



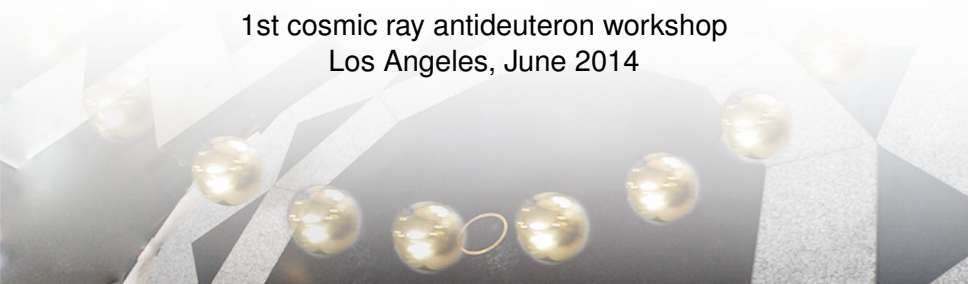
# Hadronization Dependence in Antideuteron Production

Based on arXiv:1207.4560 [hep-ph], arXiv:1402.6259 [hep-ph]

Lars A. Dal

Department of Physics, University of Oslo

Antideuteron 2014  
1st cosmic ray antideuteron workshop  
Los Angeles, June 2014





- Formation of atomic nuclei not handled by Monte Carlos. Coalescence model currently state of the art in computing the antideuteron flux
  - Simple model: Nucleons with  $\Delta p < p_0$  coalesce to form a nucleus
  - Ibarra, Wild: Additional condition: Close in position space – weakly decaying particles considered stable [arXiv:1209.5539](https://arxiv.org/abs/1209.5539) [hep-ph]
  - $p_0$  calibrated against experimental data, large spread in best fit  $p_0$ -values



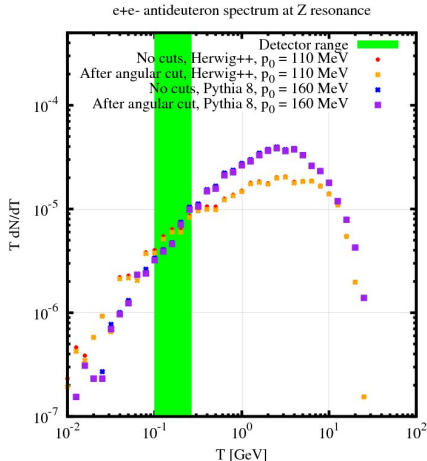
Best fit  $p_0$ -values [MeV] for various experiments

Experiment	Process	Pythia 6	Pythia 8	Herwig++
ALEPH	$e^+e^-$	—	192	159
CLEO	$e^+e^-$	—	133	145
ZEUS	$ep$	236	—	150
CERN ISR	$pp$	—	152	221
ALICE	$pp$	230	—	154

Table from arXiv:1402.6259 [hep-ph]. Pythia 6/8 values are from arXiv:1209.5539 [hep-ph].

- Why the difference between experiments, and why the difference between the Monte Carlos?

# Calibration: A closer look

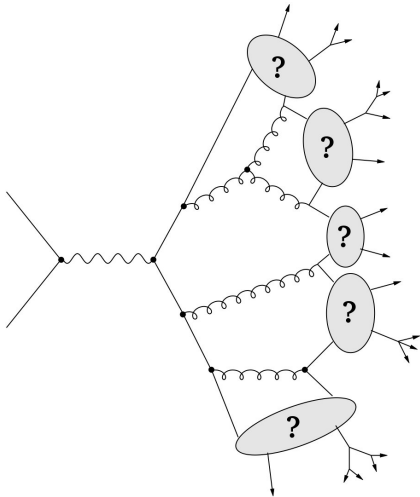


- Calibration: ALEPH ( $e^+e^- \rightarrow Z$ ):  
Herwig++:  $p_0 = 110$  MeV, Pythia:  
160 MeV [arXiv:1207.4560](https://arxiv.org/abs/1207.4560) [hep-ph] \*
- Isotropic coalescence:  
 $dN/dT \propto p_0^3$ ;  $p_0$  only gives the  
normalization
- No calibration of  $p_0$  can make the  
shapes of the spectra agree
- Problem: 2-particle correlations

\* Note: Weak decays were included, thus the low numeric values of  $p_0$

# The issue of hadronization

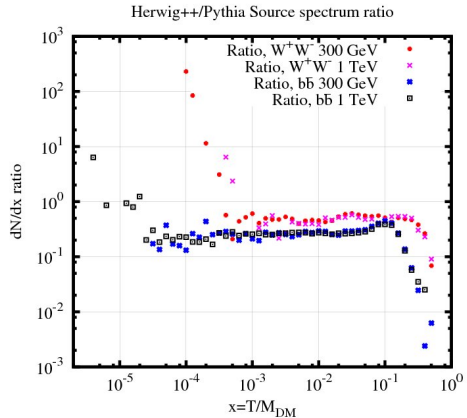
- $p_0 \sim 100 \text{ MeV} \lesssim \Lambda_{\text{QCD}}$ , sensitive to hadronization effects
- Perturbation theory for QCD breaks down at low energies, must resort to phenomenological models
- Monte Carlo: Several free parameters in these models tuned to fit experimental data
- Not specifically tuned to produce correct (anti)nucleon spectra



## Uncertainty on spectrum due to hadronization?

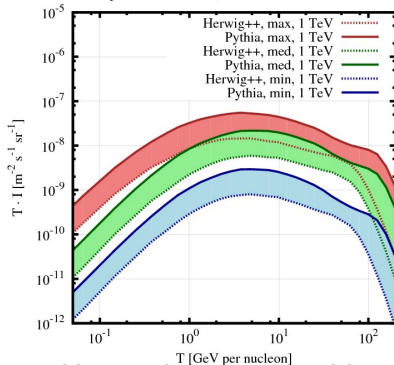
Dal, Kachelrieß arXiv:1207.4560 [hep-ph]

- ALEPH calibration:  
Herwig++:  $p_0 = 110$  MeV,  
Pythia 8:  $p_0 = 160$  MeV
- Comparison of  
antideuteron spectra  
generated with Herwig++  
and Pythia
- Large discrepancies,  
especially at high and low  
energies

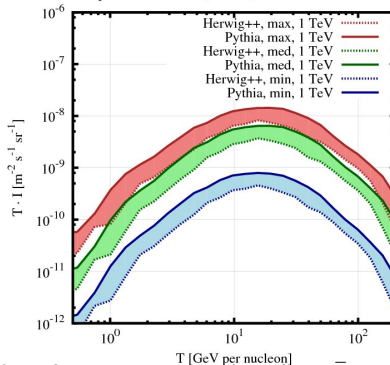


# Uncertainty in the final flux

Spectrum near Earth,  $b\bar{b}$ , with solar modulation



Spectrum near Earth,  $W^+W^-$ , with solar modulation



- Uncertainty comparable to that from propagation for  $b\bar{b}$  at high energies
- Uncertainty induced by the discrepancy seen at low  $x = T/M_{\text{DM}}$  expected to appear in the  $W^+W^-$  channel for higher DM masses



- The idea: Tune hadronization parameters specifically to reproduce antideuteron spectrum
- Uncertainties in the parameters allow us to find corresponding uncertainty on antideuteron flux
- What if we break processes that we don't tune against?
- Tuning Herwig++: Dal, Raklev arXiv:1402.6259 [hep-ph]
  - Re-tune most important Herwig++ hadronization parameters together with  $p_0$
  - Tune against antideuteron spectra from ALEPH ( $e^+e^- \rightarrow Z$ ), ZEUS ( $ep$ ) and CLEO ( $\Upsilon(1S)$  decay)
  - Also tune against (anti)proton spectra from ALEPH and OPAL for consistency
  - 4-dimensional parameter space, each parameter point costs  $\sim 120$  CPU core hours



# Best Fit Parameters



Some 40000 CPU core hours later...

Parameter	Default value	Best fit value	Uncertainty ( $1\sigma$ )*
$p_0$ [MeV]	—	143.2	+6.2 −5.5
ClMaxLight	3.25	3.03	+0.18 −0.15
PSplitLight	1.20	1.31	+0.19 −0.32
PwtDIquark	0.49	0.48	+0.15 −0.04

Best fit  $\chi^2/\text{d.o.f} = 10.6/14.2$

- Highly correlated parameters, challenging to locate best fit point
- Default parameters are reasonably close to best fit point

\* Non-parabolic uncertainty calculated using the MINOS algorithm in Minuit

# Application: Gravitino Dark Matter



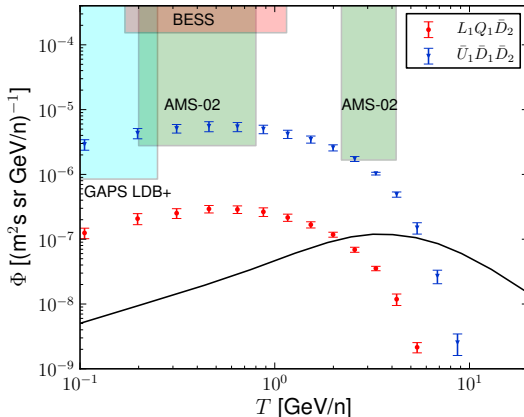
- Gravitino: Supersymmetric partner of the graviton
- R-parity violation: Gravitino unstable but long lived, good DM candidate
- RPV operators of interest:  $\lambda'_{ijk} L_i Q_j \bar{D}_k$ ,  $\lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k$
- $\Phi_{\tilde{d}} \propto \Gamma \propto \lambda^2$ ; fluxes can easily be re-scaled to any value of  $\lambda$
- Goal: Set limits on trilinear RPV couplings  $\lambda$  and Gravitino masses  $m_{\tilde{G}}$

# Antideuteron Spectrum Near Earth

UiO Department of Physics  
The Faculty of Mathematics and Natural Sciences



- Propagation: NFW DM density profile, 'med' set of diffusion parameters

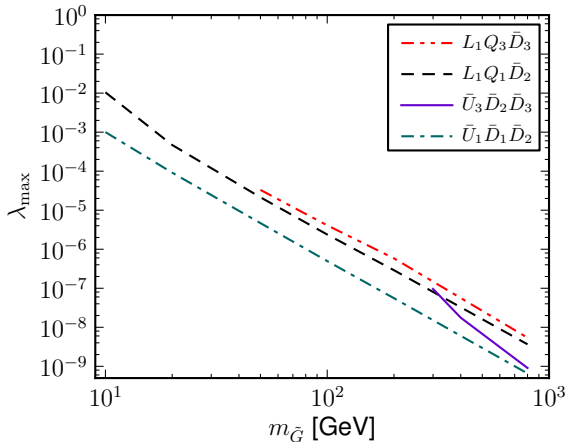


- $m_{\tilde{G}} = 50 \text{ GeV}, \lambda = 10^{-5}$
- Flux increases with increasing mass and RPV coupling
- Can set limits on mass and RPV coupling from experiments

# Limits on RPV couplings



## Prospective upper limits from GAPS



- 95% CL exclusion limits assuming 0 observed events
- Factor 2 – 4 Stronger than existing limits on RPV couplings from PAMELA  $\bar{p}$  data



- Antideuteron spectrum is highly sensitive to hadronization model
- Difference of factor  $\sim 3$  in antideuteron spectrum between Herwig++ and Pythia at most energies, rapidly increasing towards high/low energies
- Tuning necessary for giving a consistent description
- Uncertainty from tuned parameters of factor  $< 2$  after re-tuning
- Antideuterons can be used to set stronger limits on RPV couplings, in particular for  $\bar{U}\bar{D}\bar{D}$ -operators



## Backup Slides



## Tuned Herwig++ hadronization parameters:

- `ClMaxLight`: Involved in specifying mass threshold for fission of clusters of light quarks
- `PSplitLight`: Controls mass distribution of clusters (of light quarks) produced in cluster fission
- `PwtDIquark`: Controls the probability of creating a diquark pair during cluster decay

# Experiments: Number of bins



Experiment	$N_{bins}$
ALEPH	1
CLEO	5
ZEUS	3
CERN ISR	4+4
ALICE	9
ALEPH, $p/\bar{p}$	26

$\chi^2$  from ALEPH proton data weighted down by factor 1/25 to keep it from dominating the parameter determination





- Thermal production of Gravitinos during reheating can give the right relic density

$$\Omega_{\tilde{G}} h^2 \simeq 0.21 \left( \frac{T_R}{10^{10} \text{ GeV}} \right) \left( \frac{100 \text{ GeV}}{m_{\tilde{G}}} \right) \left( \frac{m_{\tilde{g}}(\mu)}{1 \text{ TeV}} \right)^2$$

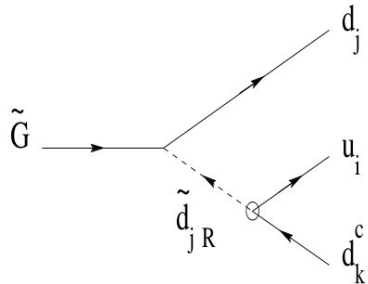
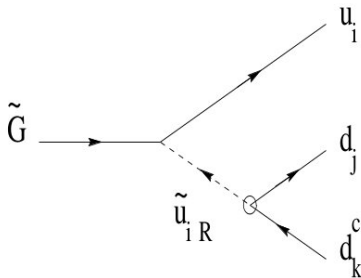
Bolz, Brandenburg, Buchmuller; arXiv:hep-ph/0012052

- The reheating temperature  $T_R$  is weakly constrained, thus so is  $m_{\tilde{G}}$

# Gravitino RPV decays



Tree-level Feynman diagrams for decays through  $\bar{U}_i \bar{D}_j \bar{D}_k$ -operators

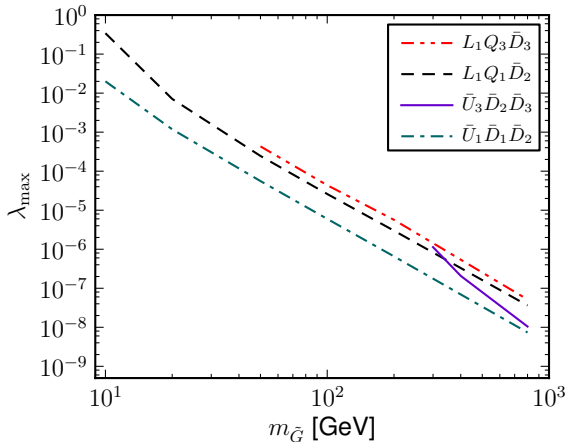


Circle indicates RPV coupling

# Coupling limits: BESS



## Current upper limits from GAPS

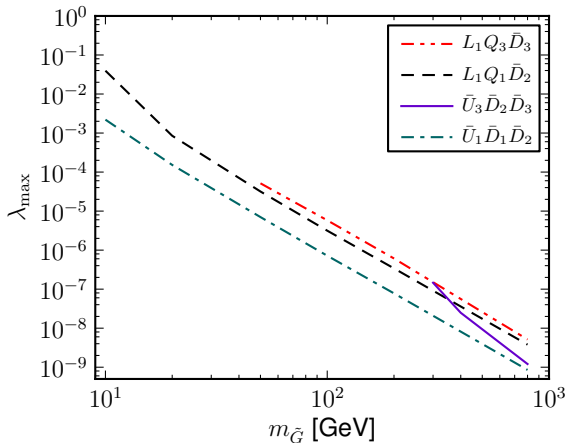


- 95% CL exclusion limits assuming 0 observed events
- Somewhat weaker than existing limits on RPV couplings from PAMELA  $\bar{p}$  data

# Coupling limits: AMS-02



## Prospective upper limits from AMS-02



- 95% CL exclusion limits assuming 0 TOF events and 1 RICH event
- $\lesssim 1$  expected background event in the RICH detector
- $L_i Q_j \bar{D}_k$ : Slightly weaker than  $\bar{p}$  limits at low energies, roughly equal above a few hundred GeV
- $\bar{U}_i \bar{D}_j \bar{D}_k$ : Factor  $\sim 1.5$  Stronger than  $\bar{p}$  limits