

Lattice results on Charmonium

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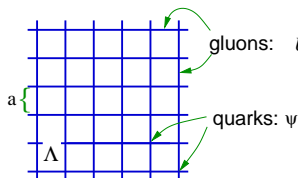
Honolulu,
May 14 2012



- 1 Lattice QCD - Basics and recent developments
 - Recent progress
 - Lattice systematics - a word of caution
 - Excited states, why are they difficult?
- 2 Recent charmonium simulations
 - Fermilab lattice and MILC collaborations
 - Hyperfine-splitting from HPQCD
 - Fermilab method with Wilson-Clover sea quarks
 - Hadron spectrum collaboration
 - BMW and QCDSF collaborations
- 3 Concluding remarks - Where will the lattice go?

- Regularization of QCD on a space-time grid

$$S_{\text{QCD}}[\psi, \bar{\psi}, A] = \sum_{f=1}^{N_f} \int d^4x \bar{\psi}^{(f)}(x) (\not{D} + m^{(f)}) \psi^{(f)}(x) + \frac{1}{2g^2} \text{Tr}[F_{\mu\nu}(x)F_{\mu\nu}(x)]$$



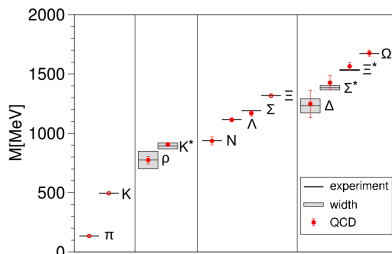
- Lattice $\Lambda = \{ \vec{n} = (n_1, n_2, n_3, n_4) | n_i \in \{0, 1, \dots, L_i - 1\} \}$
- $\vec{x} \rightarrow a\vec{n}$ with lattice spacing a
- $\int d^4x \dots \rightarrow a^4 \sum_{\vec{n} \in \Lambda} \dots$
- $D_\mu \psi(\vec{x})$ covariant lattice derivative

- Enables simulations using well established Monte Carlo methods

Recent developments

- Light mesons and baryons

Ground states have been calculated with full control of systematic uncertainties
Example from Dürer et al. Science 322 (2008)



- Calculations close to

PACS-CS Collaboration Phys. Rev. D79 (2009) 034503

or at the physical pion mass

PACS-CS Phys.Rev. D81 (2010) 074503

BMW Collaboration Phys.Lett. B701 (2011) 265

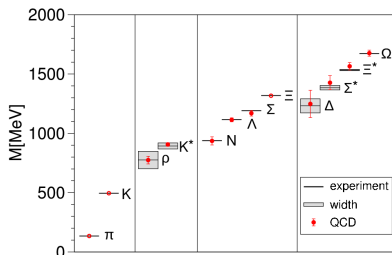
We will soon see more (MILC,...)

- First dynamical 2+1+1 flavor simulations (ETMC, MILC)

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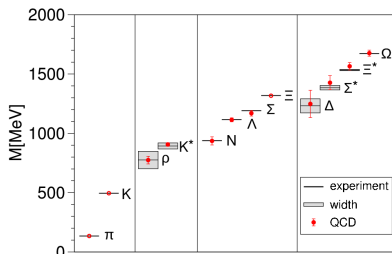
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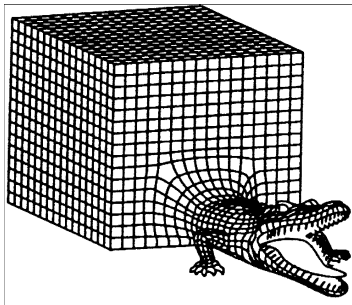
- First dynamical 2+1+1 flavor simulations (ETMC, MILC)

Needed: Multiple extrapolations

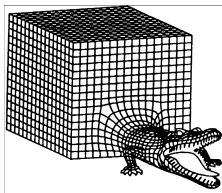
- *Continuum limit:* $a(g, m) \rightarrow 0$
 - Need simulations at multiple different lattice spacings
- *Thermodynamic limit:* $L \rightarrow \infty$ ($L \cdot a = \text{const.}$)
 - Hadron physics in a small box \rightarrow finite volume effects
Typical volume $\approx 2.5\text{fm}$
- Calculation at physical quark masses or extrapolation to the *Chiral limit:* $m \rightarrow m_0$ ($M_\pi \rightarrow M_{\pi, \text{exp}}$)
 - Physical u, d quark masses small \rightarrow Simulation very expensive!
 - Chiral Perturbation Theory (χ PT) \leftrightarrow Lattice QCD

The dangerous animals of lattice QCD

- While progress in the field is fast, studies taking into account all sources of uncertainty are still lacking in many cases.
- Qualitative studies provide/implement crucial insights/improvements and pave the way



Identifying lattice crocodiles - an example

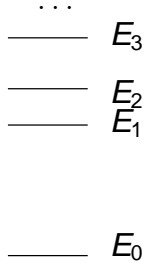


- 1S Hyperfine splitting $M_{HFS} = M_{J/\psi} - M_{\eta_c}$
- traditionally spin-dependent quantities came out too small in simulations
- This has been blamed on many things: quenching/large pion masses, 2 flavor simulations, neglected annihilation diagrams, discretization effects
- I will argue that controlling discretization effects is crucial and that high precision calculations are possible

The problem with excited states

From the analysis of Euclidean correlators:

$$\left\langle \hat{O}_2(t) \hat{O}_1(0) \right\rangle_T \propto \sum_n e^{-tE_n} \langle 0 | \hat{O}_2 | n \rangle \langle n | \hat{O}_1 | 0 \rangle$$

- The whole tower of states contributes
 - Ground state is dominant at large t
 - Excited states appear as sub-leading exponentials
 - Noisy background from limited statistics
- 
- Energy level diagram showing a tower of states E_0, E_1, E_2, E_3 with an ellipsis above E_3 .
- For a single correlator, fit to several exponentials leads to poor results
- Advanced methods needed for excited states!

Method of choice: The variational method

Matrix of correlators projected to fixed momentum (will assume 0)

$$C(t)_{ij} = \sum_n e^{-tE_n} \langle 0 | O_i | n \rangle \langle n | O_j^\dagger | 0 \rangle$$

Solve the generalized eigenvalue problem:

$$C(t) \vec{\psi}^{(k)} = \lambda^{(k)}(t) C(t_0) \vec{\psi}^{(k)}$$
$$\lambda^{(k)}(t) \propto e^{-tE_k} \left(1 + \mathcal{O} \left(e^{-t\Delta E_k} \right) \right)$$

At large time separation: only a single state in each eigenvalue.
Eigenvectors can serve as a fingerprint.

Michael Nucl. Phys. B259, 58 (1985)

Lüscher and Wolff Nucl. Phys. B339, 222 (1990)

Blossier et al. JHEP 04, 094 (2009)

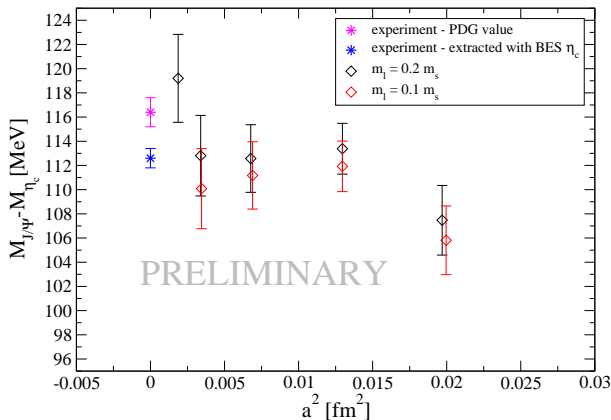
FNAL/MILC program overview

- **Main goal:** Precision study of the low-lying states including higher spin states with a full error budget
- **Analysis campaigns:**

Study	sea	Charm	Ops	States	Comment
2009 FNAL/MILC	2 + 1	FNAL	2	low S, P	$a \geq 0.09$ fm
2012 FNAL/MILC	2 + 1	FNAL	many	many	$a \geq 0.045$ fm
Current FNAL/MILC	2 + 1 + 1	HISQ	2	low S, P	$a \geq 0.06$ fm
Future FNAL/MILC	2 + 1 + 1	HISQ	many	many	$a \geq 0.06$ fm

- I will present selected *preliminary* results for 2+1 flavor asqtad
- Data for many other states (including higher spin) is available but not yet fully analyzed
- This is a progress report

Selected results: Charmonium hyperfine splitting



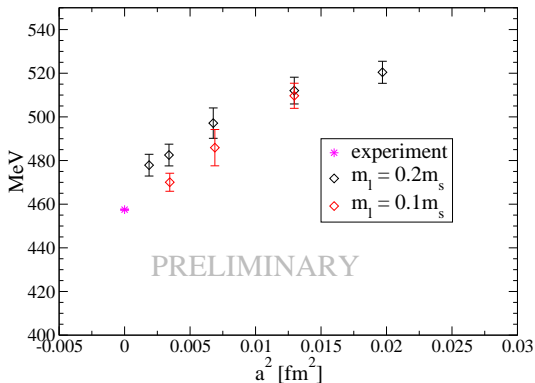
- Use of an improved heavy-quark action → mild discretization effects
- Quark annihilation contributions not yet included

Selected results: $\overline{1S} - \overline{1P}$ splitting

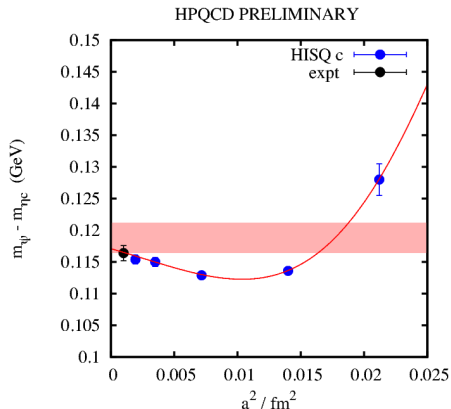
- Defining the mass centroids

$$M_{\overline{1S}} = \frac{1}{4}(M_{\eta_c} + 3M_{J/\psi})$$

$$M_{\overline{1P}} = \frac{1}{9}(M_{\chi_{c0}} + 3M_{\chi_{c1}} + 5M_{\chi_{c2}})$$



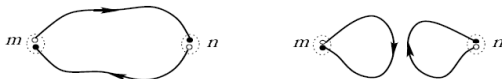
Hyperfine-splitting from HPQCD



from E. Follana PoS (Lattice 2010) 305

- The red band indicates the preliminary result: 118.8 ± 2.4 MeV

Disconnected contributions?



- For HPQCD hyperfine-splitting, contributions from disconnected diagrams are estimated from perturbation theory which leads to a shift enlarging the hyperfine splitting.
- MILC estimated the effects by direct calculations and find that disconnected diagrams **reduce** the hyperfine splitting by 1-4MeV
Levkova and DeTar, Phys.Rev. D83 (2011) 074504
- Novel methods should lead to a more precise determination (see later) Liu et al. arXiv:1204.5425

Two exploratory studies

- $N_f = 2 + 1$ Wilson-Clover sea and Fermilab charm quarks at almost physical pion masses

$N_L^3 \times N_T$	$a[\text{fm}]$	$L[\text{fm}]$	#configs	$m_\pi[\text{MeV}]$
$32^3 \times 32$	0.0907(13)	2.9	≥ 198	$156 \leq m_\pi \leq 702$

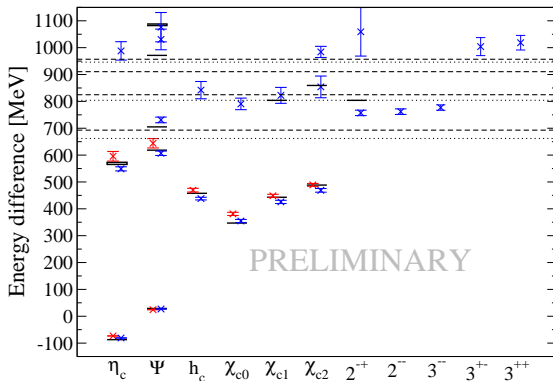
DM and R. M. Woloshyn Phys.Rev. D84 (2011) 054505

- $N_f = 2$ smeared Wilson-Clover sea and Fermilab charm quarks

$N_L^3 \times N_T$	$a[\text{fm}]$	$L[\text{fm}]$	#configs	$m_\pi[\text{MeV}]$
$16^3 \times 32$	0.1239(13)	1.98	280	266(3)(3)

DM, S. Prelovsek and R. M. Woloshyn (to be published)

Results from $N_f = 2$ smeared Wilson-Clover sea and Fermilab charm



DM and R. M. Woloshyn Phys.Rev. D84 (2011) 054505

DM, S. Prelovsek and R. M. Woloshyn (to be published)

- High statistical precision has been achieved through the *distillation* technique Peardon et al. Phys.Rev. D80 (2009) 054506

Good statistical precision for many quantities!

Mass difference	This study [MeV]	Experiment [MeV]
$\overline{1P} - \overline{1S}$	$441.7 \pm 4.0 \pm 4.6$	457.5 ± 0.3
$\overline{2S} - \overline{1S}$	$592.3 \pm 4.9 \pm 6.2$	606.1 ± 1.0
1S hyperfine	$107.9 \pm 0.3 \pm 1.1 \pm_{0}^{2.2}$	116.6 ± 1.2
1P spin-orbit	$39.7 \pm 2.1 \pm 0.4$	46.6 ± 0.1
1P tensor	$11.02 \pm 0.87 \pm 0.12$	16.25 ± 0.07
1P hyperfine	3.7 ± 2.7	-0.10 ± 0.22
2S hyperfine	57.9 ± 2.0	49 ± 4
2P spin-orbit	$24.6 \pm 15.7 \pm 0.3$	-
2P tensor	2.2 ± 4.3	-
$\overline{2P} - \overline{1S}$	$836.4 \pm 30.5 \pm 8.8$	-
$2M_{D_s} - M_{c\bar{c}}$	1065.6	1084.8 ± 0.8

Excited charmonium states from anisotropic Wilson-Clover lattices

- The *Hadron Spectrum Collaboration* has recently studied ground and excited charmonium states

Liu et al. arXiv:1204.5425

- They use anisotropic Wilson-Clover lattices with volumes $16^3 \times 128$ and $24^3 \times 128$ at a single lattice spacing ($a_s \approx 0.12 \text{ fm}$ and $\xi = 3.5$)
- They simulate at a single pion mass $m_\pi = 396 \text{ MeV}$
- A large basis containing interpolating fields with up to three covariant derivatives is used

Excited charmonium states from anisotropic Wilson-Clover lattices

Notable features:

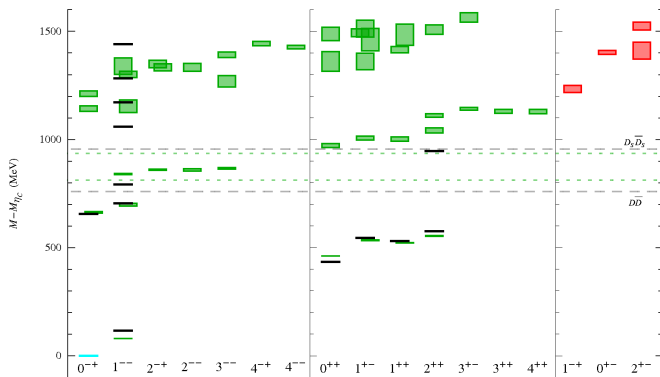
- Use of the *distillation* technique allows calculations with small statistical uncertainty
- Advanced spin identification techniques allows unambiguous spin assignment.

Peardon et al. Phys.Rev. D80 (2009) 054506

Dudek et al. Phys.Rev. D82 (2010) 034508

This allows to identify Spin 4 states and states with **exotic** quantum numbers.

Excited charmonium states from anisotropic Wilson-Clover lattices



from Liu et al. arXiv:1204.5425

What about systematics?

- Hyperfine splitting:

$$M_{HFS} = 80 \pm 1 \text{ MeV}$$

Illustrates the typical size of discretization effects.

- This can be improved by the use of an improved heavy quark action or by continuum extrapolation
- Higher states will be influenced by unphysical thresholds. These can be included systematically within the distillation technique!

Charmonium using $N_f = 2 + 1$ BMW and CQDSF configurations

- **BMW configurations:** Smearred $N_f = 2 + 1$ Wilson Clover

$$0.054\text{fm} \leq a \leq 0.092\text{fm}$$

$$120\text{MeV} \leq M_\pi \leq 520\text{MeV}$$

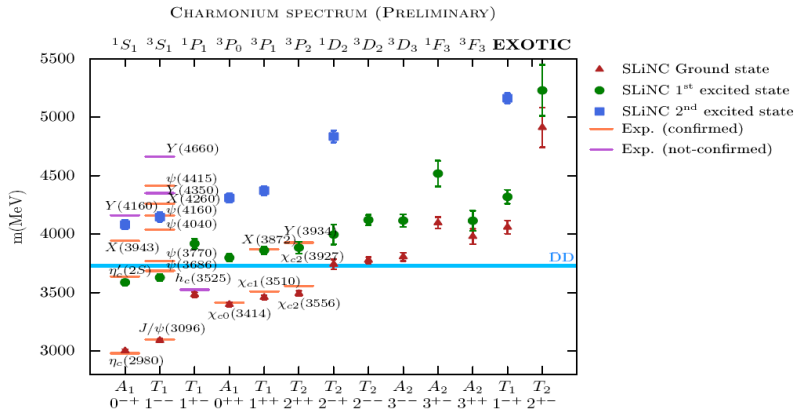
$$L^3 \times T = 32^3 \times 64, 64^3 \times 144$$

- **QCDSF configurations:** $N_f = 2 + 1$ SLiNC fermions
So far preliminary results from the following ensembles:

$N_L^3 \times N_T$	$a[\text{fm}]$	$L[\text{fm}]$	$m_\pi[\text{MeV}]$
$24^3 \times 48/32^3 \times 64$	0.0795(3)	1.9/ 2.5	442
$24^3 \times 48/32^3 \times 64$	0.0795(3)	1.9/ 2.5	348

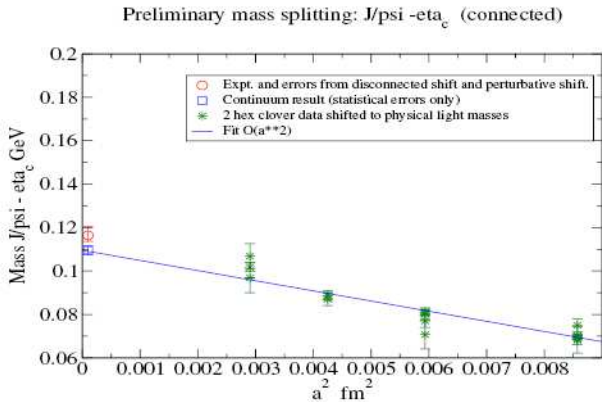
- In both cases the same action is used for light and charm quarks.

BMW/QCDSF - Preliminary results for the charmonium spectrum



from Bali et al. arXiv:1108.6147

BMW/QCDSF - Preliminary charmonium hyperfine splitting



from Bali et al. arXiv:1108.6147

Concluding remarks - Where will the lattice go?

- The era of precision calculations using Lattice QCD has barely begun and there already is remarkable success
- Calculations at or around physical pion masses will become commonplace in the near future
- There are multiple collaborations generating gauge configurations with 2+1 or 2+1+1 dynamical flavors suitable for charmonium physics
- For states below inelastic thresholds, the extraction of resonance properties from the lattice is within reach
- This necessitates the inclusion of DD, DD*, etc. states. For a first step see
from Bali, Collins, Ehmman Phys.Rev. D84 (2011) 094506
- There is much work ahead with regard to hadron excitations

Thank you!