



Effects of the $D\bar{D}\pi$ threshold on the $X(3872)$

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Discovery of $X(3872)$ in B decay

Belle collaboration, August 2003

hep-ex/0309032

$$B^\pm \longrightarrow X + K^\pm, \quad X \longrightarrow J/\psi \pi^+ \pi^-$$

confirmed by Babar

hep-ex/0406022

Observation of $X(3872)$ in $p\bar{p}$ collisions

CDF collaboration, September 2003

hep-ex/0312021

$$p\bar{p} \longrightarrow X + \text{anything}, \quad X \longrightarrow J/\psi \pi^+ \pi^-$$

confirmed by D0

hep-ex/0405004

Mass and Width of the $X(3872)$

Mass

Belle, CDF, Babar, D0

$$M_X = 3871.9 \pm 0.5 \text{ MeV}$$

Width

Belle, hep-ex/0309032

$$\Gamma_X < 2.3 \text{ MeV} \quad (90\% \text{ C.L.})$$

Mass is extremely close to the $D^{*0}\bar{D}^0$ threshold

$$M_X - (M_{D^{*0}} + M_{D^0}) = -0.31 \pm 0.64 \text{ MeV}$$

using recent precision determination of D^0 mass by CLEO

Quantum Numbers of the $X(3872)$

$$X \longrightarrow J/\psi \gamma \\ \implies C = +$$

Belle, hep-ex/0505037

$$X \longrightarrow J/\psi \pi^+ \pi^-$$

Belle, hep-ex/0505038

angular correlations \implies not $J^{PC} = 0^{++}, 0^{-+}$

$\pi\pi$ mass distribution \implies not $1^{-+}, 2^{-+}$

1^{++} “strongly favored”, but 2^{++} “not ruled out”

$$X \longrightarrow D^0 \bar{D}^0 \pi^0$$

Belle, preliminary

8 MeV above threshold $\implies 2^{++}$ ruled out

$$J^{PC} = 1^{++}$$

Decays of the $X(3872)$

Measured branching ratios

$$\frac{\text{Br}[X \rightarrow J/\psi \pi^+ \pi^-]}{\text{Br}[X \rightarrow D^0 \bar{D}^0 \pi^0]} \approx 0.059 \pm 0.025 \quad ??$$

Belle, preliminary ??

$$\frac{\text{Br}[X \rightarrow J/\psi \pi^+ \pi^- \pi^0]}{\text{Br}[X \rightarrow J/\psi \pi^+ \pi^-]} = 1.0 \pm 0.4 \pm 0.3$$

Belle, hep-ph/0505038

$$\frac{\text{Br}[X \rightarrow J/\psi \gamma]}{\text{Br}[X \rightarrow J/\psi \pi^+ \pi^-]} = 0.14 \pm 0.05$$

Belle, hep-ex/0505037

What is the $X(3872)$?

Belle: D-wave charmonium? $D^*\bar{D}$ molecule?

- charmonium

Close and Paige (9-03); Pakvasa and Suzuki (9-03); Barnes and Godfrey (11-03); Eichten, Lane, and Quigg (1-04); Olsen (7-04); Meng, Gao, and Chao (6-05)
motivated by J/ψ among decay products

- $D^*\bar{D}$ molecule

Tornqvist (8-03, 2-04); Close and Paige (9-03); Pakvasa and Suzuki (9-03); Voloshin (9-03, 8-04, 9-05, 5-06); Wong (11-03); Braaten and Kusunoki (11-03, 2-04, 12-04, 6-05, 7-05, 9-06); Swanson (11-03, 6-04); Braaten, Kusunoki, and Nussinov (4-04); Kalashnikova (6-05); AIFiky, Gabbiani, and Petrov (6-05); ElHady (3-06)
motivated by proximity to $D^*\bar{D}$ threshold

- threshold cusp from interactions with $D^*\bar{D}$

Bugg (10-04)

motivated by proximity to $D^*\bar{D}$ threshold

What is the $X(3872)$? (cont.)

- tetraquark: $(cq)_3^* (\bar{c}\bar{q})_3$

Vijande, Fernandez, and Valcarce (7-04); Maiani, Piccinini, Polosa, and Riquer (12-04); Ishida, Ishida, and Maeda (9-05, 10-06); Ebert, Faustov, and Galkin (12-05); Karliner and Lipkin (1-06); Chiu and Hsieh (3-06)

- tetraquark: $(c\bar{c})_8 (q\bar{q})_8$

Hogassen, Richard, and Sorba (11-05)

- charmonium hybrid: $c\bar{c}g$

Close and Paige (9-03); Li (10-04)

- glueball: gg

Seth (11-04)

See also Ko (5-04); Rosner (8-04)

What is the $X(3872)$?

Experimental inputs

- $J^{PC} = 1^{++}$
 \implies S-wave coupling to $D^{*0}\bar{D}^0$, $D^0\bar{D}^{*0}$
- Mass is extremely close to the $D^{*0}\bar{D}^0$ threshold

$$M_X - (M_{D^{*0}} + M_{D^0}) < -0.31 \pm 0.64 \text{ MeV}$$

Assumption: $M_X < M_{D^{*0}} + M_{D^0}$

Conclusion

$X(3872)$ is a weakly-bound charm meson molecule

$$X = \frac{1}{\sqrt{2}} (D^{*0}\bar{D}^0 + D^0\bar{D}^{*0})$$

with universal properties that are insensitive to details of QCD

What is the X ? (cont.)

Quantum Mechanics

2-body system with **short-range interactions**

and **S-wave bound state** sufficiently close to **threshold**
has **universal** properties

that depend only on the **large scattering length** a

“**Universality** of Few-Body Systems with **Large Scattering Length**”

Braaten and Hammer, Physics Reports (cond-mat/0410417)

- **^4He atoms**: $a \approx +100 \text{ \AA}$
- **spin-polarized ^3H atoms**: $a_t \approx -82 \text{ \AA}$
- **alkali atoms** near Feshbach resonance: $-\infty < a(B) < \pm\infty$
- **nucleons**: $a_s \approx -23.8 \text{ fm}$, $a_t \approx +5.4 \text{ fm}$
- **charm mesons** $D^{*0}\bar{D}^0$ and $D^0\bar{D}^{*0}$!!

What is the X ? (cont.)

Long-distance interaction between low-energy D^{*0} and \bar{D}^0 are dominated by pion exchange

Natural scales from pion exchange

length: $1/m_\pi = 1.4$ fm

momentum: $m_\pi = 140$ MeV

energy: $m_\pi^2/(2M_{D^*D}) = 10$ MeV

\implies large scattering length: $|a| \gg 1.4$ fm

\implies small binding energy: $E_X \ll 10$ MeV

Binding energy of $X(3872)$

Belle, Babar, CDF, D0, CLEO

$$E_X = 0.31 \pm 0.64 \text{ MeV}$$

What is the X ? (cont.)

Since X decays,

there must be inelastic scattering channels for $D^{*0}\bar{D}^0$, $D^0\bar{D}^{*0}$

\implies Scattering length a is complex

$$\frac{1}{a} = \gamma_{\text{re}} + i\gamma_{\text{im}}$$

Real and imaginary parts of $1/a$ determine

binding energy : $E_X = (\gamma_{\text{re}}^2 - \gamma_{\text{im}}^2)/(2M_{D^*D})$

decay width : $\Gamma_X = 2\gamma_{\text{re}}\gamma_{\text{im}}/M_{D^*D}$

Constraints on E_X and Γ_X

$$0 < E_X < 1.1 \text{ MeV}$$

$$0.07 \text{ MeV} < \Gamma_X < 2.3 \text{ MeV}$$

What is the X ? (cont.)

Universal results

- elastic $D^{*0}\bar{D}^0$ cross section at low energy E

$$\sigma(E) = \frac{\pi}{\gamma_{re}^2 + (\gamma_{im} + \sqrt{2M_{D^*D}E})^2}$$

huge cross section at threshold: $\sigma(E=0) > 60 \text{ fm}^2$!!

- shallow S-wave bound state: $X(3872)$

wavefunction: $\psi(r) = \exp(-\gamma_{re} r)/r$

$$\implies \langle r \rangle = 1/(2\gamma_{re})$$

huge mean separation: $\langle r \rangle_X > 2 \text{ fm}$!!

What is the X ? (cont.)

$X(3872)$ is a weakly-bound charm meson molecule

$$X = \frac{1}{\sqrt{2}} (D^{*0}\bar{D}^0 + D^0\bar{D}^{*0})$$

with universal properties that depend on E_X , Γ_X
but are insensitive to other details of QCD

insensitive to mechanism for binding

- $\chi_{c1}(2P)$ near $D^{*0}\bar{D}^0$ threshold?
- 1^{++} tetraquark near $D^{*0}\bar{D}^0$ threshold?
- depth of pion-exchange potential between charm mesons?
- coupling of charmonium to charm meson pairs?

If couplings to $D^{*0}\bar{D}^0$ and $D^0\bar{D}^{*0}$ are taken into account,
 $X(3872)$ will be a weakly-bound charm meson molecule
regardless of the mechanism for binding

Decays of the $X(3872)$

Long-distance decays

$$\begin{aligned} X &\longrightarrow D^0 \bar{D}^0 \pi^0 \\ &\longrightarrow D^0 \bar{D}^0 \gamma \end{aligned}$$

dominated by $D^{*0} \bar{D}^0$, $D^0 \bar{D}^{*0}$ components of X
proceed through decays of constituent D^{*0} or \bar{D}^{*0}
universal, determined by E_X , Γ_X only

Short-distance decays

$$\begin{aligned} X &\longrightarrow \text{charmonium} + \text{light hadrons} \\ &\longrightarrow D^+ D^- \gamma \\ &\longrightarrow \text{light hadrons} \end{aligned}$$

dominated by short-distance (“core”) components of X
factorization formula with universal long-distance factor

Long-distance Decays of $X(3872)$

$$X(3872) \longrightarrow D^0 \bar{D}^0 \pi^0, D^0 \bar{D}^0 \gamma$$

Voloshin (9-03); El-Hady (3-06)

dominated by $D^{*0} \bar{D}^0, D^0 \bar{D}^{*0}$ components

proceeds through decay of constituent: $D^{*0} \longrightarrow D^0 \pi^0, D^0 \gamma$
 $\bar{D}^{*0} \longrightarrow \bar{D}^0 \pi^0, \bar{D}^0 \gamma$

- absolute predictions of partial widths

$$\begin{aligned} \Gamma[X \rightarrow D^0 \bar{D}^0 \pi^0] &\approx \Gamma[D^{*0} \rightarrow D^0 \pi^0] \\ \Gamma[X \rightarrow D^0 \bar{D}^0 \gamma] &\approx \Gamma[D^{*0} \rightarrow D^0 \gamma] \end{aligned}$$

- lower bound on total width

$$\Gamma[X] > \Gamma[D^{*0}] \approx 70 \pm 15 \text{ keV}$$

- momentum distributions for $D^0 \bar{D}^0 \pi^0, D^0 \bar{D}^0 \gamma$
determined primarily by E_X, Γ_X

Factorization

Braaten and Kusunoki (6-05), Braaten and Lu (6-06)

separate **short** distances $\lesssim 1/m_\pi, 1/\Lambda_{\text{QCD}}, 1/m_c$
from **long** distances $\gtrsim |a|$

Factorization formulas for amplitudes

$$\mathcal{A} = \mathcal{A}_{\text{short}} \times \mathcal{A}_{\text{long}}$$

$\mathcal{A}_{\text{short}}$ has smooth limit as $a \rightarrow \infty$

$\mathcal{A}_{\text{long}}$ is determined by a only

Amplitudes for $X(3872)$

$D^{*0}\bar{D}^0$ near threshold

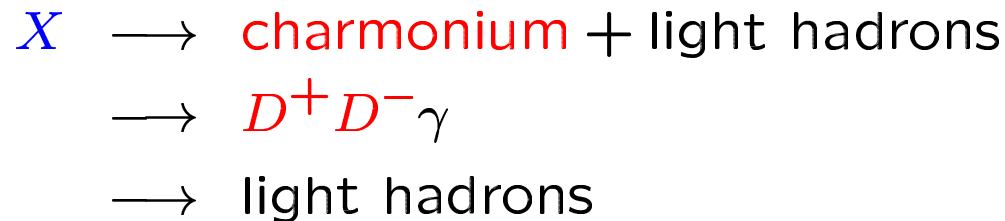
$D^0\bar{D}^{*0}$ near threshold

same **short-distance** factor $\mathcal{A}_{\text{short}}$

long-distance factors $\mathcal{A}_{\text{long}}$ determined by E_X and Γ_X

Short-distance Decays of $X(3872)$

dominated by **short-distance** components of X



Factorization formula

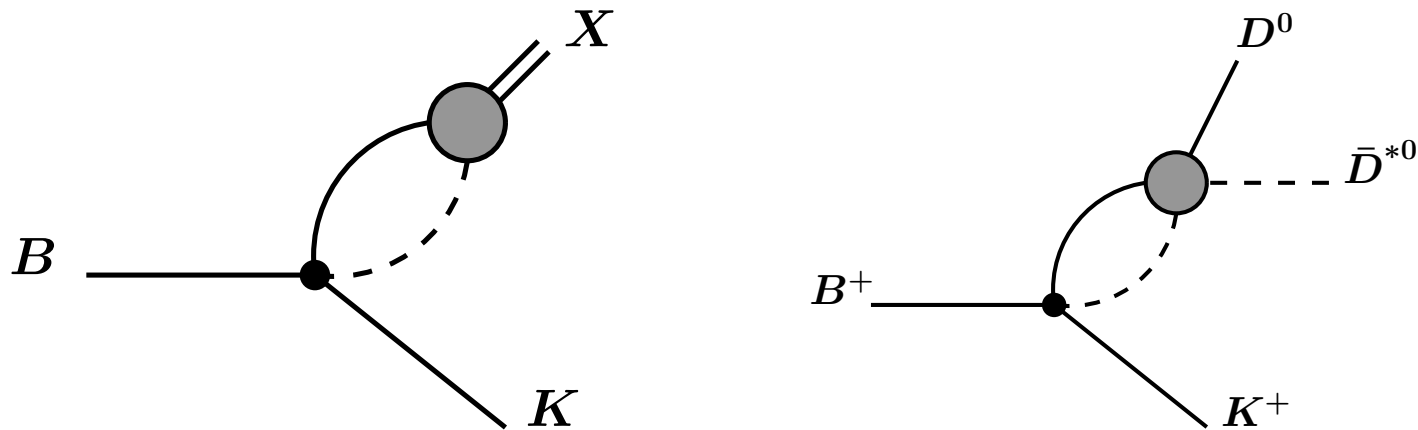
Braaten and Kusunoki (7-05); Braaten and Lu (6-06)

$$\Gamma[X \rightarrow H] = \Gamma_{\text{short}}[H] \times \left(E_X + \Gamma_X^2 / 16 E_X \right)^{1/2}$$

- **long-distance** factor cancels in ratio of decay rates
 \implies ratios are insensitive to E_X, Γ_X

Production of $X(3872)$, $D^{*0}\bar{D}^0$

Braaten, Kusunoki, and Nussinov (4-04), Braaten and Kusunoki (12-04)



Factorization formulas

$$\Gamma[B \rightarrow X + K] = \Gamma_{\text{short}} \times \frac{\pi}{M_{D^*D}} [\gamma_{\text{re}}^2 + \gamma_{\text{im}}^2]^{1/2}$$

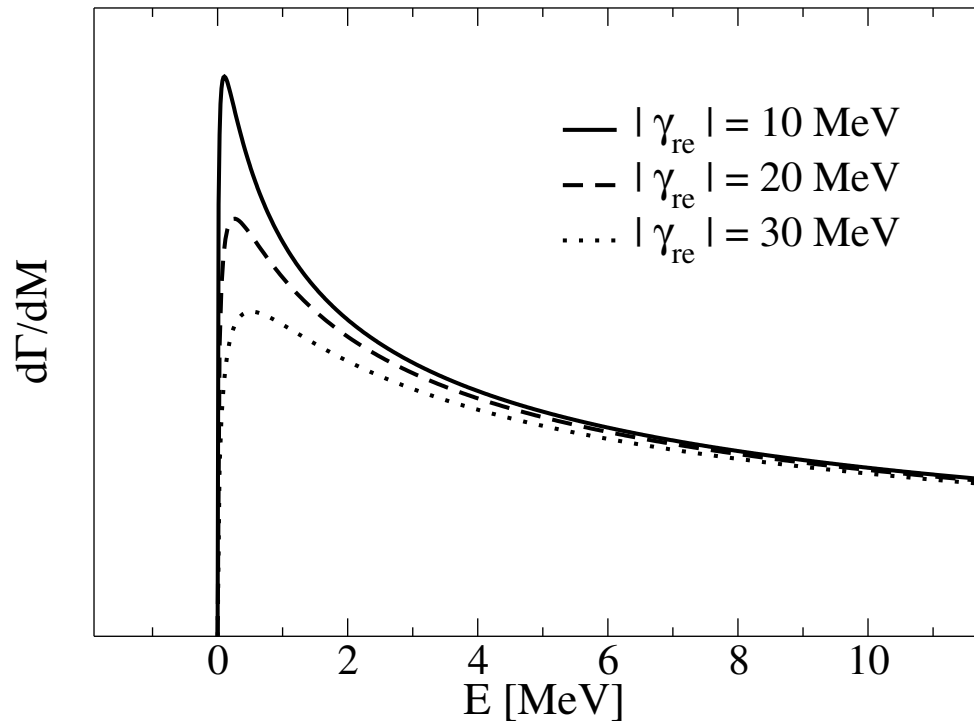
$$\frac{d\Gamma}{dM} [B \rightarrow D^{*0}\bar{D}^0 + K] = \Gamma_{\text{short}} \times \frac{\sqrt{2M_{D^*D}E}}{\gamma_{\text{re}}^2 + (\gamma_{\text{im}} + \sqrt{2M_{D^*D}E})^2}$$

$$M = (M_{D^0} + M_{D^{*0}}) + E$$

$B \rightarrow K + X, D^* \bar{D}$ (cont.)

Decay $B^+ \rightarrow K^+ + D^{*0} \bar{D}^0$

$$\frac{d\Gamma}{dM} [B \rightarrow K + D^{*0} \bar{D}^0] = \Gamma_{\text{short}} \times \frac{\sqrt{2M_{D^*D}E}}{\gamma_{\text{re}}^2 + (\gamma_{\text{im}} + \sqrt{2M_{D^*D}E})^2}$$



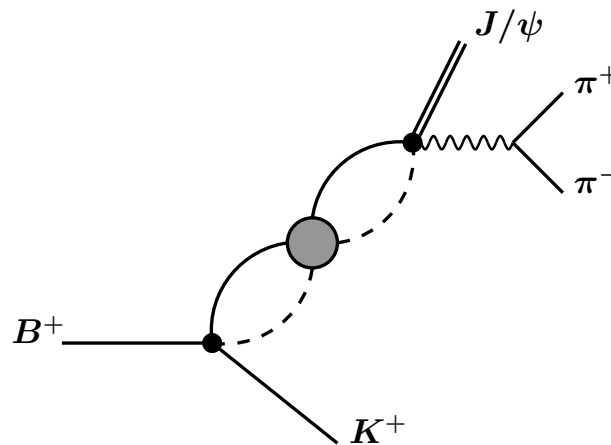
$D^{*0} \bar{D}^0$ invariant mass distribution peaks at $E = E_X + \Gamma_X^2 / (16E_X)$

Line Shape of X

invariant mass distribution of decay products of X

short-distance decay channel: $X \rightarrow H$

$$H = J/\psi \pi^+ \pi^-, J/\psi \pi^+ \pi^- \pi^0, J/\psi \gamma, \dots$$



Factorization formula

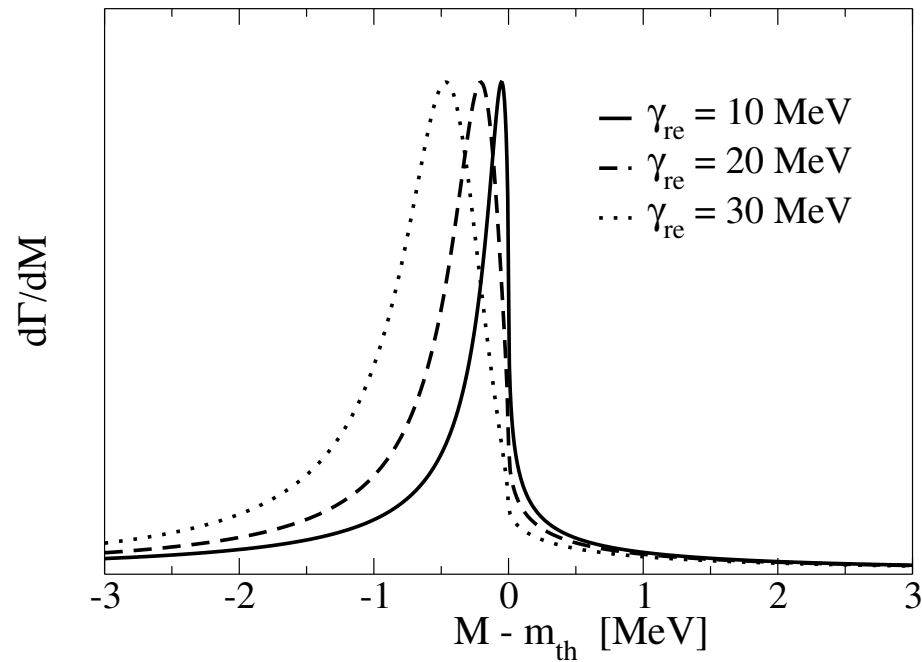
$$\frac{d\Gamma}{dM} [B \rightarrow K + H] = \Gamma_{\text{short}} \times \left| \frac{1}{\gamma_{\text{re}} + i\gamma_{\text{im}} - \sqrt{-2M_{D^*D}E}} \right|^2$$

$$M = (M_{D^*0} + M_{D0}) + E$$

Line Shape of X (cont.)

invariant mass of H : $M = (M_{D^*0} + M_{D0}) + E$

$$\frac{d\Gamma}{dM} [B \rightarrow K + H] = \Gamma_{\text{short}} \times \left| \frac{1}{\gamma_{\text{re}} + i\gamma_{\text{im}} - \sqrt{-2M_{D^*D}E}} \right|^2$$



non-Breit-Wigner resonance near $E = -E_X$

Nearby Meson Thresholds

2-meson and 3-meson thresholds

that can be reached by $D^* \longleftrightarrow D\pi$

		order [†] in g
$D^{*+}D^{-}, D^{+}D^{*-}$	+8.1 MeV	g^2
$D^{+}D^{-}\pi^0$	+2.5 MeV	g^3
$D^{+}\bar{D}^0\pi^{-}, D^0\bar{D}^{-}\pi^{+}$	+2.3 MeV	g
$D^{*0}\bar{D}^0, D^0\bar{D}^{*0}$	0	
$D^0\bar{D}^0\pi^0$	-7.1 MeV	g

[†] g = coupling constant for $D^* \leftrightarrow D\pi$

Nearby Meson Thresholds (cont.)

What is the effect of nearby $D^0\bar{D}^0\pi^0$ threshold?

Voloshin

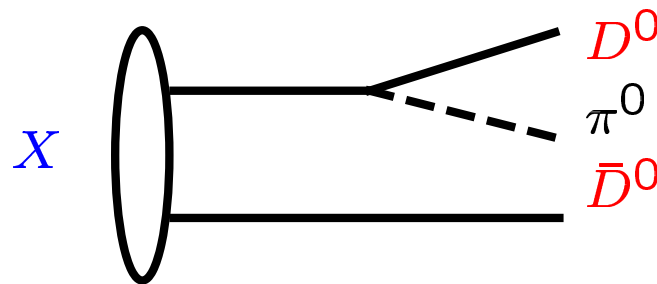
Phys. Lett. B **579**, 316 (2004)

“Interference and binding effects in decays of possible molecular component of $X(3872)$ ”

X decays into $D^0\bar{D}^0\pi^0$ through decay of constituent:

$$D^{*0} \longrightarrow D^0\pi^0 \quad \text{or} \quad \bar{D}^{*0} \longrightarrow \bar{D}^0\pi^0$$

Calculate momentum distribution of $D^0\bar{D}^0\pi^0$
using wavefunction for $D^{*0}\bar{D}^0 / D^0\bar{D}^{*0}$
and treating decay as 1st order perturbation



Problem: effects of $D^0\bar{D}^0\pi^0$ become increasingly nonperturbative as $M_{D^*} - M_D - m_\pi \longrightarrow 0$

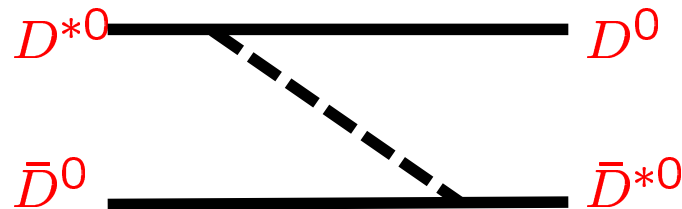
Nearby Meson Thresholds (cont.)

What is the effect of nearby $D^0\bar{D}^0\pi^0$ threshold?

Suzuki

Phys. Rev. D **72**, 114013 (2005)

“The $X(3872)$ boson: Molecule or charmonium”



exchanged π^0 can be on-shell, has momentum of order

$$\mu \equiv \left(2m_{\pi^0}(M_{D^{*0}} - M_{D^0} - m_{\pi^0})\right)^{1/2} \approx 44 \text{ MeV}$$

Static potential:

$$V(r) = -\text{constant} \times \delta^3(\mathbf{r}) + \text{terms suppressed by } \mu^2$$

“one-pion exchange produces practically no force
between D^* and \bar{D} , ... incapable of binding D^* and \bar{D} ”

Nearby Meson Thresholds

What is the effect of nearby $D\bar{D}\pi$ thresholds?

Meson Potential Models

2-body channels only, no 3-body channels

Tornquist (8-03, 2-04)

one-pion exchange

coupled channels: $D^*\bar{D}$, S-wave and P-wave

- predicted possible existence of $X(3872)$

Swanson (11-03, 6-04)

one-pion exchange + quark-exchange

5 coupled channels: $D^{*0}\bar{D}^0 + D^0\bar{D}^{*0}$ S-wave and P-wave

$D^{*+}D^- + D^+D^{*-}$ S-wave

$J/\psi \rho$ S-wave; $J/\psi \omega$ S-wave

- correctly predicted $\text{Br}[X \rightarrow J/\psi \pi^+\pi^-\pi^0] \sim \text{Br}[X \rightarrow J/\psi \pi^+\pi^-]$
- incorrectly predicted $\text{Br}[X \rightarrow J/\psi \pi^+\pi^-\pi^0] \gg \text{Br}[X \rightarrow D^0\bar{D}^0 \pi^0]$
- predicts weakly-bound $0^{++} D^*\bar{D}^*$ molecule

Nearby Meson Thresholds

Question: Can effects of nearby $D\bar{D}\pi$ thresholds be treated as perturbations?

Unnecessary complications

- spin-1 constituent: D^{*0} or \bar{D}^{*0}
- $C = + \implies X = (D^{*0}\bar{D}^0 + D^0\bar{D}^{*0})/\sqrt{2}$
- nearby 2-particle thresholds: $D^{*+}D^-$, D^+D^{*-}
- derivative coupling of pions: $(D^{*\mu})^\dagger \vec{\tau} \bar{D} \cdot \partial_\mu \vec{\pi}$

Question can be addressed in simpler model.

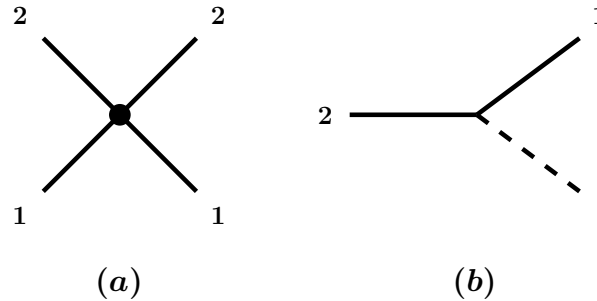
Scalar Meson Model

Scalar mesons: D_1, D_2, ϕ

Masses: $M_2 - (M_1 + m_\phi) \ll m_\phi \ll M_1 < M_2$

Channels near D_1D_2 threshold: $D_1D_2, D_1D_1\phi$

Interactions

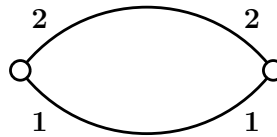


- D_1D_2 coupling constant: λ_0
 must be treated **nonperturbatively**
 requires **renormalization**
 real, imaginary parts determined by E_X, Γ_X
- $D_2 \rightarrow D_1\phi$ coupling constant: g
 determined by decay rate Γ_2 for $D_2 \rightarrow D_1\phi$
perturbative ??

Scalar Meson Model (cont.)

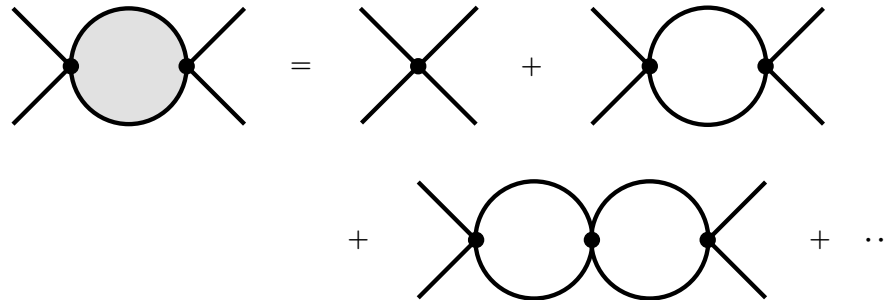
0th Order in g

amplitude for propagation of $D_1 D_2$ between contact interactions



$$L_0(E) = \infty + \kappa, \quad \kappa = \sqrt{-2M_{12}(E - M_1 - M_2 + i\epsilon)}$$

amplitude for $D_1 D_2 \rightarrow D_1 D_2$: sum to all orders in λ_0



$$\mathcal{A}_0(E) = \frac{1}{1/\lambda_0 - [\infty + \kappa]} = \frac{1}{1/\lambda - \kappa}$$

reproduces **universal** properties

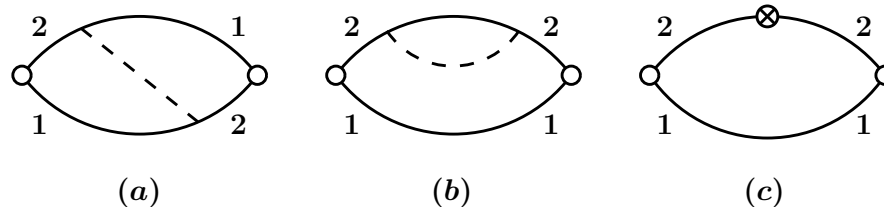
for particles with **large scattering length**

Scalar Meson Model (cont.)

2nd Order in g

corrections to amplitude

for propagation of $D_1 D_2$ between contact interactions



$$L_2(E) = \infty + g^2 F(\kappa), \quad \kappa = \sqrt{-2M_{12}(E - M_1 - M_2 + i\epsilon)}$$

renormalization requires summing geometric series to all orders

$$A_2(E) = \frac{1}{1/\lambda_0 - [\infty + \kappa] - [\infty + g^2 F(\kappa)]} = \frac{1}{1/\lambda - \kappa - g^2 F(\kappa)}$$

Infrared Problem:

$F(\kappa)$ diverges as $1/\kappa$ at the $D_1 D_2$ threshold ($\kappa \rightarrow 0$)

\implies nonperturbative as $\kappa \rightarrow 0$

Scalar Meson Model (cont.)

Infrared Problem at 2nd order in g

$$\mathcal{A}_2(E) = \frac{1}{1/\lambda - \kappa - g^2 F(\kappa)}$$

divergence in $F(\kappa)$ as $\kappa \rightarrow 0$

Resolution: divergence is proportional to width Γ_2 of D_2

$$g^2 F(\kappa) \longrightarrow -i \frac{M_{12} \Gamma_2}{2\kappa}$$

can be resummed into leading-order amplitude

by shifting mass: $M_2 \longrightarrow M_2 - i\Gamma_2/2$

$$\mathcal{A}_2(E) = \left(1/\lambda - \sqrt{-2M_{12}(E - M_1 - M_2 + i\Gamma_2/2)} - \left[g^2 F(\kappa) + iM_{12}\Gamma_2/(2\kappa) \right] \right)^{-1}$$

Resummed amplitude has same expansion to order g^2
smooth limit as $\kappa \rightarrow 0$

Summary

$X(3872)$ is weakly-bound charm meson molecule

$$X(3872) = \frac{1}{\sqrt{2}} (D^{*0}\bar{D}^0 + D^0\bar{D}^{*0})$$

has **universal** properties that depend on E_X , Γ_X
but are insensitive to other **details of QCD**

Unique feature of $X(3872)$ is proximity to **3-meson thresholds**:

$$D^0\bar{D}^0\pi^0, \quad D^+\bar{D}^0\pi^-, \quad D^0D^-\pi^+$$

Effect of **3-meson threshold** can be calculated
using (resummation of) **perturbation theory**

Applications

- momentum distribution for $X \longrightarrow D^0\bar{D}^0\pi^0$
- line shape of X in $J/\psi \pi^+\pi^-$
- ...

The Truth is Out There