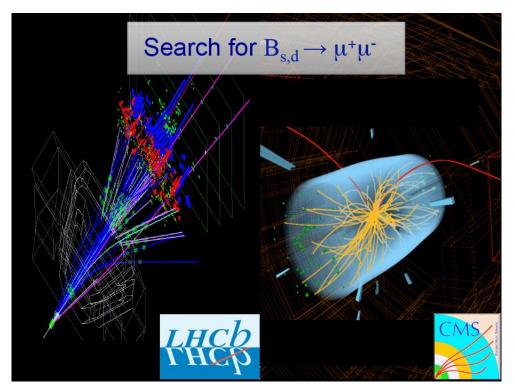


Yutaro Sato, R. Itoh et al

Charged Higgs: A strong case at Belle II

Rare B Decays

Two event displays

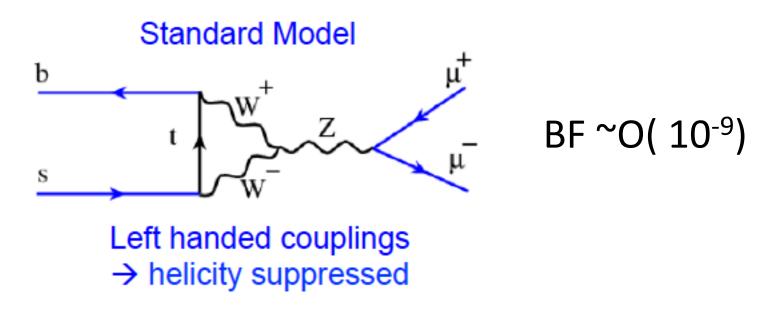


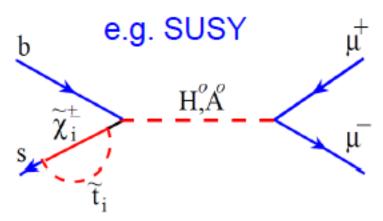
J. Albrecht



Goa, India

LHC found the rarest B decay; $B_s \rightarrow \mu + \mu$ -





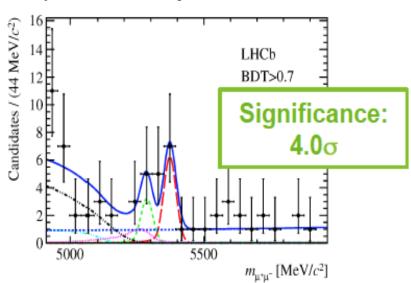
N. B. Here and in $b \rightarrow s l^+ l^-$ all the heavy particles of the SM enter as virtual particles in the Feynman diagrams

LHCb

- Update: full dataset: 3fb-1
 - Improved BDT

PRL 111 (2013) 101805

Expected sensitivity: 5.0 o

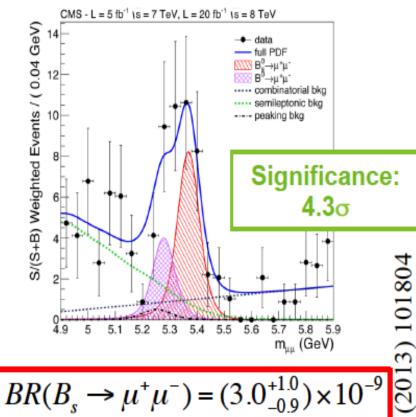


$$BR(B_s \to \mu^+ \mu^-) = (2.9^{+1.1}_{-1.0}) \times 10^{-9}$$

$$BR(B^0 \to \mu^+ \mu^-) = (3.7^{+2.4}_{-2.1}) \times 10^{-10}$$

 $BR(B^0 \to \mu^+ \mu^-) < 7 \times 10^{-10} @ 95\%CL$

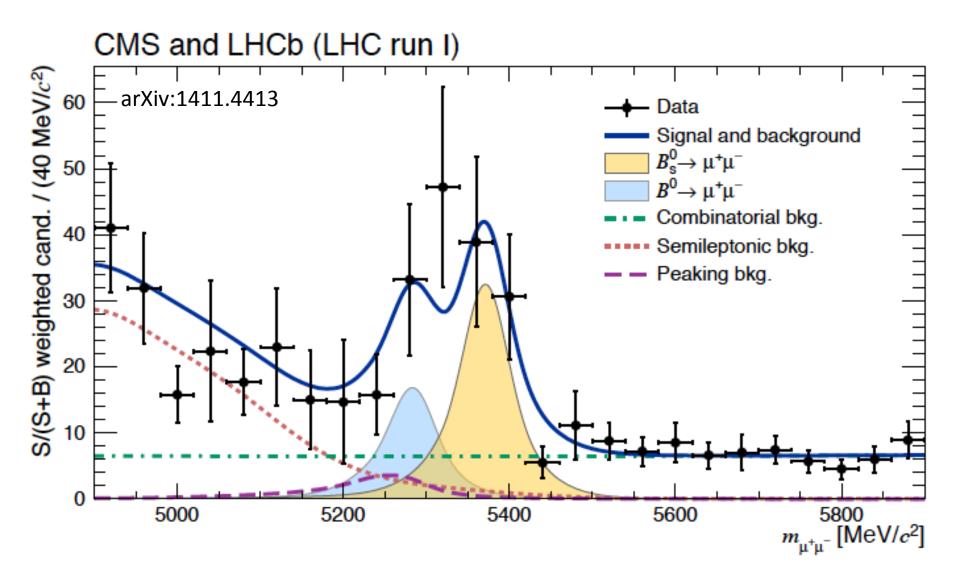
CMS



$$BR(B_s \to \mu^+ \mu^-) = (3.0^{+1.0}_{-0.9}) \times 10^{-9}$$

$$BR(B^0 \to \mu^+ \mu^-) = (3.5^{+2.1}_{-1.8}) \times 10^{-10}$$

 $BR(B^0 \to \mu^+ \mu^-) < 11 \times 10^{-10} @ 95\%CL$



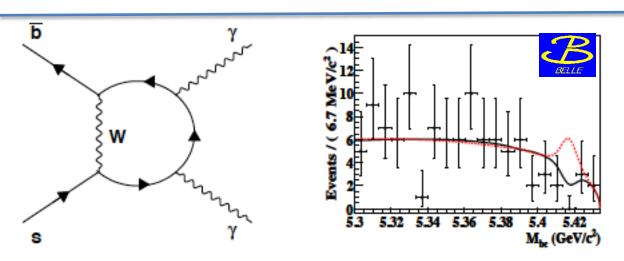
Published in Nature: June 4, 2015

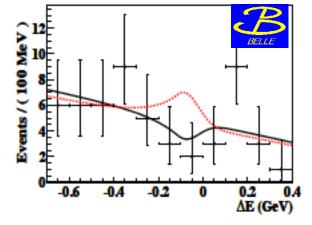
$$\mathcal{B}(B_s^0 \to \mu^+ \mu^-) = (2.9 \pm 0.7) \times 10^{-9}$$

 $\mathcal{B}(B^0 \to \mu^+ \mu^-) = (3.6 ^{+1.6}_{-1.4}) \times 10^{-10}$

SM: BR(B_s) = (3.65±0.23) 10^{-9} BR(B⁰) = (1.1±0.1) 10^{-10} PRL 112 101801 (2014)

Combining evidence from two LHC experiments (LHCb and CMS), $B_s \rightarrow \mu^+ \mu^-$ is observed with 6.2 σ significance. The corresponding B_d decay is not clearly seen yet.





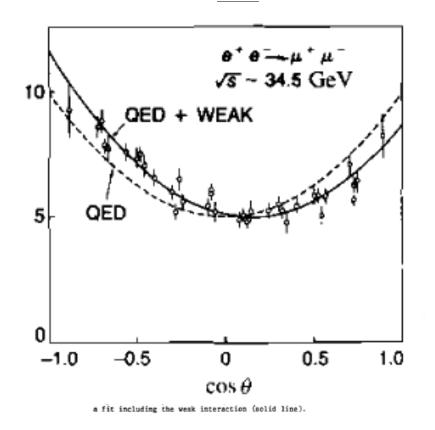
BF ($B_s \rightarrow \gamma \gamma$) <3.1 x 10⁻⁶

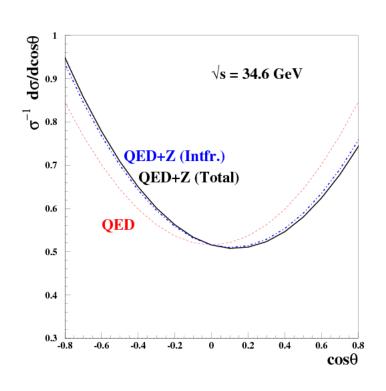
D. Dutta et al (Belle)
Phys. Rev. D 91, 011101(R)

Red Hot Flavor Physics



High Energy Physics History: finding NP in A_{FB} (using interference)

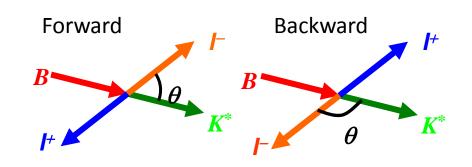




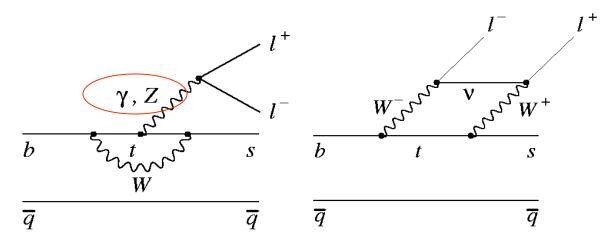
Conclusion: There is a Z boson at higher energy even though colliders of the time did not have enough \sqrt{s} to produce it

$$A_{FB}(B \rightarrow K^*l^+l^-)(q^2)$$

The SM forward-backward asymmetry in $b \rightarrow s l^+ l^-$ arises from the <u>interference</u> between γ and Z^0 contributions.

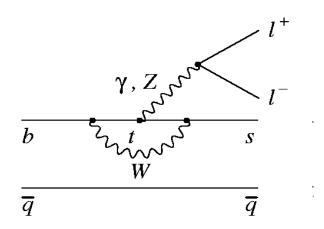


$$A_{FB}(B \to K^* \ell^+ \ell^-) = -C_{10} \xi(q^2) \left[Re(C_9) F_1 + \frac{1}{q^2} C_7 F_2 \right]$$
Ali, Mannel, Morozumi, PLB273, 505 (1991)



Note that all the heavy particles of the SM (W, Z, top) enter in this decay.

More on $A_{FB}(B \rightarrow K^*l^+l^-)(q^2)$

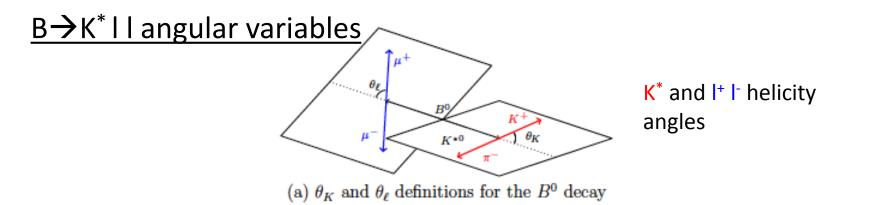


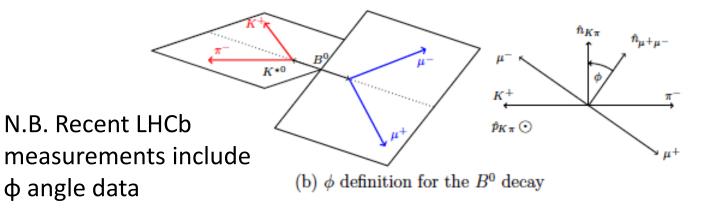
Can in effect vary \sqrt{s} for \sqrt{NP}

 A_{FB} depends on $q^2 = M^2(I^+I^-)$

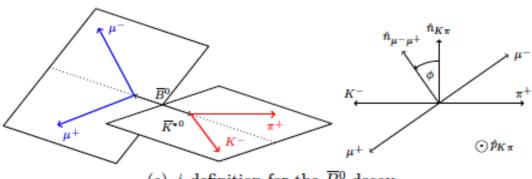
$$A_{FB}(B \to K^* \ell^+ \ell^-) = -C_{10} \xi(q^2) \left[Re(C_9) F_1 + \frac{1}{q^2} C_7 F_2 \right]$$
Ali, Mannel, Morezumi, PLB273, 505 (1991)

The "zero-crossing" of A_{FB} depends only on a ratio of form factors and is a *clean* observable.





Angle between <u>the</u> <u>normals</u> to the two decay planes.



(c) ϕ definition for the \overline{B}^0 decay

From the 2013 LHCb paper

$B \rightarrow K^*l^+l^-(q^2)$ bootcamp

Angular dependence



(-) means the term is only in G - G

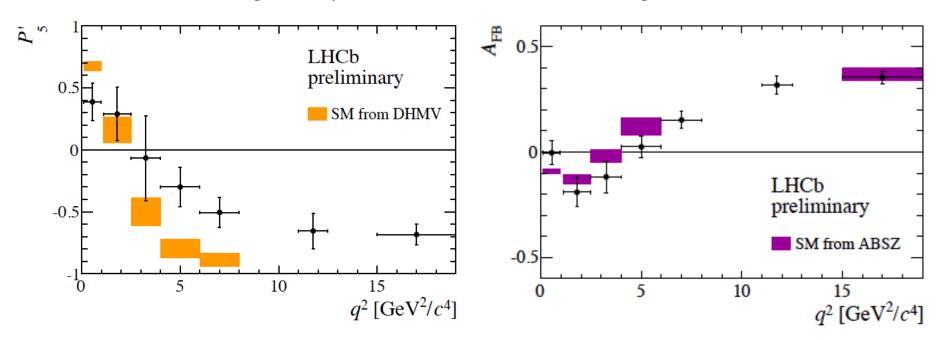
$$\frac{1}{d(\Gamma + \overline{\Gamma})/dq^2} \frac{d(\Gamma + \overline{\Gamma})}{d\overline{\Omega}} = F_L \text{ is the longitudinal polarization fraction.}$$

$$= \begin{bmatrix} \frac{3}{4}(1 - F_L)\sin^2\vartheta_K + F_L\cos^2\vartheta_K \\ + \frac{1}{4}(1 - F_L)\sin^2\vartheta_K\cos2\vartheta_L \\ -F_L\cos^2\vartheta_K\cos2\vartheta_L + S_3\sin^2\vartheta_K\sin^2\vartheta_L\cos2\varphi \\ + S_4\sin2\vartheta_K\sin2\vartheta_L\cos\varphi + \\ + S_7\sin2\vartheta_K\sin2\vartheta_K\sin2\vartheta_L\sin\varphi \end{bmatrix}$$

Introduce $P_{4,5} = S_{4,5} / sqrt[F_L(1-F_L)]$ to reduce dependence on form factors

New LHCb 3fb⁻¹ results on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

Angular Asymmetries based on 2398±57 signal events



"The P_5 ' measurements are only compatible with the SM prediction at a level of 3.7σ A mild tension can also be seen in the A_{FB} distribution, where the measurements are systematically <=1 σ below the SM prediction in the region $1.1 < q^2 < 6.0 \text{ GeV}^2$ "

Blank regions are the J/ ψ and ψ '

Theory from http://arxiv.org/abs/1407.8526

Experiment from LHCb-CONF-2015-002

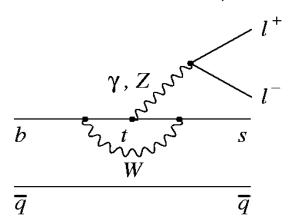
vetos

New LHCb results on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

<u>Is HEP History repeating itself</u>? [Make sure this is not a tricky SM form factor effect.]

Why does NP appear first in this mode

(and not others)?





Possible answer: All the heavy particles of the SM (t, W, Z) and maybe NP (except the Higgs) appear here. Sensitive to NP via interference (linear effects).

Theory issues on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

"To better understand the compatibility of the data with the SM a global analysis of the data, taking into account correlations between observables is necessary"-LHCb conference paper.

- → Check dependence on light-cone form factors (some checks already done by Lattice QCD group, or *fit form factors from data a la Mandal and Sinha and check for consistency*)
- → Check binning effects (do an unbinned analysis of NP)
- → Can tails of large B → K* [c-cbar] produce the anomalies found in the angular distributions?



B→K* I+ I- form factor ratios determined from data disagree with theory

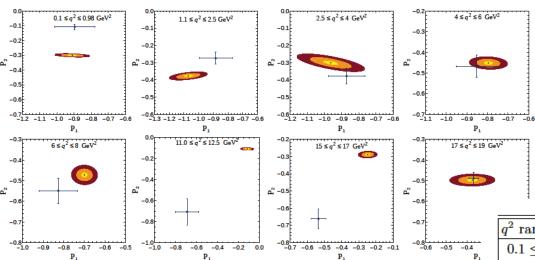


FIG. 1. (color online). The allowed region for P₁ versus P₂ plane. The innermost yellow (lightest), the middle of and outer most red (dark) contours represent 1σ , 3σ and 5σ regions, respectively. The theoretically estimate Ref. [7] for $q^2 \leq 8 \text{ GeV}^2$ and Ref. [10] for $q^2 \geq 11 \text{ GeV}^2$ are shown as points with error bars. In most cases, the agreement between the theoretical values and those obtained from data. However, for the ranges $0.1 \le$ $11.0 \le q^2 \le 12.5 \text{ GeV}^2$ and $15 \le q^2 \le 17 \text{ GeV}^2$ there are significant disagreements. It is convenient to define P_1 and P_2 as,

$$\mathsf{P}_1 = \frac{\mathcal{F}_\perp}{\mathcal{F}_\parallel}, \qquad \mathsf{P}_2 = \frac{\mathcal{F}_\perp}{\mathcal{F}_0}. \tag{8}$$

The observables F_{\perp} , F_L , $A_{\rm FB}$, A_5 and A_4 can be written [2] as

$$F_{\perp} = u_{\perp}^2 + 2\zeta \tag{9}$$

$$F_L \mathsf{P}_2^2 = u_0^2 + 2\zeta \tag{10}$$

$$A_{\rm FB}^2 = \frac{9\zeta}{2P_1^2} (u_{\parallel} \pm u_{\perp})^2 \tag{11}$$

$$A_5^2 = \frac{9\zeta}{4\mathsf{P}_2^2} (u_0 \pm u_\perp)^2 \tag{12}$$

$$A_4 = \frac{\sqrt{2}}{\pi P_1 P_2} \left(2\zeta \pm u_0 u_{\parallel} \right) \tag{13}$$

SKIP TODAY

Three form factors here

q^2 range in GeV ²	$V(q^2)$	$A_1(q^2)$	$A_{12}(q^2)$
$0.1 \le q^2 \le 0.98$	$0.704 \pm 0.404 \ (0.81\sigma)$	$0.538 \pm 0.309 \ (0.79\sigma)$	$0.246 \pm 0.141 \ (1.27\sigma)$
$1.1 \le q^2 \le 2.5$	$0.624 \pm 0.081 \ (2.48\sigma)$	$0.384 \pm 0.051 \ (1.42\sigma)$	$0.331 \pm 0.052 \\ (0.72\sigma)$
$2.5 \le q^2 \le 4.0$	$0.318 \pm 0.185 \\ (0.70\sigma)$	$0.204 \pm 0.119 \ (0.89\sigma)$	$0.270 \pm 0.177 \ (1.56\sigma)$
$4.0 \le q^2 \le 6.0$	$0.556 \pm 0.026 \ (1.42\sigma)$	$0.398 \pm 0.020 \ (2.02\sigma)$	$0.359 \pm 0.032 \ (1.28\sigma)$
$6.0 \le q^2 \le 8.0$	$0.597 \pm 0.017 \ (0.83\sigma)$	$0.437 \pm 0.014 \ (2.74\sigma)$	0.394 ± 0.022 (2.18σ)
$11.0 \le q^2 \le 12.5$	$0.172 \pm 0.006 \ (5.65\sigma)$	$0.539 \pm 0.027 \ (2.43\sigma)$	$0.462 \pm 0.028 \ (2.82\sigma)$
$15.0 \le q^2 \le 17.0$	$0.713 \pm 0.004 \ (6.25\sigma)$	$0.638 \pm 0.026 \ (3.36\sigma)$	$0.505 \pm 0.016 \ (4.64\sigma)$
$17.0 \le q^2 \le 19.0$	$1.936 \pm 0.007 \\ (4.38\sigma)$	$0.678 \pm 0.025 \ (3.82\sigma)$	$0.498 \pm 0.014 \ (4.64\sigma)$

TABLE I. The form factor values obtained from fit to 3 fb⁻¹ of LHCb data [4]. Round brackets indicate the standard deviation between fitted values and theoretical estimates [7, 10]. We find significant discrepancies for several values, especially for the large q^2 region.

B \rightarrow K* I⁺ I⁻ form factor ratios determined from data are inconsistent: violate HQET equalities at q^2_{max}

q^2 range in GeV^2	u_0	ı
$15 \le q^2 \le 17$	0.001 ± 0.015	0.013 :
$17 \le q^2 \le 19$	0.137 ± 0.013	0.002
$15 \le q^2 \le 19$	0.068 ± 0.005	0.002 :

TABLE II. The values of u_0 , u_{\parallel} and to $3 \, \text{fb}^{-1}$ of LHCb data [4]. In large equality $u_0 = u_{\parallel} = u_{\perp}$ is expected to he charm loop contributions with negligible value of u_{\parallel} for the larger q^2 bin is unexpany conclusions. Significant discrepanciate to be solely due to non-factorizable chare observed between the values of u_{\perp}

Cannot recover by adding resonances or large non-perturbative effects. But NP works.

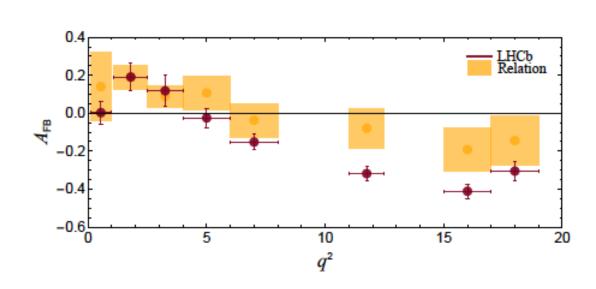


FIG. 3. (color online). The mean values and 1σ uncertainties for $A_{\rm FB}$ calculated using the 'Relation' (Eq. (17)) are shown by yellow (light) bands. We emphasize that these bands are derived using only experimentally measured observables and do not depend on any theoretical input. The error bars in red (dark) correspond to the experimentally measured central values and errors in $A_{\rm FB}$ for the respective q^2 bins. See text for details.

Still confirmation and more data is needed to close the case

A_{FB} (q²) for <u>Inclusive</u> b \rightarrow s l⁺ l⁻

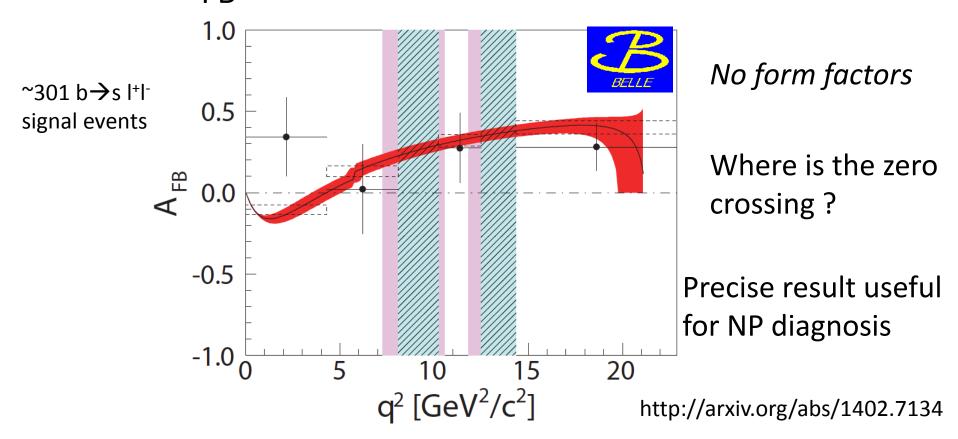


TABLE II. Fit results for the four q^2 bins. For \mathcal{A}_{FB} , the first uncertainty is statistical and the second uncertainty is systematic. \mathcal{A}_{FB} values predicted by the SM [4, 7] are also shown with systematic uncertainties. For the signal yields, only statistical uncertainties are shown. The uncertainties of α and β are due to the statistical uncertainties of the MC.

		1st bin	2nd bin	3rd bin	4th bin
	$(B \to X_s e^+ e^-)$ $(B \to X_s \mu^+ \mu^-)$	[0.2,4.3]	[4.3,7.3] [4.3,8.1]	[10.5,11.8] [10.2,12.5]	[14.3, 25.0]
$\mathcal{A}_{\mathrm{FB}}$		$0.34 \pm 0.24 \pm 0.02$	$0.04 \pm 0.31 \pm 0.05$	$0.28 \pm 0.21 \pm 0.01$	$0.28 \pm 0.15 \pm 0.01$
A_{FB} (theor	ry)	-0.11 ± 0.03	0.13 ± 0.03	0.32 ± 0.04	0.40 ± 0.04
$N_{ m sig}^{ee}$		45.6 ± 10.9	30.0 ± 9.2	25.0 ± 7.0	39.2 ± 9.6
$N_{ m sig}^{ee} \ N_{ m sig}^{\mu\mu} \ lpha^{ee}$		43.4 ± 9.2	23.9 ± 10.4	30.7 ± 9.9	62.8 ± 10.4
$\alpha^{\stackrel{ie}{e}}$		1.289 ± 0.004	1.139 ± 0.003	1.063 ± 0.003	1.121 ± 0.003
$lpha^{\mu\mu}$		2.082 ± 0.010	1.375 ± 0.003	1.033 ± 0.003	1.082 ± 0.003
$oldsymbol{eta}$		1.000	1.019 ± 0.003	1.003 ± 0.000	1.000

TABLE I: Projections for the statistical uncertainties on the $B \to K^{(*)} \nu \bar{\nu}$ branching fractions.

Mode	$\mathcal{B} [10^{-6}]$	Efficiency Belle	711 fb^{-1}	$N_{\rm Sig-exp.} 711 { m fb}^{-1}$	50 ab ⁻¹	50 ab ⁻¹		Total Error
		$[10^{-4}]$	Belle	Belle	Belle II	Belle II	$50 { m ab}^{-1}$	
$B^+ \to K^+ \nu \bar{\nu}$	3.98	5.68	21	3.5	2960	245	23%	24%
$B^0 o K^0_{ m S} u ar{ u}$	1.85	0.84	4	0.24	560	22	110%	110%
$B^+ \to K^{*+} \nu \bar{\nu}$	9.91	1.47	7	2.2	985	158	21%	22%
$B^0 \to K^{*0} \nu \bar{\nu}$	9.19	1.44	5	2.0	704	143	20%	22%
$B \to K^* \nu \bar{\nu}$ combined							15%	17%

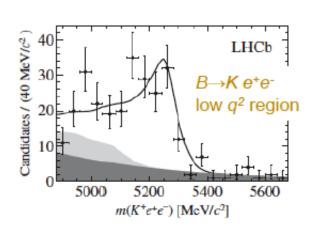
Ans: Verify hint of lepton universality breakdown at Belle II (good electron eff)

Control region gives R_K consistent with unity. Interesting, low q^2 region gives:

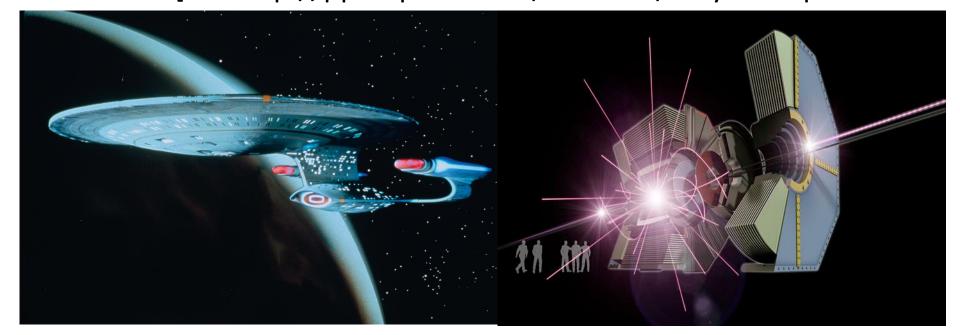
$$R_K = 0.745^{+0.090}_{-0.074}(\text{stat}) \pm 0.036(\text{syst})$$

which is 2.6σ from unity, 3σ if BaBar included.

R. Aaij et al. (LHCb collab); PRL 113, 151601 (2014)



2014 was the 50th anniversary of the discovery of CP violation in the kaon sector [see http://pprc.qmul.ac.uk/research/50-years-cp-violation



The Next Generation Belle II and the LHCb upgrade

US P5 report (p. v): "Explore the unknown: new particles, interactions, and physical principles"

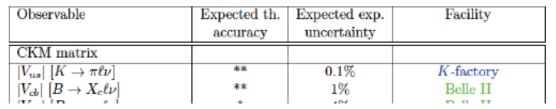
Physics Reach of Belle II and the LHCb upgrade

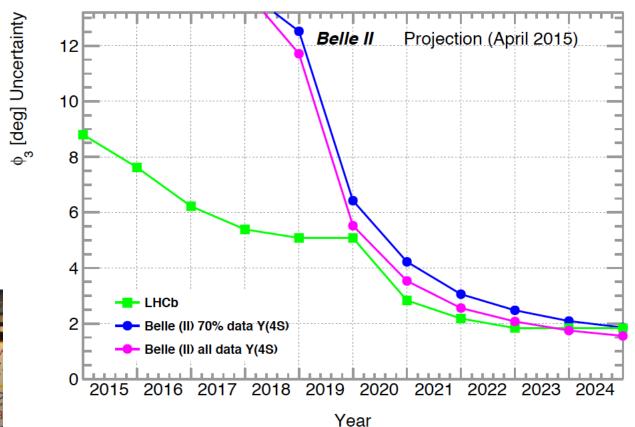
Competition and complementarity



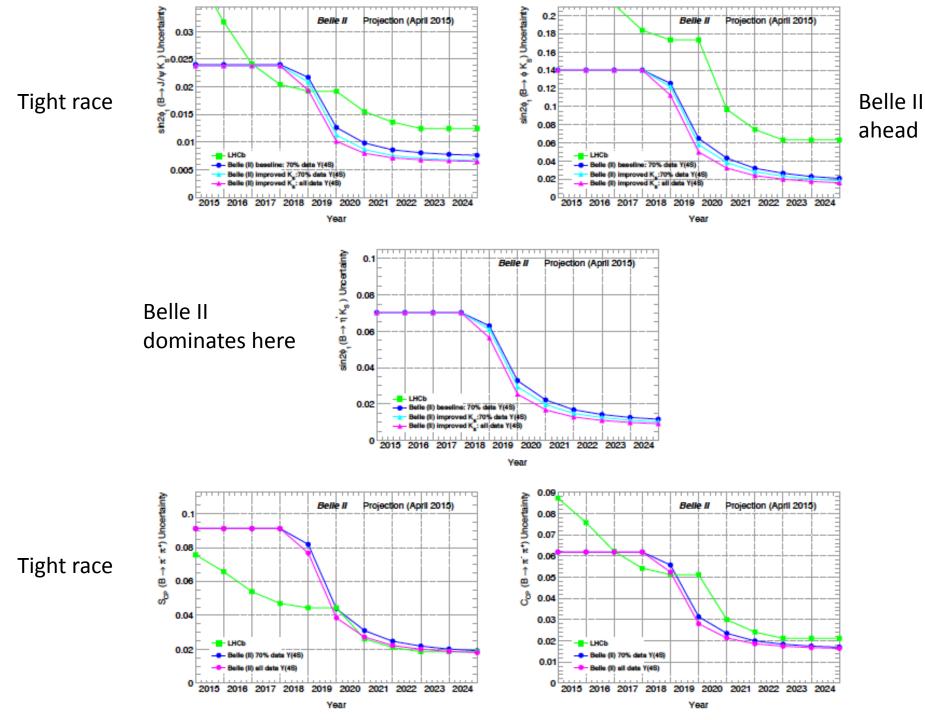


Tofu Gelato?

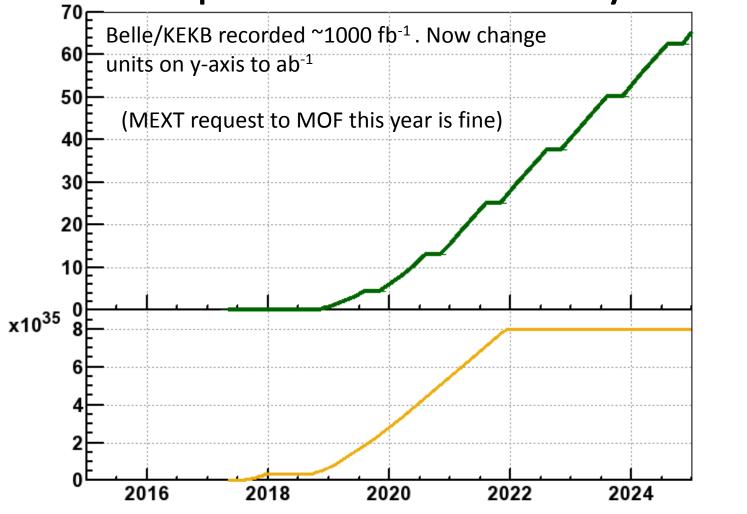




$\mathcal{B}(K \to e\pi\nu)/\mathcal{B}(K \to \mu\pi\nu)$	ककक	0.1%	K-factory
charm and τ			
$B(\tau \rightarrow \mu \gamma)$	***	$3 \cdot 10^{-9}$	Belle II
$ q/p _D$	***	0.03	Belle II
$arg(q/p)_D$	***	1.5°	Belle II
		1	F 2



Latest SuperKEKB Luminosity Profile



N.B. To realize this steep turn-on, requires close cooperation between Belle II and SuperKEKB [and international collaboration on the accelerator].

Also assumes full and stable operation funding profile.

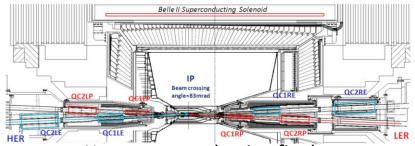
Compare the Parameters for KEKB and SuperKEKB

	KEKB Design	KEKB Achieved : with crab	SuperKEKB Nano-Beam
Energy (GeV) (LER/HER)	3.5/8.0	3.5/8.0	4.0/7.0
β _y * (mm)	10/10	5.9/5.9	0.27/0.30
β _x * (mm)	330/330	1200/1200	32/25
ε _x (nm)	18/18	18/24	3.2/5.3
ε _y /ε _x (%)	1	0.85/0.64	0.27/0.24
σ _y (mm)	1.9	0.94	0.048/0.062
σ_{y}	0.052	0.129/0.090	0.09/0.081
σ_{z} (mm)	4	6 - 7	6/5
I _{beam} (A)	2.6/1.1	1.64/1.19	3.6/2.6
N _{bunches}	5000	1584	2500
Luminosity (10 ³⁴ cm ⁻² s ⁻¹)	1	2.11	80





2015: Basic hardware (except final focus) now in place

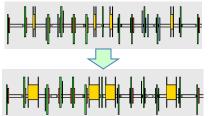


New superconducting final focusing magnets near the IP

e⁻ 2.6A

e+ 3.6A

Redesign the lattice to reduce the emittance (replace short dipoles with longer ones, increase wiggler cycles) (all magnets installed 8/2014)



Replace beam pipes with TiNcoated beam pipes with antechambers (*installed*)



KEKB to SuperKEKB

- lacktriangle Nano-Beam scheme extremely small eta_y^* low emittance
- ◆Beam current X 2

$$L = \frac{g_{\pm}}{2er_{e}} \stackrel{\text{R}}{\circ} 1 + \frac{S_{y}^{*} \stackrel{\text{O}}{\circ} I_{\pm} X_{\pm y}}{S_{x}^{*} \stackrel{\text{O}}{\circ}} \frac{R_{L}}{b_{y}^{*}} \stackrel{\text{O}}{\circ} \frac{R_{L}}{o} \stackrel{\text{O}}{\circ}$$

40 times higher luminosity $2.1 \times 10^{34} \rightarrow 8 \times 10^{35} \text{ cm}^{-2} \text{s}^{-1}$



Reinforce RF systems for higher beam currents

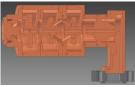
Improve monitors and control system

Injector Linac upgrade

Upgrade positron capture section



Low emittance RF electron gun

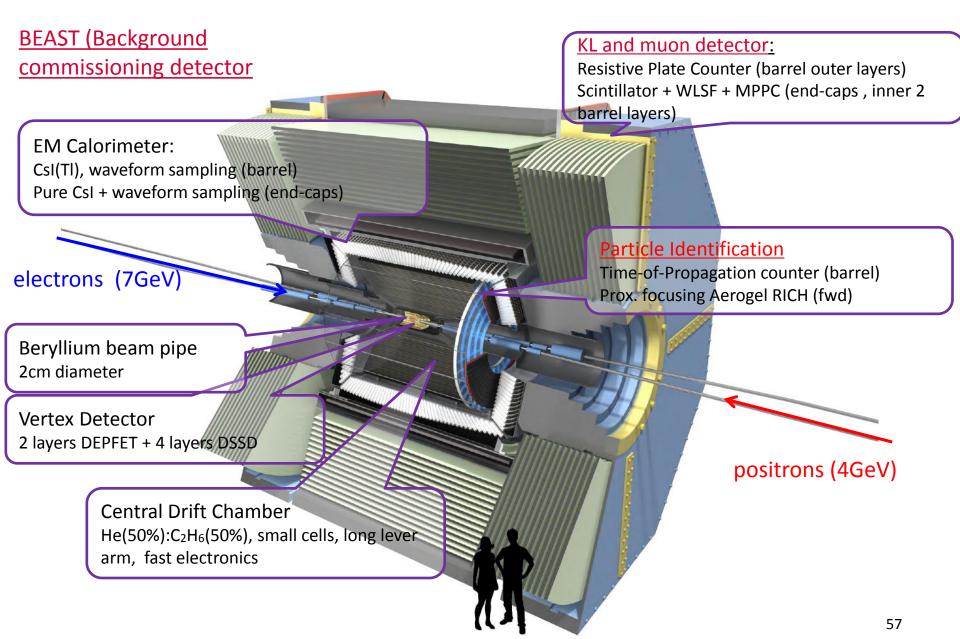






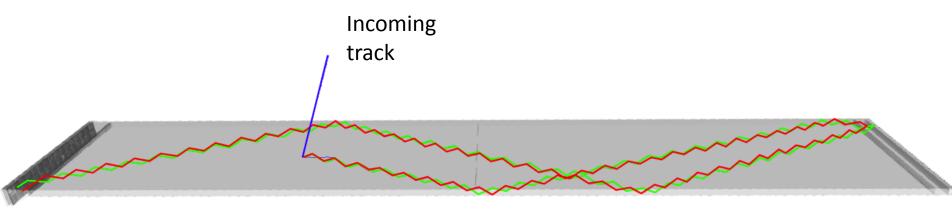
New e+ Damping Ring constructed

Belle II Detector

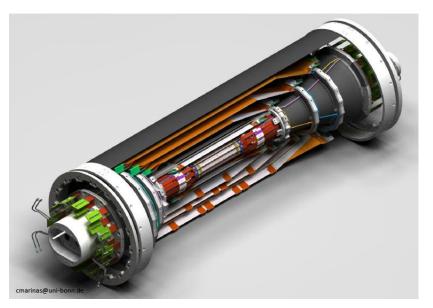


Barrel Particle Identification

A GEANT4 event display of a 2 GeV pion and kaon interacting in a TOP quartz bar. (Japan, US, Slovenia, Italy)



Vertexing/Inner Tracking



```
Beampipe r= 10 mm

DEPFET pixels (Germany, Czech Republic, Spain...)

Layer 1 r=14 mm

Layer 2 r= 22 mm

DSSD (double sided silicon detectors) FWD/BWD

Layer 3 r=38 mm (Australia) Italy

Layer 4 r=80 mm (India)

Layer 5 r=115 mm (Austria)

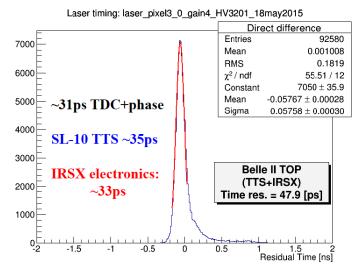
Layer 6 r=140 mm (Japan)
```

Belle II iTOP at Fuji Hall/Hawaii



Module 04 assembly at Fuji Hall

Module 01 assembly at Fuji Hall



Production testing of readout with single photo-electron laser pulses in Hawaii; electronics resolution ~35ps

All quartz and electronics in hand; now testing and assembling.

BEAST Phase 1 Installation Sequence



Install IP bridge: Completed

2. Install 6km cables: Completed June 25-28



Install IP chamber: Completed June 29th



Install support structure and sensors: Aug 17- Sep 22



3.

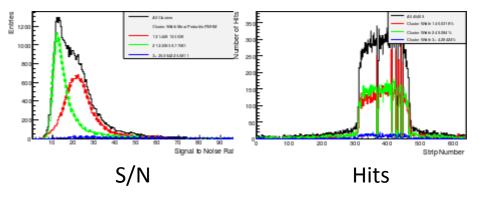


Highlights of Belle II construction

SuperKEKB hardware is being finalized.

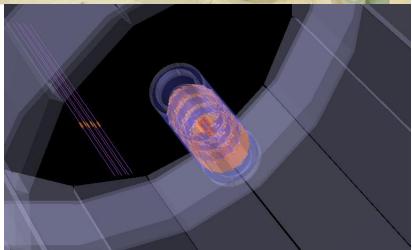


BEAST PHASE I beampipe installed



Final Belle II SVD ladder in CERN beam in June (working well!)





(a) CDC arriving at Tsukuba Hall; (b) first cosmics with partly instrumented electronics (6 layers)

Innovative Technologies in Belle II

Pixelated photo-sensors play a central role



MCP-PMTs in the iTOP HAPDs in the ARICH SiPMs in the KLM



Waveform sampling with precise timing is "saving our butts". Front-end custom ASICs (Application Specific Integrated Circuits) for all subsystems → a 21st century HEP experiment.

KL/muon detector (TARGETX ASIC)

Electromagnetic calorimeter

(New waveform sampling backend with good timing)

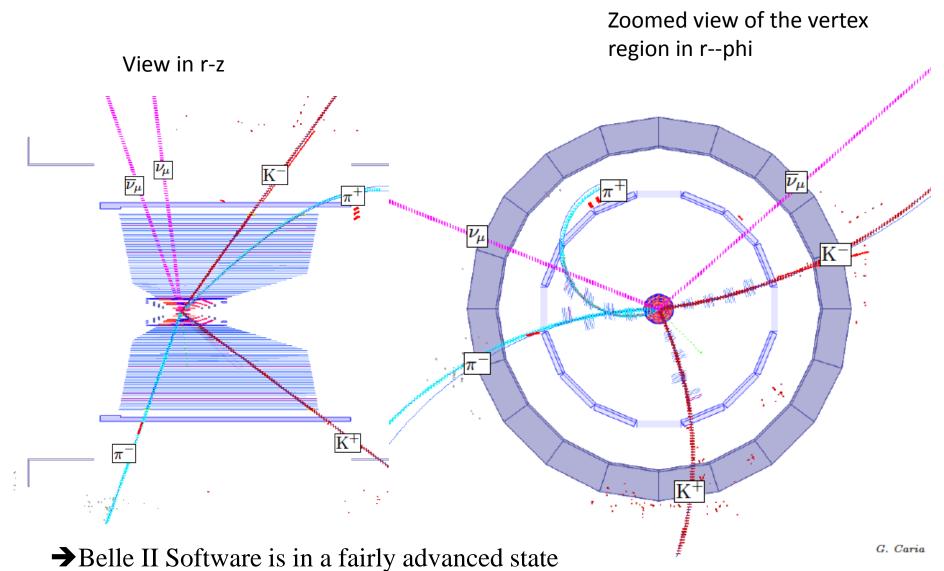
iTOP particle identification (IRSX ASIC)

Aerogel RICH (KEK custom ASIC)

Central Drift Chamber (KEK custom ASIC)

SVD (APV2.5 readout chip adapted from CMS)

"Missing Energy Decay" in a Belle II GEANT4 MC simulation Signal B \rightarrow K $\nu \nu$ tag mode: B \rightarrow D π ; D \rightarrow K π

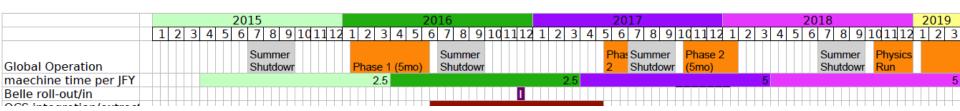


Some Belle II jargon

BEAST PHASE I: Simple background commissioning detector (diodes, TPCs, crystals). No final focus. Only single beam background studies possible [starts in Jan 2016].

<u>BEAST PHASE II</u>: More elaborate inner background commissioning detector. <u>Full Belle II outer</u> <u>detector</u>. Full superconducting final focus. *No vertex detectors*.

Belle II Schedule (Zoom-in on operations)



When do we start Belle II?

BEAST PHASE I: Starts in Jan 2016

BEAST PHASE II: Starts ~May 2017 [some

limited physics without vertex detectors]

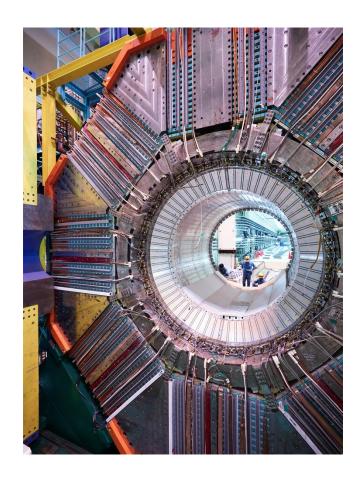
Belle II Physics Running: Fall 2018

Conclusion/Next Generation

- The e+ e- B factories confirmed that the KM phase is responsible for most of the observed CPV [Physics Nobel Prize 2008]
- Nevertheless, 10-20% NP effects are consistent with all current flavor data.
- LHCb has ruled out large CPV phases from NP in the B_s sector.
- "Missing energy B decays" provide important high—mass sensitivity to the charged Higgs in the multi-TeV range.
- LHC exps (CMS, LHCb) measured $B_s \rightarrow \mu^+ \mu^-$
- Angular anomalies in B→K*l+l- from LHCb with 3 fb-1
- <u>Flavor physics is exciting and fundamental.</u> (Did we just find NP? This may be the path for the future of HEP.)

SuperKEKB commissioning starts in January. Belle II Physics runs in 2018 and the LHCb upgrade in ~2020. <u>These facilities will</u> <u>inaugurate a new era of flavor physics and the study of CP violation.</u>

Backup slides



Belle II construction status



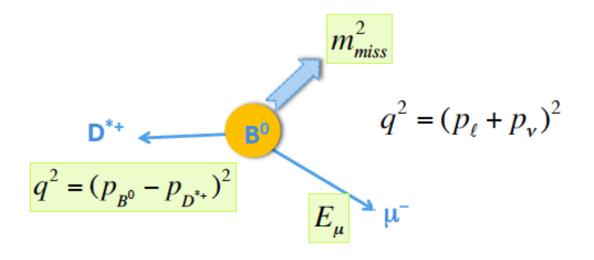
"Tensions are high, tempers are short."

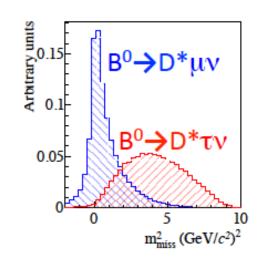
Belle I Drift Chamber and Vertex Detector in the Ueno Science Museum in Tokyo

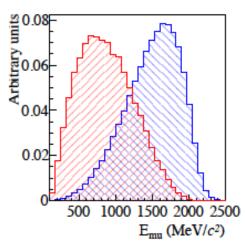


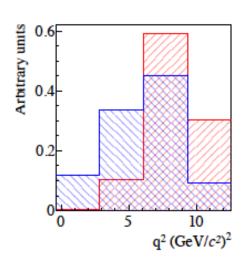
Separating $B^0 \rightarrow D^*\tau\nu$ from $B^0 \rightarrow D^*\mu\nu$

3 key kinematic variables computed in the B rest frame







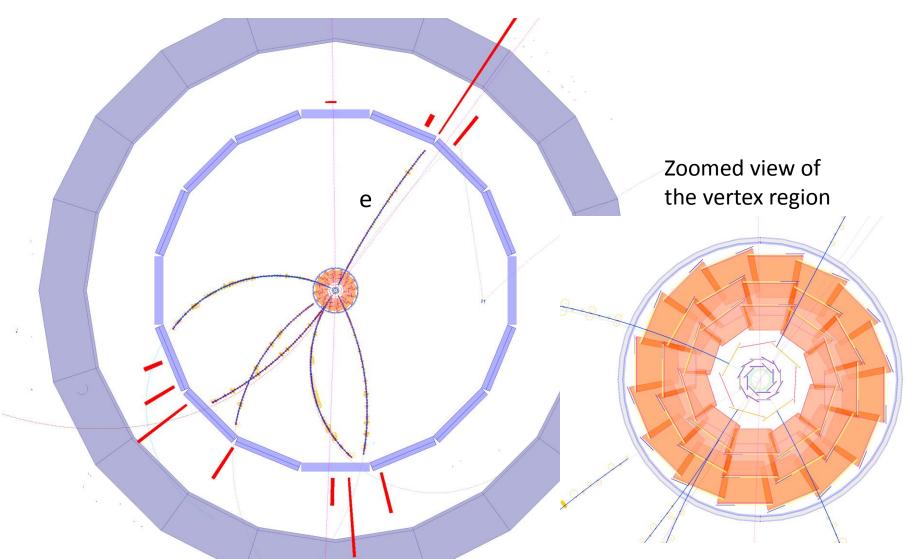


LHCb simulation

approximate B rest frame

"Missing Energy Decay" in a Belle II GEANT4 MC simulation





Updated projections for $B \rightarrow K(*)$ nu nubar modes

TABLE I: Projections for the statistical uncertainties on the $B \to K^{(*)} \nu \bar{\nu}$ branching fractions.

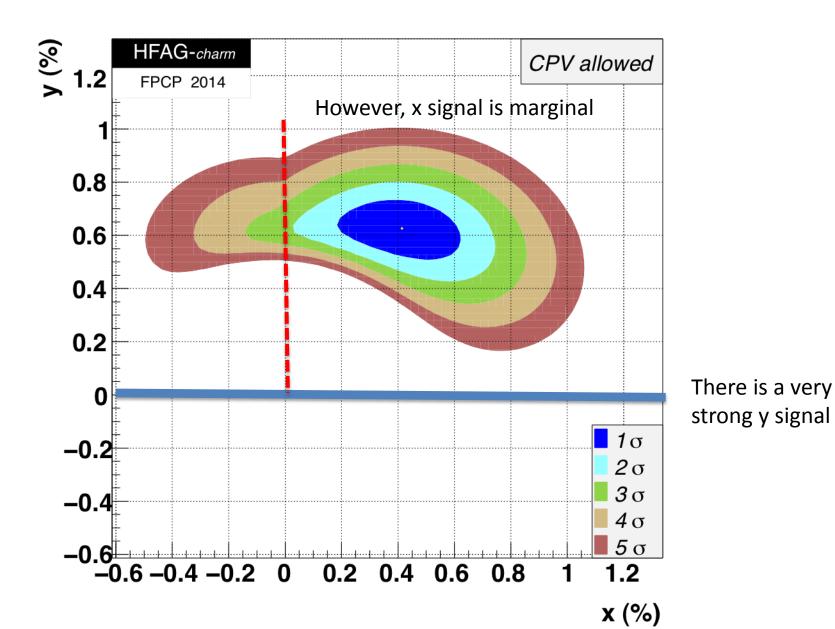
Mode	$\mathcal{B} [10^{-6}]$	Efficiency Belle		$N_{\rm Sig-exp.} 711 { m fb}^{-1}$			Statistical error	Total Error
		$[10^{-4}]$	Belle	Belle	Belle II	Belle II	50 ab^{-1}	
$B^+ o K^+ u \bar{ u}$	3.98	5.68	21	3.5	2960	245	23%	24%
$B^0 o K^0_{ m S} u ar{ u}$	1.85	0.84	4	0.24	560	22	110%	110%
$B^+ \to K^{*+} \nu \bar{\nu}$	9.91	1.47	7	2.2	985	158	21%	22%
$B^0 \to K^{*0} \nu \bar{\nu}$	9.19	1.44	5	2.0	704	143	20%	22%
$B \to K^* \nu \bar{\nu}$ combined							15%	17%

A. J. Buras, J. Girrbach-Noe, C. Niehoff and D. M. Straub, JHEP 1502, 184 (2015) [arXiv:1409.4557 [hep-ph]].

O. Lutz et al. [Belle Collaboration], Phys. Rev. D 87, no. 11, 111103 (2013) [arXiv:1303.3719 [hep-ex]].

^[3] T. Kuhr, " $B \rightarrow h^{(*)} \nu \bar{\nu}$ ", KEK-FF Workshop (2013).

Mixing and CP violation in the D system



D mixing: Another new physics phase!

$$\varphi \sim \frac{2\eta A^2 \lambda^5}{\lambda} \sim O(10^{-3})$$

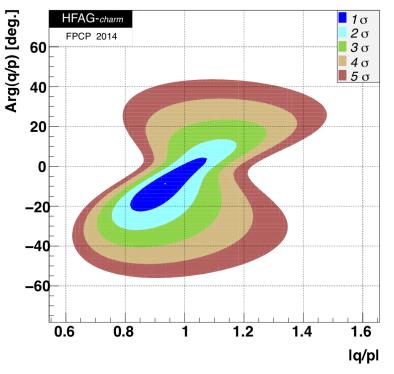
CPV in D system negligible in SM

CPV in interf. mix./decay:

$$\operatorname{Im} \frac{q}{p} \frac{\overline{A}_f}{A_f} \equiv (1 + \frac{A_M}{2})e^{i\varphi} \neq 0; \varphi \neq 0$$

The existence of D mixing (if x is non-zero) allows us to look for another poorly constrained new physics phase but this time from up-type quarks.

(c.f. CPV in B_s mixing)

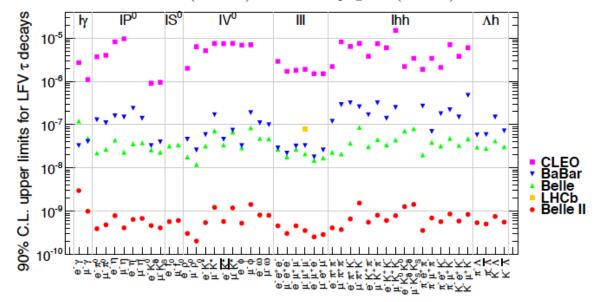


Current WA sensitivity $\sim \pm 20^{\circ}$, 50 ab⁻¹ go below 2°

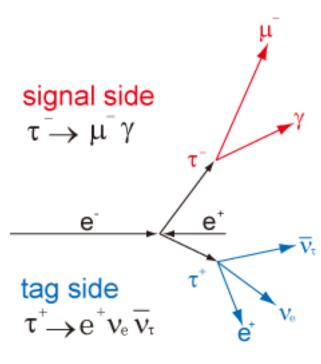


Tau Lepton Flavor Violation

FIG. 5: LFV UL (90% C.L.) results from CLEO, BaBar and Belle, and extrapolations for Belle II (50 ab⁻¹) and LHCb updgrade (50 fb⁻¹).



Example of the decay topology



Belle II will push many limits below 10⁻⁹; LHCb has very limited capabilities.

CPV in the charged lepton sector

 There is mixing in the neutrino (neutral lepton) sector. CP violation is possible too.

BaBar rate anomaly ??

Is this interesting or ruled out by other constraints?

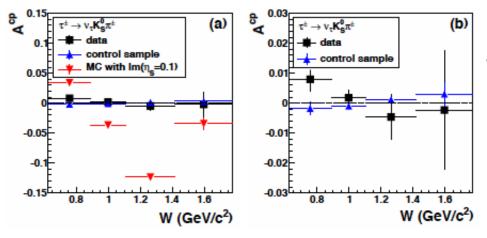


FIG. 2. (a) Measured CP violation asymmetry after background subtraction (squares). The vertical error bars are the statistical error and systematic errors added in quadrature. The CP asymmetry measured in the control sample is indicated by the blue triangles (statistical errors only) and the inverted red triangles show the expected asymmetry for $\Im(\eta_S) = 0.1 \ [\Re(\eta_S) = 0]$. (b) Expanded view (the vertical scale is reduced by a factor of five).

Can we explore at Belle II?

Theoretical predictions for $\Im(\eta_S)$ can be given in context of a MHDM with three or more Higgs doublets [4, 5]. In such models η_S is given by [12]

$$\eta_S \simeq \frac{m_\tau m_s}{M_{H^\pm}^2} X^* Z \tag{10}$$

if numerically small terms proportional to m_u are ignored. Here, $M_{H^{\pm}}$ is the mass of the lightest charged Higgs boson and the complex constants Z and X describe the coupling of the Higgs boson to the τ and ν_{τ} and the u and s quarks, respectively (see [5, 12]). The limit $|\Im(\eta_S)| < 0.026$ is therefore equivalent to

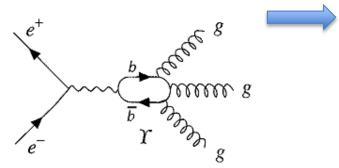
$$|\Im(XZ^*)| < 0.15 \frac{M_{H^{\pm}}^2}{1 \,\text{GeV}^2/c^4}.$$
 (11)

M. Bischofberger et al, Phys. Rev. Lett. 107, 131801 (2011)

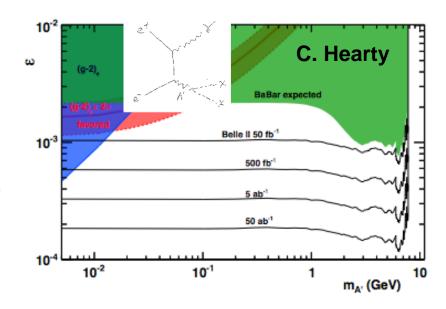
Beast Phase II & New Triggers

- Update to First-physics report: <u>BELLE2-NOTE-PH-2015-003</u> Y(2S), Y(3S), Y(6S), Scan proposals
- Beast Phase II Physics Task Force formed to study physics with this configuration (B. Fulsom).

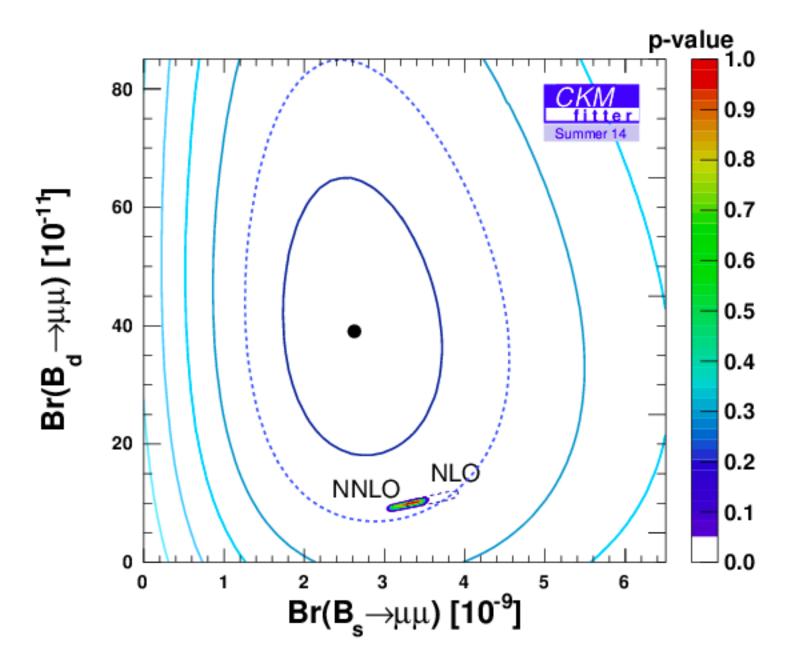
Belle Y(1S) decay data used for Pythia 8
 MC tuning in Belle II (U. Tamponi).



 HLT & L1 Trigger Menu under design. Evolving <u>Trigger Menu</u> (Link).



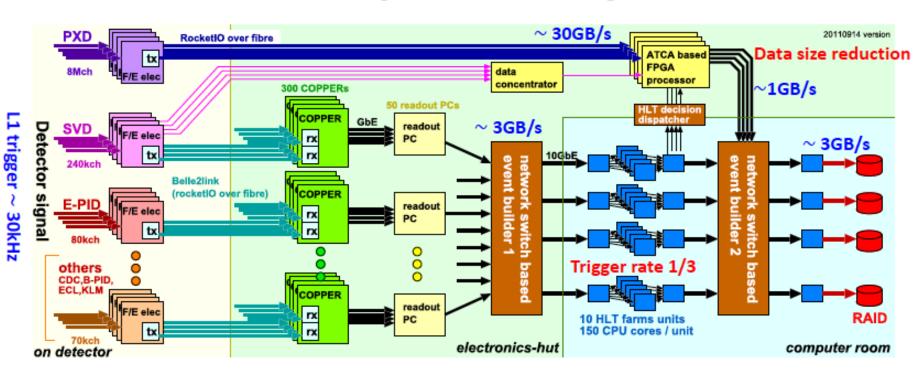
Trigg	gers	Some Ideas C-H. Li				
Single Photon (γ)		Cascade: different thresholds with separate pre-scale factors Use different pre-scale factors for Barrel and Endcap				
e+e-		 two Bhabha triggers, "accept" and "veto" "accept": flattening scheme "veto": 2D→3D ECL Bhabha is being investigated salvage: retain a pre-scaled sample of physics triggers without veto 				
μ+μ-		independent CDC and KLM triggers for luminosity systematics				
γγ		reduce pre-scale to 10 instead of 100				
	γe ⁺ e ⁻ [hlt]	dedicated triggers for calibration (CDC,ECL)				
γ+ 2 trks	γμ+μ-	dedicated triggers for detectors study (CDC, ECL, KLM)				
	γh+h-	 high efficiency for all γ energies and h⁺h⁻ invariant masses one high energy cluster in ECL, one track in opposite hemisphere 				
Additional trigger information		CDC-TOP-ECL-KLM Matching More detectors information				



Belle II

Belle II

Data acquisition system



<u>Discussion Topic</u>: What additional *Theoretical Work* is required to determine whether NP is present in B decays?

<u>Participants</u>: Wolfgang Altmannshofer, Christoph Bobeth, Jorge Martin Camalich, Robert Fleischer, Zoltan Ligeti, Rahul Sinha



Nagoya FPCP15 roundtable

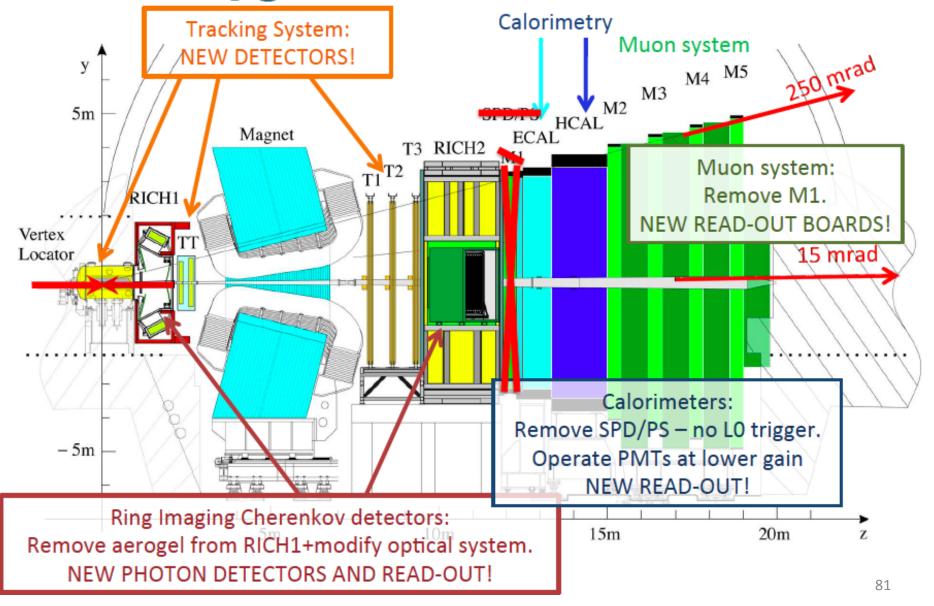




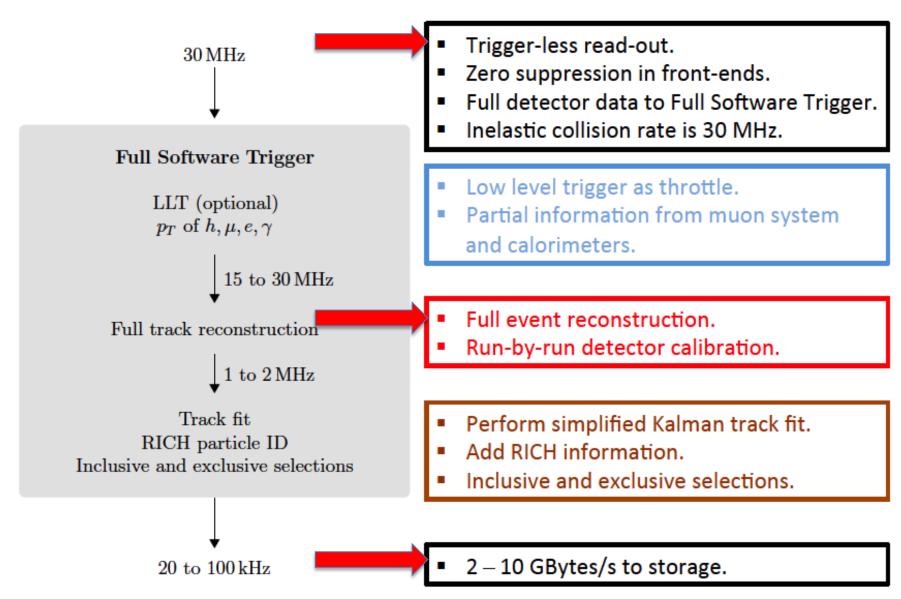
Discussed in Nakada-san's talk



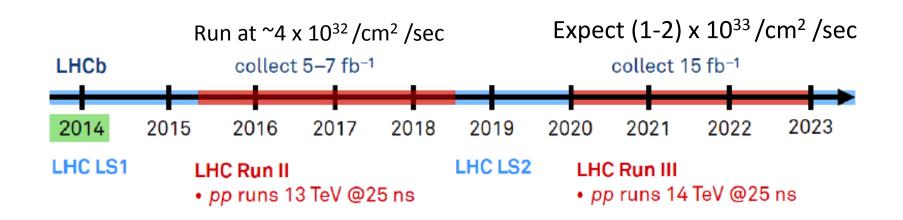
Upgraded LHCb detector



LHCb Upgrade: Key Feature is Trigger-less readout



LHCb upgrade timeline



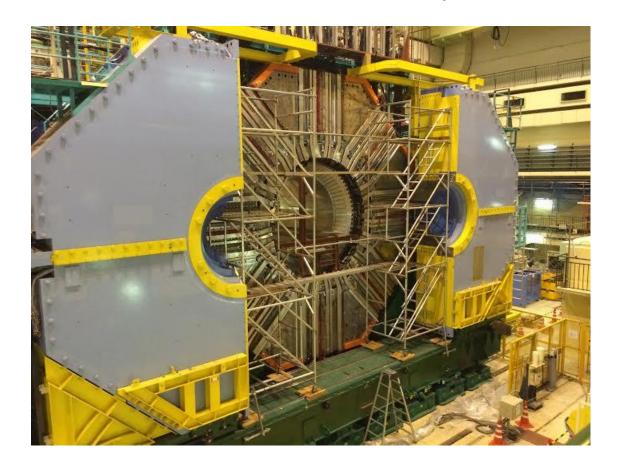
HL-LHC

- Collect 50 fb⁻¹ after upgrade.
- Continue taking data during HL-LHC.

Upgraded trigger and DAQ is the key feature



Belle II@Tsukuba Hall, KEK



Belle II detector with barrel KLM upgrade as well as forward and backward muon endcap upgrades now installed.

Belle II at IPMU in 柏の葉, Japan

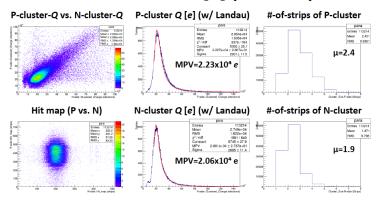
Constructing two layers of the Belle II SVD detector in the clean room on the 1st floor. Dr T. Higuchi is the leader.

Japan (Layer 6) and India/Tata Institute (Layer 4)

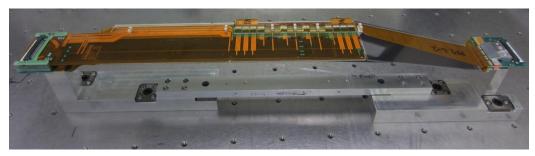
Test Production by late Sept;
Detector production starts ~ Nov 2015



90Sr Source Test [2] (SBW990)



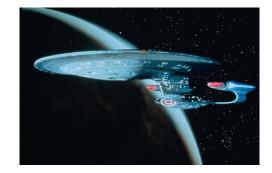




L4 mechanical prototype

New Reference for the Next Generation

The Physics of the B Factories http://arxiv.org/abs/1406.6311



This work is on the Physics of the B Factories. Part A of this book contains a brief description of the SLAC and KEK B Factories as well as their detectors, BaBar and Belle, and data taking related issues. Part B discusses tools and methods used by the experiments in order to obtain results. The results themselves can be found in Part C.

Comments: 928 pages

Subjects: High Energy Physics - Experiment (hep-ex); High

Energy Physics - Phenomenology (hep-ph)

Report number: SLAC-PUB-15968, KEK Preprint 2014-3

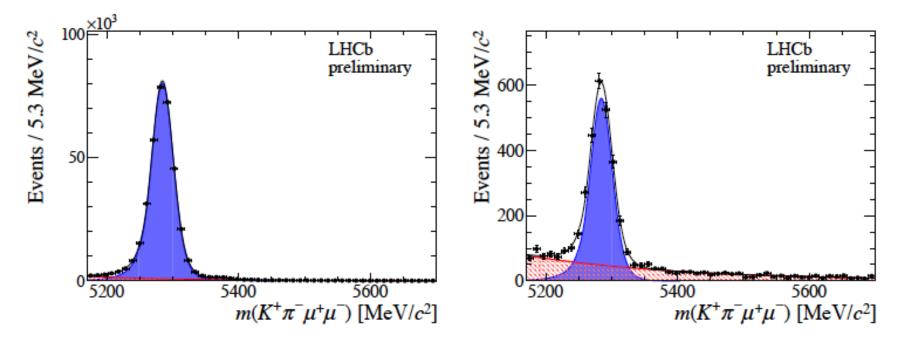
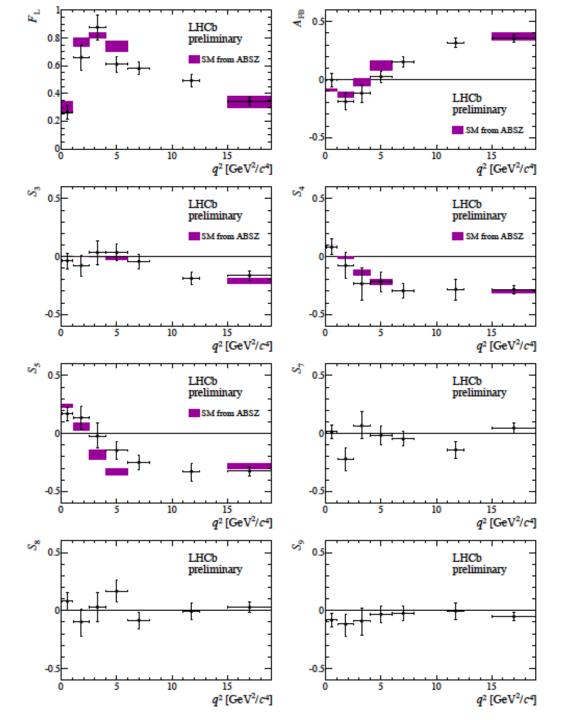
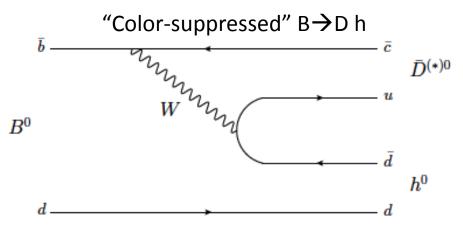


Figure 3: Invariant mass $m(K^+\pi^-\mu^+\mu^-)$ for (left) the control decay $B^0 \to J/\psi K^{*0}$ and (right) the signal decay $B^0 \to K^{*0}\mu^+\mu^-$, integrated over the full q^2 range. The $B^0 \to K^{*0}\mu^+\mu^-$ signal yield integrated over q^2 is determined to be 2398 ± 57 . Overlaid are the projections of the total fitted distribution (black line) and its different components. The signal is shown by the blue component and the background is shown by the red hatched component.



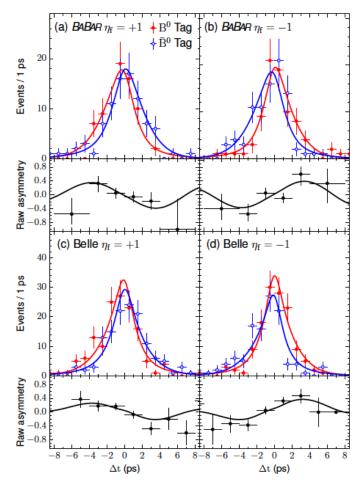
B factories: Check CP violation in $b \rightarrow c$ [ubar d] processes

2015: First joint BaBar-Belle data analysis м. Rohrken et al



where D^0 is a CP eigenstate and $h^0=\pi^0$, η , ω

Combining Belle and BaBar datasets, ~1260 signal events, obtain a 5.4 σ CP violation signal \rightarrow First observation sin(2 β_{eff})=0.66±0.10(stat)±0.06(sys)



Phase of V_{td} again

Conclusion: CP violation in $b \rightarrow c$ ubar d modes is the same as in $b \rightarrow c$ cbar s modes (e.g. $B \rightarrow J/\psi K_S$)

Efficiency at low q² for Belle II?

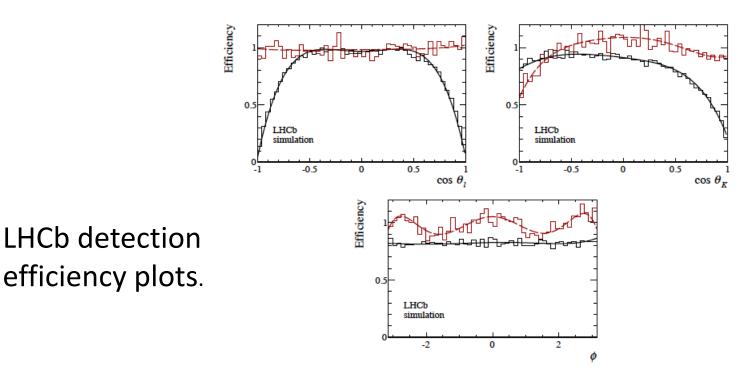


Figure 2: Angular efficiency in $\cos\theta_{\ell}$, $\cos\theta_{K}$ and ϕ , as determined from a principal moment analysis of simulated three-body $B^0 \to K^{*0} \mu^+ \mu^-$ phase-space decays. The efficiency is shown for the regions $0.1 < q^2 < 0.98 \, \text{GeV}^2/c^4$ (black solid line) and $18.0 < q^2 < 19.0 \, \text{GeV}^2/c^4$ (red dashed line). The histograms indicate the distribution of simulated three-body $B^0 \to K^{*0} \mu^+ \mu^-$ phase-space decays used to determine the acceptance. The absolute normalisation of the dsitributions is arbitrary.

How important are the di-electron modes that Belle II does well? (see the effect of the <u>photon pole</u> more clearly; NP from right handed currents??)

More backup

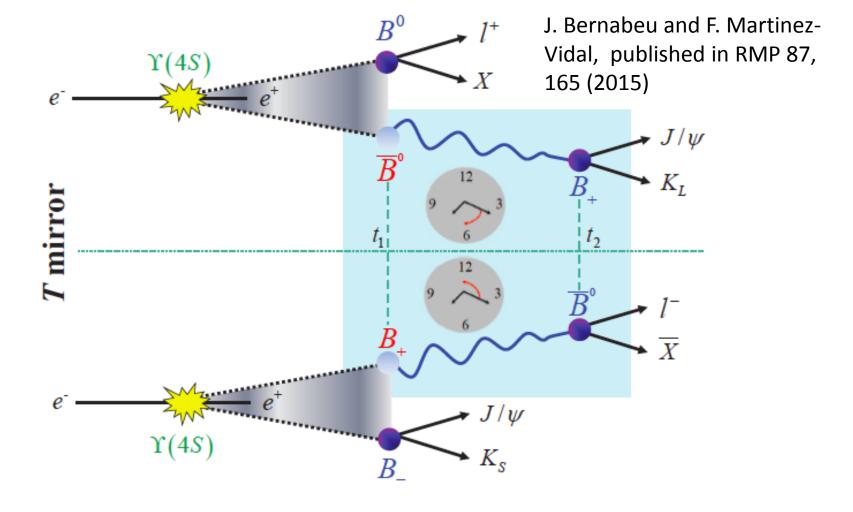
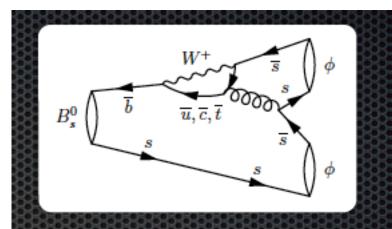


FIG. 11. Foundations of the time-reversal experiment. Electron-positron collisions at the asymmetric B factory produce $\Upsilon(4S)$ resonances, each of which decays through strong interaction in an entangled pair of B mesons. When one B meson decays at t_1 , the identity of the other is "tagged" without measuring it specifically. In the top panel, the B meson observed to decay to the final state $\ell^+ X$ at t_1 transfers information to the (still living) partner meson and dictates that it is in a $\overline{B}{}^0$ state. This surviving meson tagged as $\overline{B}{}^0$ is observed later at t_2 , encapsulating a time ordering, to decay into a final state $J/\psi K_L^0$ that filters the B meson to be in a B_+ state, a linear combination of B^0 and $\overline{B}{}^0$ states. This case corresponds to a transition $\overline{B}{}^0 \to B_+$. To study time reversal we have to compare the rate at which this transition occurs to the rate of the time-reversed transition, $B_+ \to \overline{B}{}^0$ (bottom panel). Adapted from².

M.Needham@ICHEP2014

But LHCb dominates on these B_s modes



 $B_s \rightarrow \varphi \varphi$ - Time-Dependent Results Candidates / (0.32 ps Candidates / (0.42 rad LHCb Preliminary LHCb Preliminary 300 250 200 150 10-1 10-2 10 10 Decay time [ps] Φ [rad] Candidates / (0.13) 350 LHCb Preliminary LHCb Preliminary 300 Total 250 200 200 CP-even 150 100 CP-odd 100 S-wave -0.5 0.5 -0.5 0.5 cos0, cos0. Projections are s-weighted and include acceptances, Decay time acceptance from B_s→D_sπ data,

Angular acceptance from simulated events.



Signal model

We use the same notations as LHCb [arXiv:1304.2600]:

$$\frac{d^4\Gamma(B_s(t))}{d\Theta dt} = X(\Theta, \alpha, t) = \sum_{i=1}^{10} O_i(\alpha, t).g_i(\Theta),$$

$$O_i(\alpha, t) = N_i e^{-\Gamma_S t} \left[a_i \cosh(\frac{1}{2}\Delta\Gamma_S t) + b_i \sinh(\frac{1}{2}\Delta\Gamma_S t) + c_i \cos(\Delta m_S t) + d_i \sin(\Delta m_S t) \right]$$

i	$g_i(heta_T, \psi_T, \phi_T)$	Ni	aį	b_i	c _i	d_i
1	$2\cos^2\psi_T(1-\sin^2\theta_T\cos^2\phi_T)$	$ A_0(0) ^2$	1	D	С	-5
2	$\sin^2 \psi_T (1 - \sin^2 \theta_T \sin^2 \phi_T)$	$ A_{ }(0) ^2$	1	D	С	_S
3	$\sin^2\psi_T\sin^2 heta_T$	$ A_{\perp}^{"}(0) ^{2}$	1	-D	С	S
4	$-\sin^2\psi_{\mathcal{T}}\sin2 heta_{\mathcal{T}}\sin\phi_{\mathcal{T}}$	$ A_{ }(0)A_{\perp}(0) $	$C\sin(\delta_{\perp}-\delta_{\parallel})$	$S\cos(\delta_{\perp}-\delta_{\parallel})$	$\sin(\delta_{\perp}-\delta_{\parallel})$	$D\cos(\delta_{\perp} - \delta_{\parallel})$
5	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin^2 \theta_T \sin 2\phi_T$	$ A_0(0)A_{ }(0) $	$\cos(\delta_{\parallel} - \delta_0)$	$D\cos(\delta_{\parallel}-\delta_0)$	$C\cos(\delta_{\parallel}-\delta_0)$	$-S\cos(\delta_{\parallel}-\delta_0)$
6	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin 2\theta_T \sin \phi_T$	$ A_0(0)A_{\perp}(0) $	$C\sin(\delta_{\perp}-\delta_0)$	$S\cos(\delta_{\perp}-\delta_0)$	$\sin(\delta_{\perp}-\delta_0)$	$D\cos(\delta_{\perp}-\delta_0)$
7	$\frac{2}{3}(1-\sin^2\theta_T\cos^2\phi_T)$	$ A_S(0) ^2$	1	-D	С	5
8	$\frac{1}{3}\sqrt{6}\sin\psi_T\sin^2\theta_T\sin2\phi_T$	$ A_{S}(0)A_{ }(0) $	$C\cos(\delta_{\parallel}-\delta_S)$	$S\sin(\delta_{\parallel}-\delta_S)$	$\cos(\delta_{\parallel} - \delta_S)$	$D\sin(\delta_{\parallel}-\delta_S)$
9	$\frac{1}{3}\sqrt{6}$ sin ψ_T sin $2\theta_T$ cos ϕ_T	$ A_{S}(0)A_{\perp}^{"}(0) $	$\sin(\delta_{\perp}^{"} - \delta_{S})$			$S\sin(\delta_{\perp}^{"}-\delta_{S})$
10	$\frac{4}{3}\sqrt{3}\cos\psi_{\mathcal{T}}(1-\sin^2\theta_{\mathcal{T}}\cos^2\phi_{\mathcal{T}})$	$ A_{S}(0)A_{0}(0) $	$C\cos(\delta_0-\delta_S)$	$S\sin(\delta_0-\delta_S)$	$\cos(\delta_0 - \delta_S)$	$D\sin(\delta_0-\delta_S)$

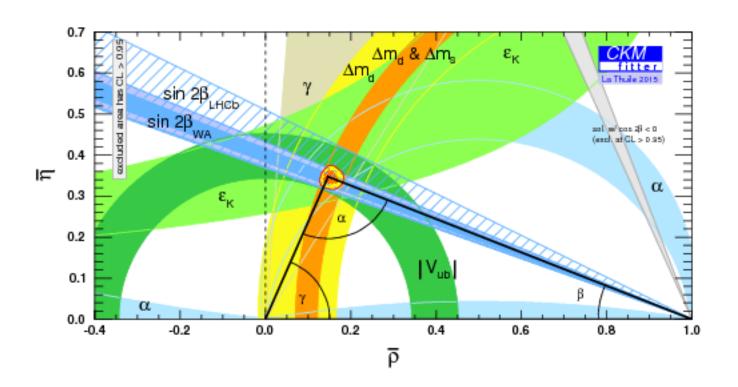
$$C = \frac{1 - |\lambda|^2}{1 + |\lambda|^2}, \qquad S = -\frac{2|\lambda|\sin\phi_s}{1 + |\lambda|^2}, \qquad D = -\frac{2|\lambda|\cos\phi_s}{1 + |\lambda|^2}$$

 $|\lambda|$ includes possible contribution from CP violation in direct decay, we assume $|\lambda|=1$ and we assign a systematics.

 $\Delta\Gamma_s>0$: we use previous LHCb results. α physics parameters $(\Delta\Gamma_s,\,\phi_s,\,c au,\,|A_0|^2,\,|A_s|^2,\,|A_s|^2,\,|A_s|^2,\,|A_s|^2,\,|A_s|^2$



CKMFitter with LHCb sin(2β) included



Efficiency at low q² for Belle II?

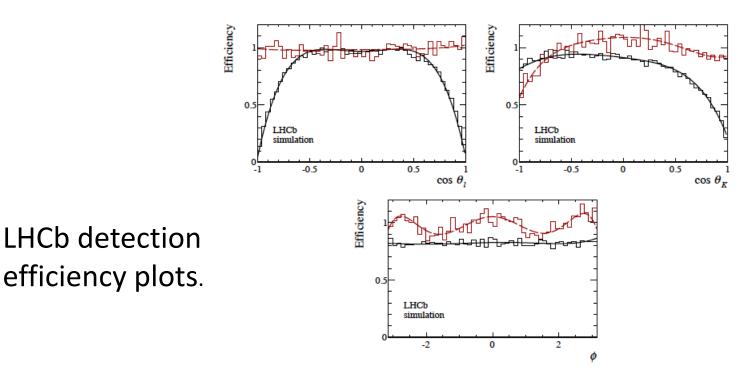
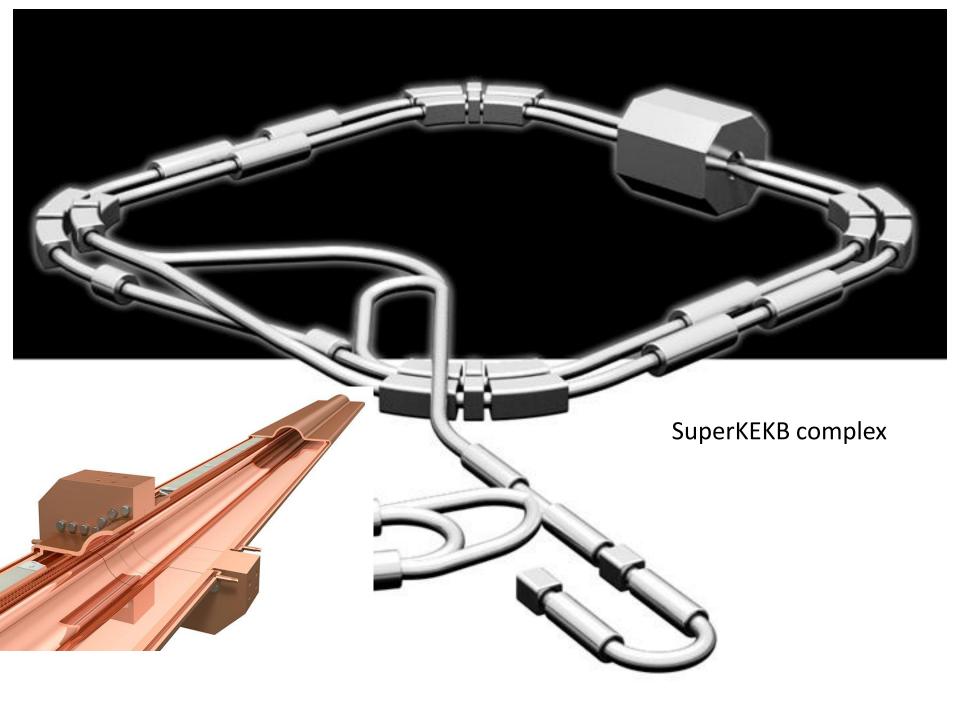
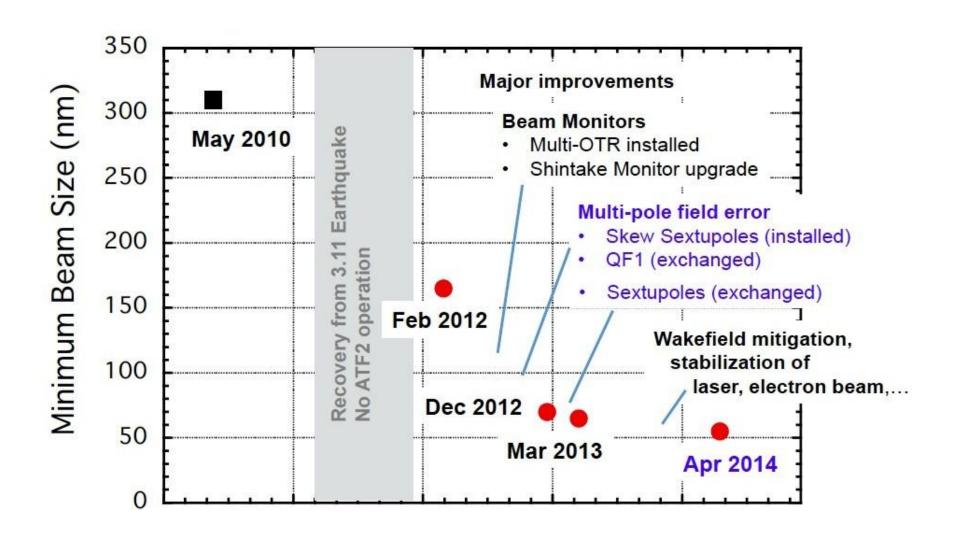


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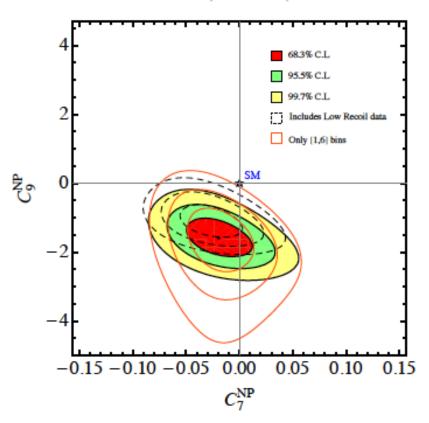


ATF2 nanobeams

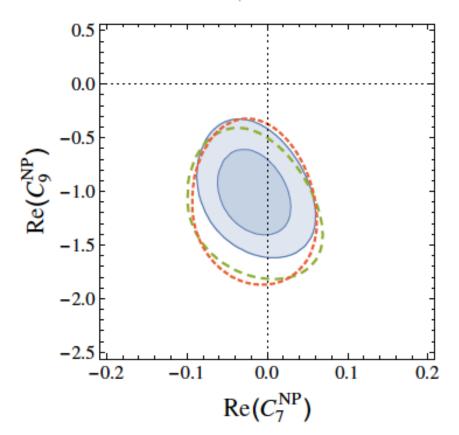


NP Fits

Descotes-Genon, Matias, JV 1307.5683



Altmannshofer, Straub 1503.06199



"Missing Energy" Decays





Tsutentaku tower, Osaka