

Nongeometric String Backgrounds

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Outline

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II. D-branes on “nongeometric”
backgrounds

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hep-th/0602025

III. Nongeometric backgrounds
and spacetime supersymmetry

A. Lawrence, R. Minasian, T. Sander, M. Schulz,
and B. Wecht, in progress

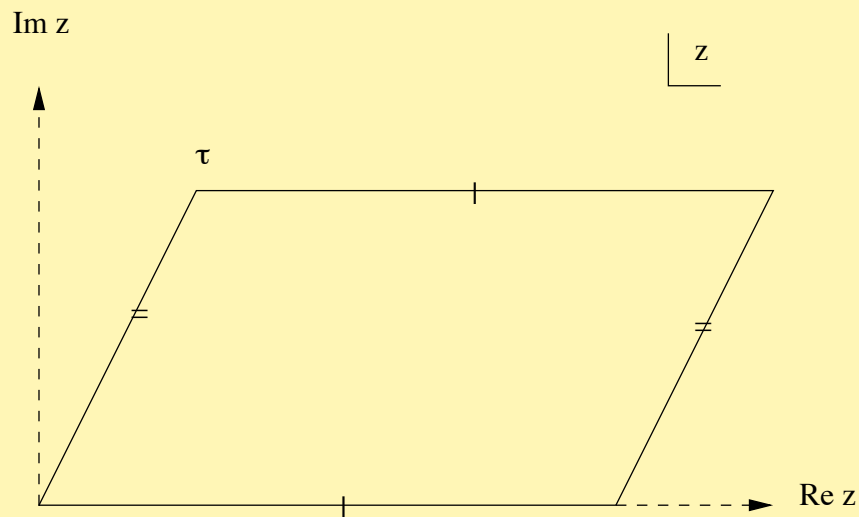
I. Introduction

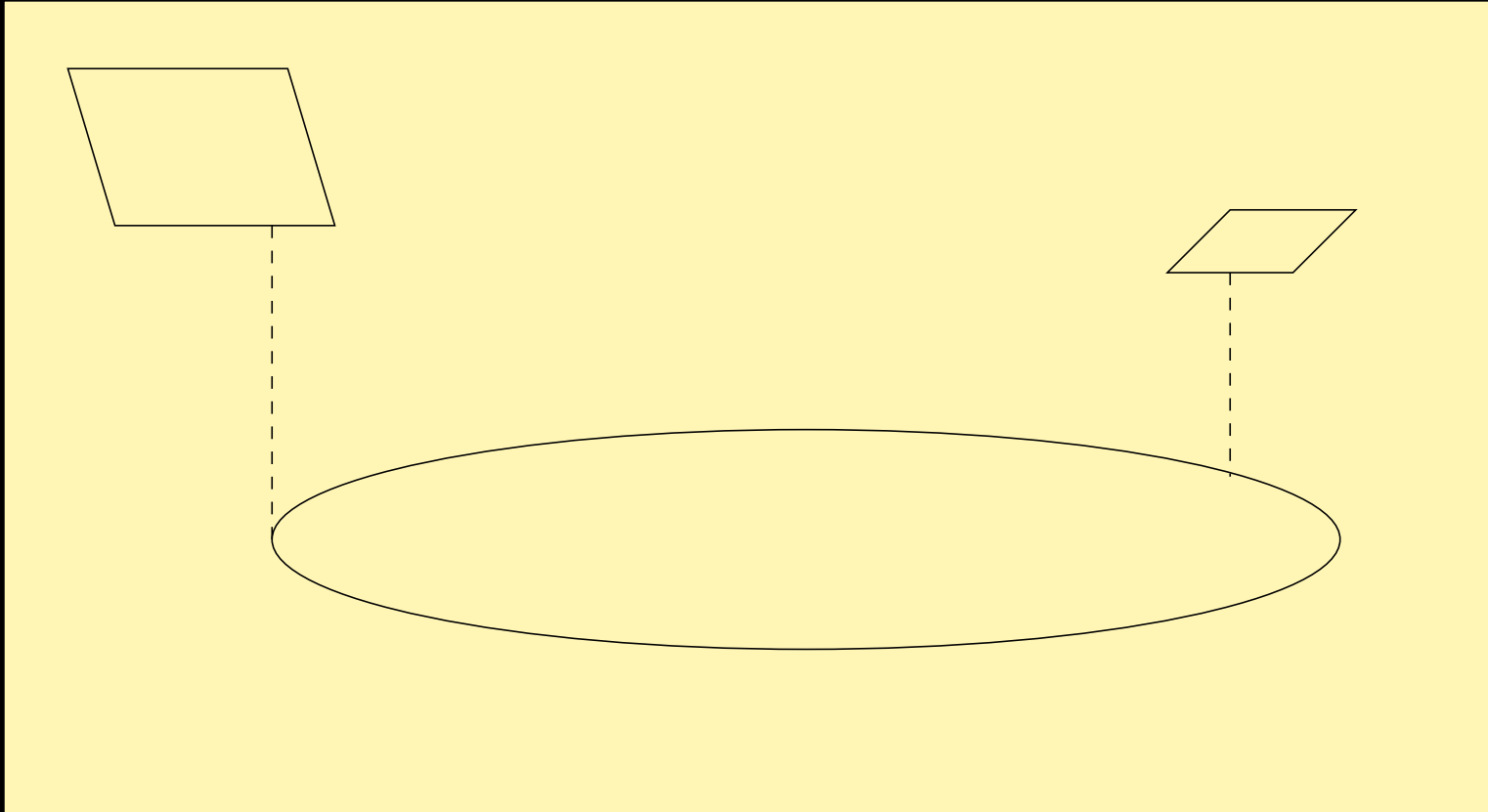
A. Example of non-geometric “T-fold” backgrounds

Compactify string theory to $d = 7$ via NS-NS fields as follows:

Compactification to $d = 8$ on a T^2 . Data:

- Complex structure τ
- Complexified volume $\rho = b + i\sqrt{G}$, $b = \int_{T^2} B_{12}$. b has period 1.





Compactify to $d = 7$ on additional S^1_R , coordinate $x \equiv x + 2\pi R$.

$$\rho \equiv \rho(x)$$

$$\tau \equiv \tau(x)$$

As $x \rightarrow x + 2\pi R$, (τ, ρ) must return to selves up to symmetry of $d = 8$ compactification. The symmetry action is called a "monodromy".

Symmetries of string compactification on T^2 :

- Isometries: $SL(2, \mathbb{Z})_\tau$:

$$\tau \rightarrow \frac{a\tau + b}{c\tau + d} ; \quad a, b, c, d \in \mathbb{Z} ; \quad ad - bc = 1$$

- Stringy symmetries: $SL(2, \mathbb{Z})_\rho$. Includes $b \rightarrow b + 1$, T-duality $\rho \rightarrow -\frac{1}{\rho}$.
- Mirror symmetry $\tau \leftrightarrow \rho$. This is a T-duality transformation $R \rightarrow 1/R$ on one cycle of torus.
- $(\tau, \rho) \rightarrow (-\tau^*, -\rho^*)$

Classification by monodromy

1. Shifts $b \rightarrow b + n$, $n \in \mathbb{Z}$ lead to "magnetic" NS-NS flux:

$$H = dB, \int_{T^3} H = n.$$

2. Shifts in $SL(2, \mathbb{Z})_\tau$ isometries lead to manifolds. "Geometric flux".

Kachru, Schulz, Tripathy, Trivedi;
Tomasiello; Shelton, Taylor, Wecht

3. T-duality shifts such as $\rho \rightarrow \frac{-1}{\rho}$, $\rho \leftrightarrow \tau$ lead to *nongeometric* compactifications. "Nongeometric flux".

KSTT; Hellerman, McGreevy, Williams;
Hull; STW.

Fiberwise T-duality (with rectangular T^2)

1. $\rho \leftrightarrow \tau$ at every z : magnetic flux \rightarrow geometric flux.
2. $\rho \rightarrow \frac{-1}{\rho}$ at every z : magnetic flux \rightarrow nongeometric flux

KSTT,STW,LSW

More general story:

1. T^n fibres: $GL(n, \mathbb{Z})$ monodromies lead to geometric models. $O(n, n; \mathbb{Z})$ monodromies lead to non-geometric compactifications
2. More general manifolds: T^n fibration over more general base manifold B .
 - a. Geometric: $GL(n, \mathbb{Z})$ transition functions
 - b. Non-geometric: $O(n, n; \mathbb{Z})$

Motivation

1. Magnetic fluxes useful for model building GKP, KKLT
2. More general class includes “nongeometric fluxes” STW
3. Will argue: important for understanding SUSY breaking
4. Nongeometric compactifications intrinsically interesting!

II. D-branes on “T-folds”

A. Motivation

1. Add open strings to nongeometric flux models
2. Wrapped D-branes: nonperturbative objects (solitons, instantons)
3. D-branes are probes of $L \sim g_s l_s$: potentially useful for understanding exotic compactifications
Shenker; Kabat&Pouliot; Douglas, Kabat, Pouliot & Shenker

B. Questions

1. What are allowed D-brane configurations?
2. What is geometry or topology of moduli space/
low-energy configuration space of D-branes?

C. D-branes and monodromy

D-branes transform nontrivially under $O(n,n;Z)$
Example: T^2

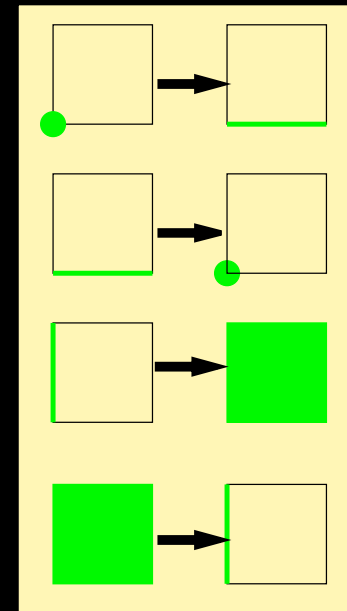
1. T-duality along a cycle exchanges Dirichlet and Neumann

a. D0 \leftrightarrow D1 along cycle

b. D1 along cycle \leftrightarrow D0

c. D1 along dual cycle \leftrightarrow D2

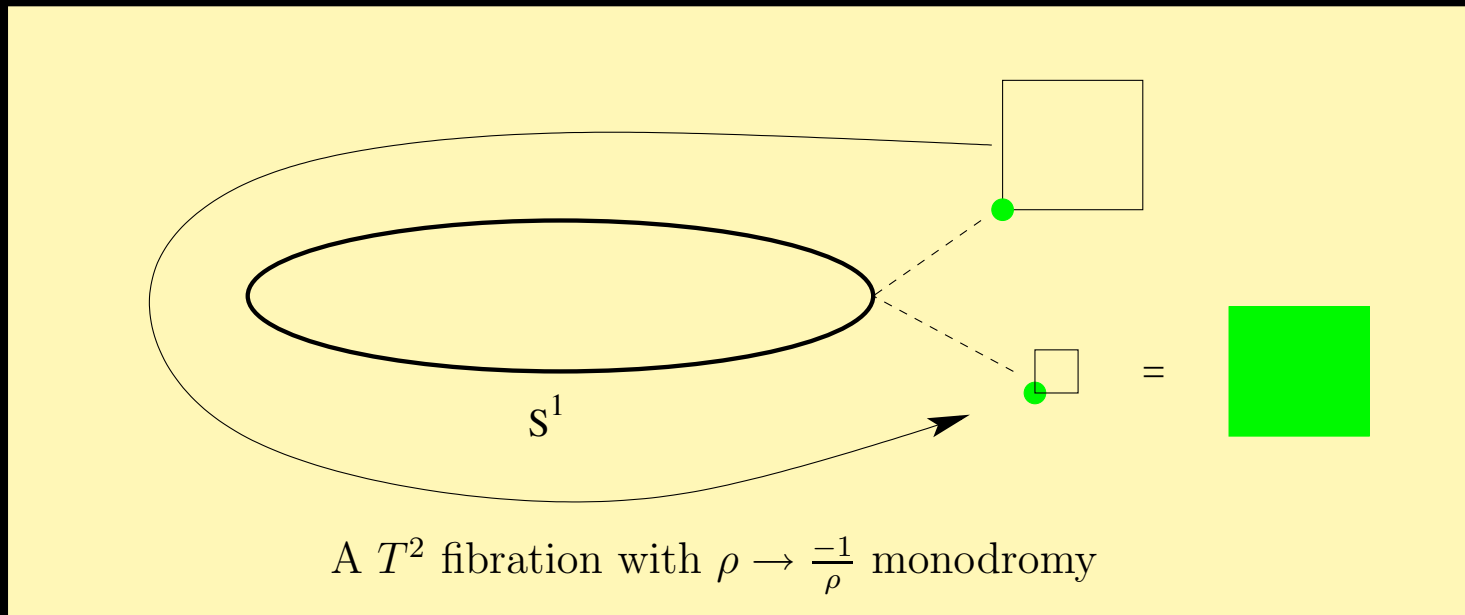
d. D2 \leftrightarrow D1 along dual cycle



2. $b \leftrightarrow b+1$ takes $D2 \leftrightarrow D2 + D0$

D. Allowed D-brane configurations

1. D-brane wrapping base circle: fibre directions must be invariant under monodromy g^n : else it does not close on itself.



This example is not allowed

If fibre directions (and brane orientation) are invariant under g^n , it may be wrapped n times around the base.

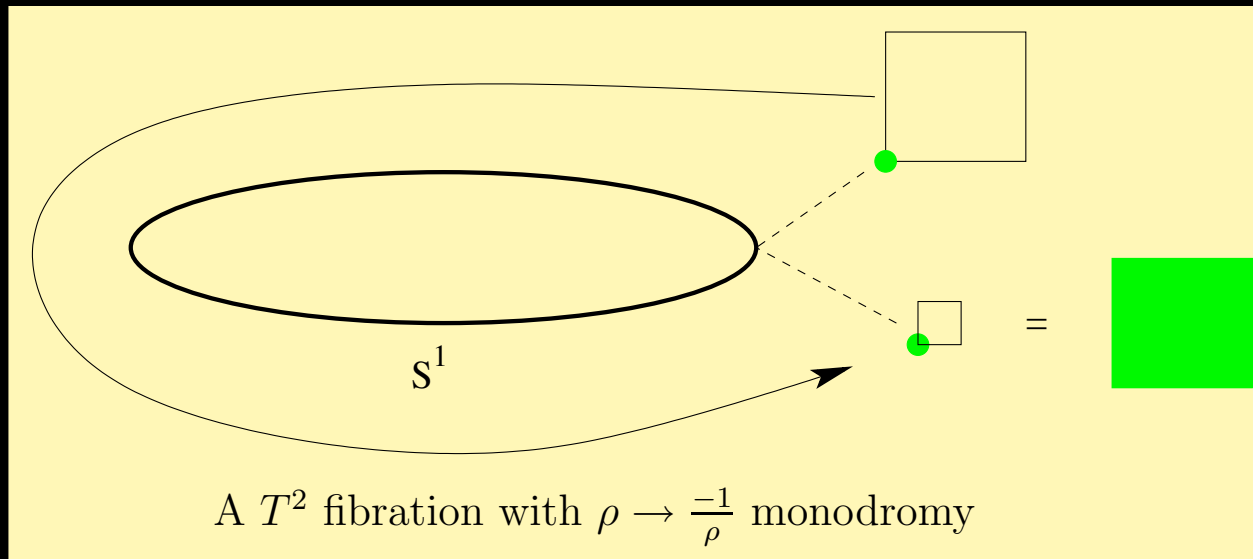
2. D-branes at points on the base are all allowed

Other examples

1. T with H-flux: D3 brane not allowed
2. Non-geometric flux: D1-branes or D3-branes wrapping base not allowed.

Conditions can be stated elegantly using Hull's
“doubled torus” formalism

E. “D-topology” of T-folds



Transport D0-brane around S^1

$$g: D0 \dashrightarrow D2$$

$$g^2: D0 \dashrightarrow D0$$

Configuration space of D0 is
geometric double cover of T-fold

Conjectures:

1. Configuration space of D-branes at point on base is **always** a geometric n -fold cover of “T-fold”
2. If g does not preserve D-brane for any n , there is a potential on the configuration space (else an infinite degeneracy of D-brane states), or brane is otherwise unstable.

(Hellerman, private correspondence)

III. Nongeometric fluxes and SUSY

A. Fluxes and soft SUSY breaking

Consider type IIB on a Calabi-Yau with D-branes.
Lagrangian for open strings (gauge bosons, charged matter) depends on closed strings

1. Perturbative superpotential for open string chiral scalar superfields: couple to complex structure moduli
2. FI D-terms, tree-level gauge couplings: couple to Kahler moduli
3. Kahler potential for open string scalars: couple to all moduli

Brunner, Douglas, Lawrence & Romelsberger;
Douglas; Lawrence & McGreevy

Auxiliary components of closed string fields: soft SUSY-breaking terms in open string Lagrangian

Closed string modes descend from N=2 multiplets.

1. Expand $\mathcal{N} = 2$ superfield in $SU(2)_R$ doublet $(\theta, \hat{\theta})$ of superspace variables.

2. Vector multiplets are chiral in $(\theta, \hat{\theta})$. $SU(2)_R$ triplet of auxiliary fields:

$$V = w + \theta\lambda + \hat{\theta}\hat{\lambda} + \theta^2 D_{++} + \hat{\theta}^2 D_{--} + \theta^\alpha \hat{\theta}^\beta \epsilon_{\alpha\beta} (D_{+-} + \sigma_{\alpha\beta}^{\mu\nu} F_{\mu\nu}) + \dots$$

Here D_{ab} are auxiliary fields.

Grimm, Sohnius & West; de Wit & van Holten; de Roo,
van Holten, de Wit & van Proeyen

3. Hypermultiplets are "twisted chiral" :

$$H = t + \theta\psi + \hat{\theta}\hat{\psi} + \theta^2 y + \hat{\theta}^2 \bar{y} + \theta^\alpha \hat{\theta}^\beta \sigma_{\alpha\beta}^\mu F_\mu + \dots$$

where $t, y, \bar{y}, F_\mu = \partial_\mu \phi$ are complex. Here y, \bar{y} are auxiliary fields.

Berkovits & Siegel

Only certain auxiliary fields are understood

Type IIB vector multiplets

- w : complex structure deformations.
- $J^\mu{}_\nu$: almost complex structure; $\omega = g_{\mu\lambda} J^\lambda{}_\nu dx^\mu \wedge dx^\nu$.
- $D_{\pm\pm}$: built from $d\omega$ and NS-NS 3-form field strength H , both $\in H^{(2,1)} \oplus H^{(1,2)}$.
- D_{+-} are built from RR 3-form F in $H^{(2,1)} \oplus H^{(1,2)}$.

Vafa;
Lawrence
&McGreevy

Type IIA hypermultiplets

- t are complex structure deformations.
- y, \bar{y} built from $d\omega$, $H \in H^{(2,1)} \oplus H^{(1,2)}$.

Lawrence
&McGreevy

What about IIB hypermultiplets, IIA vectormultiplets?

Mirror symmetry for NS-NS flux?

1. Mirror symmetry exchanges IIA and IIB, Kahler and complex structure moduli.
2. y, \bar{y} in IIB should be “mirrors of NS flux”
3. Mirror symmetry is a form of T-duality for most Calabi-Yau compactifications Strominger, Yau & Zaslow
4. T-duality applied to H-flux: geometric, non-geometric fluxes

Worksheet calculation in sigma model limit

- $\Omega \in H^{(3,0)}$ is holomorphic 3-form which (together with ω determines metric.
- y, \bar{y} built from H with all holomorphic indices and $d\Omega$ with 2 holomorphic indices. (= Particular class of "intrinsic torsion".)

Vafa; Gurrieri. Louis, Micu &
Waldrum; Gurrieri & Micu;
Fidanza, Minasian & Tomasiello;
LMSSW

Puzzles

1. How does this relate to nongeometric flux? Does that emerge globally?
2. y, \bar{y} in IIA allow for spacetime SUSY in supergravity approximation. y, \bar{y} in IIB do not.

Answer: worldsheet instantons correct IIB SUSY conditions in the presence of y, \bar{y}

LMSSW

Conclusions

1. Nongeometric compactifications lead to interesting modification of stringy topology
2. Nongeometric “flux” is generic and **important** in type II models with reduced/broken SUSY
3. Worldsheet instantons always crucial for mirror symmetry