

Determining the Neutrino Mass Hierarchy with MINOS

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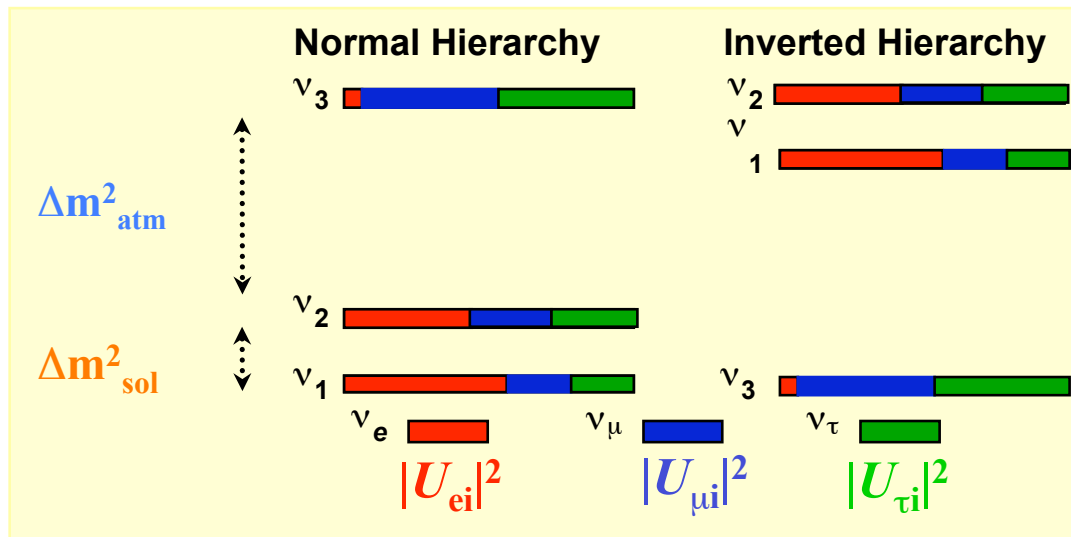
Honolulu, October 29-Nov 3 2006

Outline

- Neutrino oscillations in matter
- Matter induced neutrino-antineutrino asymmetry as a tool to determine the neutrino mass hierarchy
- Measuring the mass hierarchy at MINOS:
 - MC study
 - Results
- Summary

Neutrino Mass Spectrum

(3 mass eigenstates - assume no ν_{sterile})



Two mass splittings:

Large (Δm^2_{atm}):

Atmospheric (up-down asymmetry)
Long baseline

Small (Δm^2_{sol}):

Solar + reactor

- Solar data suggests $|U_{e2}|^2 \sim 1/3$, and thus $|U_{e1}|^2 \sim 2/3$ (unitarity and small $|U_{e3}|^2$).
- Also, data suggests $\nu_\mu \rightarrow \nu_\tau$ dominant at Δm^2_{atm} and maximal mixing angle,

$$|U_{\mu i}|^2 \sim |U_{\tau i}|^2$$

➔ $|U_{e3}|^2$ is not known (bounded by reactor expts) and neither is the sign of Δm^2_{atm} which determines whether the ν mass hierarchy is normal or inverted.

Matter Effects

- ◆ ν propagating in matter interact through W (ν_e only) or Z (all flavors) exchange, giving rise to an interaction potential:

$$V_W = +\sqrt{2} G_F N_e \quad (W \text{ exchange})$$

- ◆ The sign of V is opposite for $\bar{\nu}$

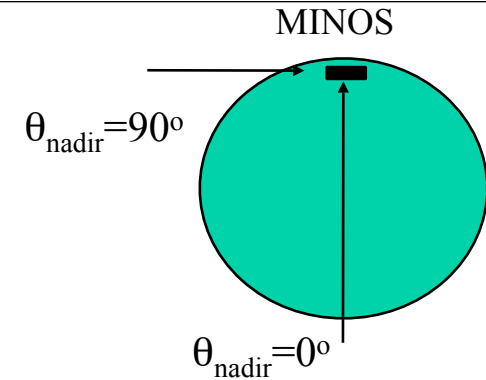
➔ Oscillation probabilities in matter are modified

$$P_M(\nu_e \rightarrow \nu_\mu) = \sin^2 2\theta_M \sin^2\left(\Delta m_M^2 \frac{L}{4E}\right) \quad x \equiv \frac{V_W/2}{\Delta m^2/4E} = \frac{2\sqrt{2}G_F N_e E}{\Delta m^2}$$

$$\sin^2 2\theta_M \equiv \frac{\sin^2 2\theta}{\sin^2 2\theta + (\cos 2\theta - x)^2} \quad \Delta m_M^2 \equiv \Delta m^2 \sqrt{\sin^2 2\theta + (\cos 2\theta - x)^2}$$

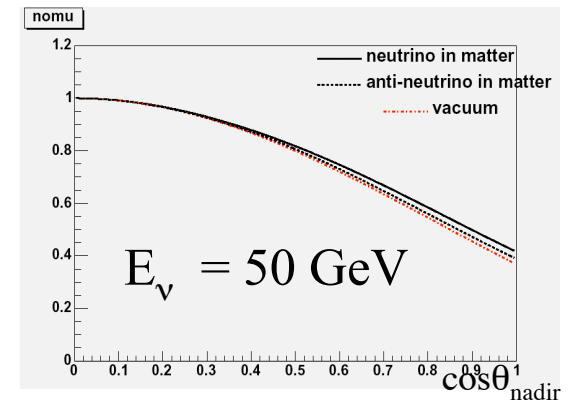
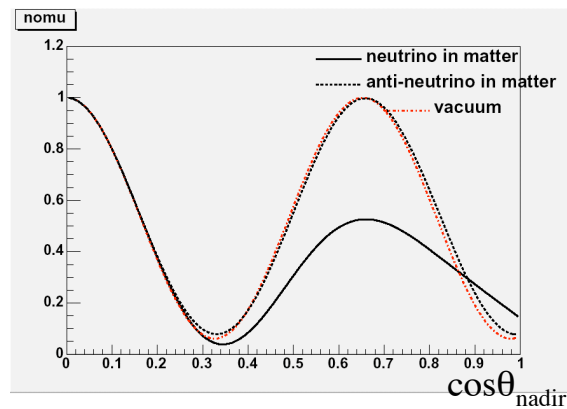
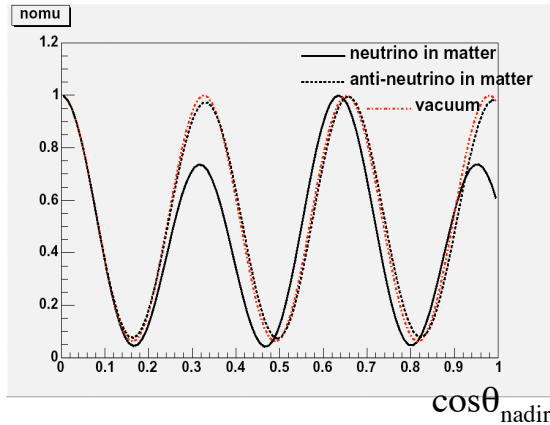
- ➔ Matter alters ν , $\bar{\nu}$ oscillations differently
- ➔ Depending on mass hierarchy, ν or $\bar{\nu}$ oscillations are enhanced (or suppressed) \implies matter induced asymmetry
- ➔ Resonant effects occur

Oscillation Probabilities in Matter (mantle only)

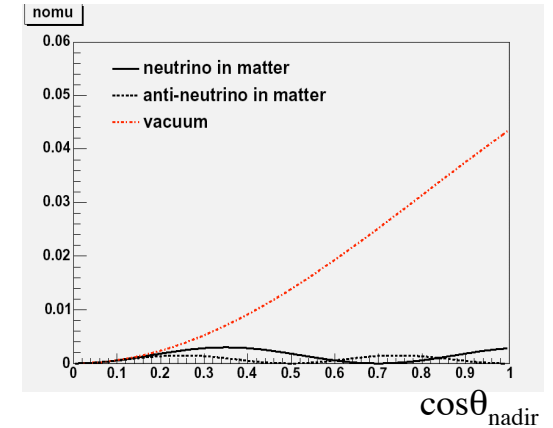
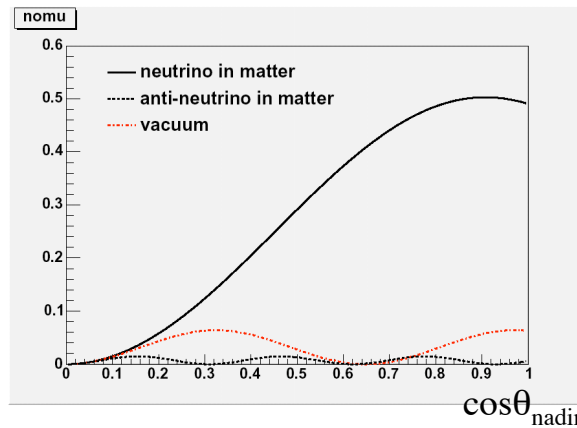
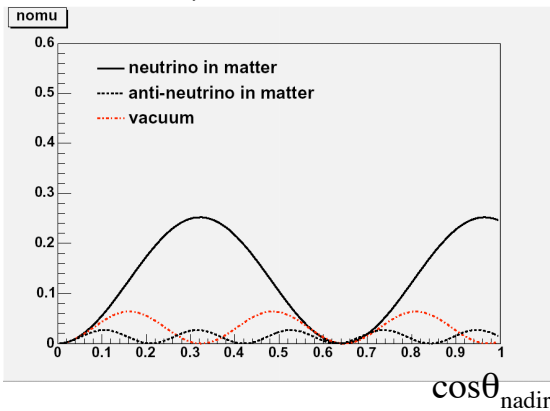


$$\nu_\mu \rightarrow \nu_\mu \quad E_\nu = 5 \text{ GeV}$$

$$E_\nu = 10 \text{ GeV}$$



$$\nu_e \rightarrow \nu_\mu$$

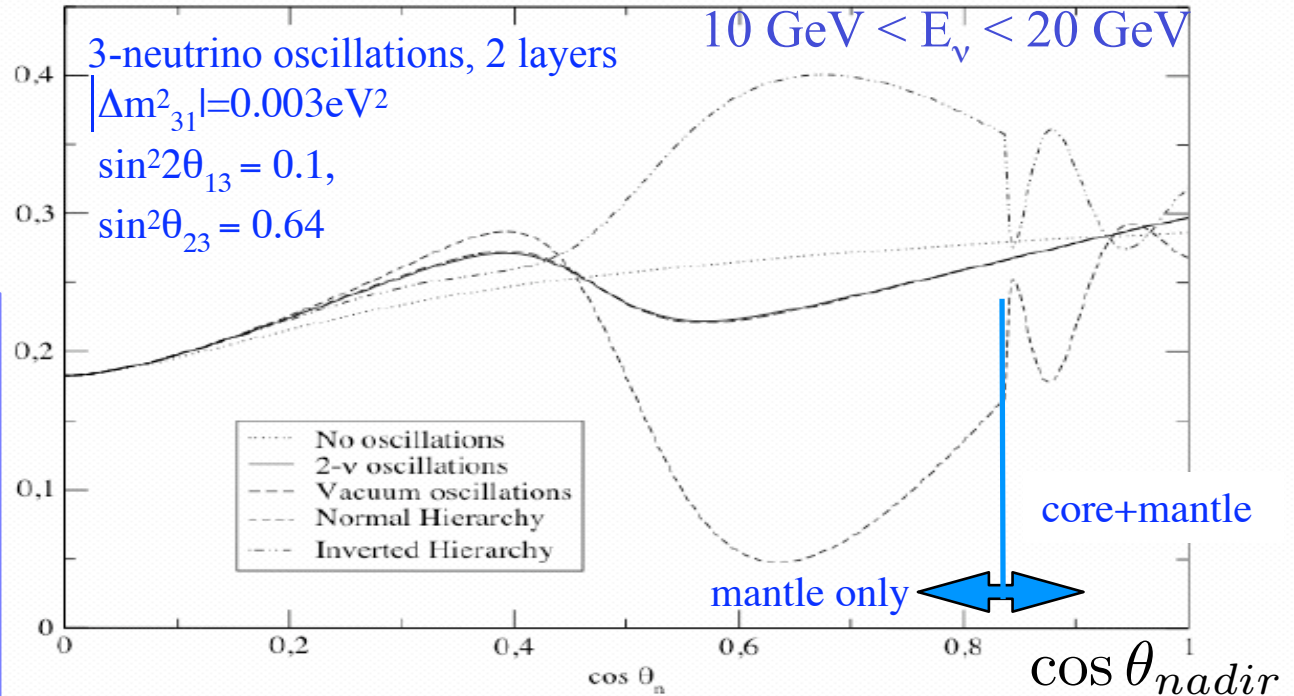
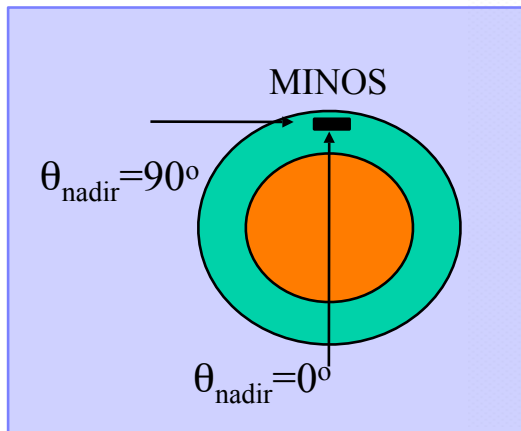


$$|\Delta m_{31}^2| = 0.003 \text{ eV}^2, \quad \sin^2 2\theta_{13} = 0.1, \quad \sin^2 \theta_{23} = 0.64$$

Charge Asymmetry-Atmospheric ν

(Ref: J. Bernabeu, S. Palomares, and S. Petcov, hep-ph/0305152)

$$\frac{N(\mu^-) - N(\mu^+)}{N(\mu^-) + N(\mu^+)}$$



- The charge asymmetry for neutrino induced muons can be significantly modified by matter effects and it could carry the signature of the underlying mass hierarchy
- Resonant enhancement occurs for $E_\nu \sim 7-11 \text{ GeV}$ and $\cos \theta_{nadir} > 0.4$

Sensitivity Study

● The MINOS far detector is the first large underground detector with a magnetic field and thus able to discriminate the muon charge

➡ **We explored the possibility of using the atmospheric ν data collected at MINOS to discriminate between normal and inverted hierarchy.**

● We determine the sensitivity using simulated atmospheric neutrino events reconstructed at the MINOS far detector

● We consider different topologies:

Contained vertex ν events: ν interaction occurs in the detector (6,500 years exposure in the MC)

Upward-going μ : ν interaction occurs in rock (2,500 years exposure in the MC)

MC Event Selection

- Select upward going CC ν_μ interactions with muon energy > 2 GeV
- Require one track. Track must:
 - ◆ pass the fitter internal consistency checks; have fit $\chi^2/ndf < 2$
 - ◆ contain 40% of the total PH in the event

Contained vtx ν :

- Track vertex in fiducial volume (50 cm from edge and 5 planes from SM boundaries)
- track crosses >10 planes and is >1 meter long
- $\sigma_{qp}/qp < 0.4$ or track length > 8 m for better charge separation

Up μ :

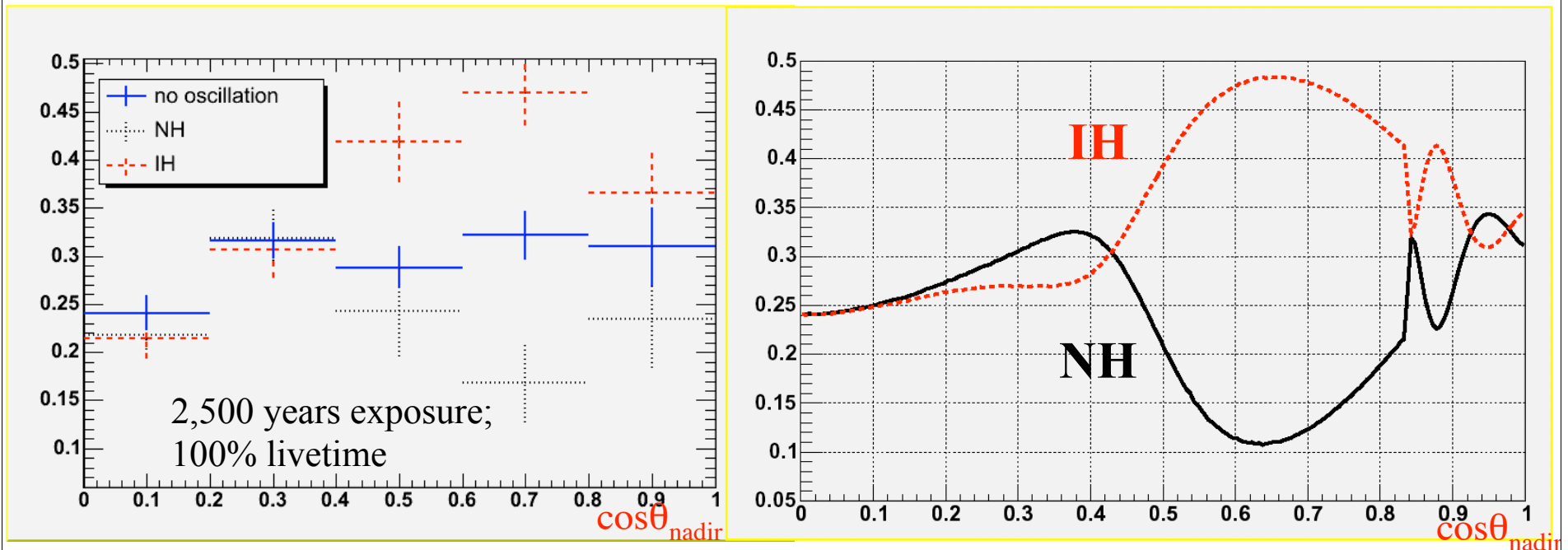
- Track vertex out of fiducial volume (30 cm from edge and 15 cm from SM boundaries)
- Track crosses >20 planes and is >2 meter long
- Track contains 40% of the total number of planes in the event
- χ_{line}^2/ndf (measures deviation of track from straight line) > 10 for better charge separation

→ No significant improvement in mass hierarchy discrimination with tighter cuts
but 15~30% reduction in statistics

Charge Asymmetry - MC

Contained vtx ν with $10 \text{ GeV} < E_\nu < 20 \text{ GeV}$

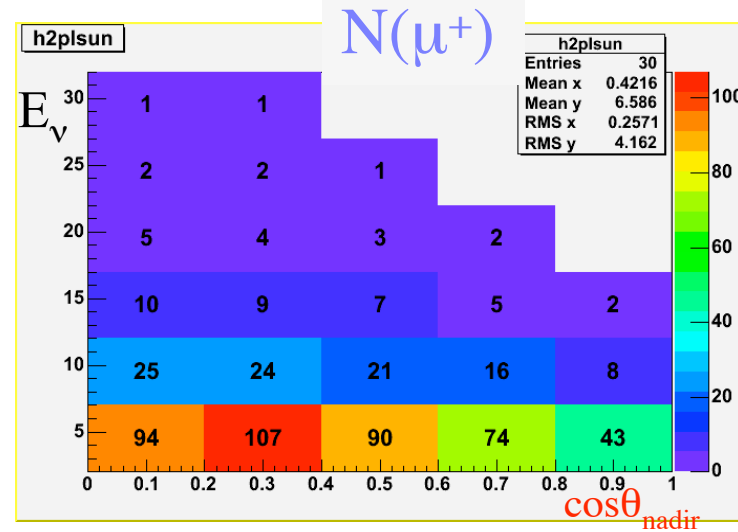
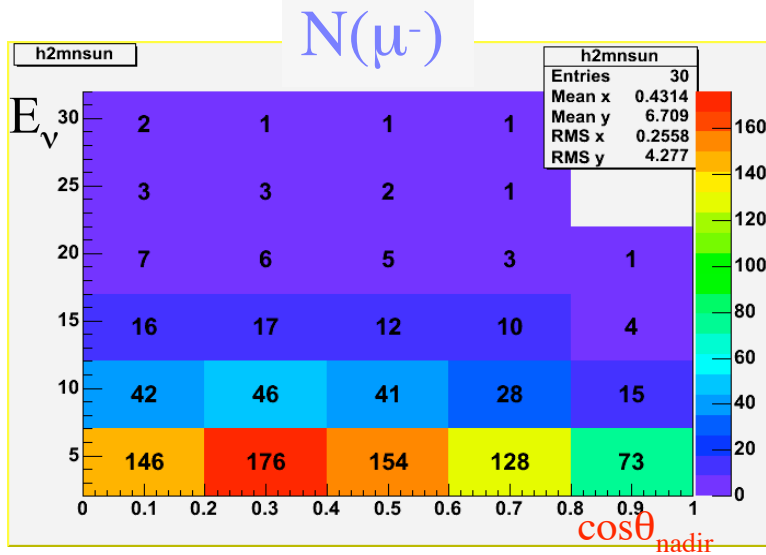
$$|\Delta m_{31}^2| = 0.003 \text{ eV}^2, \sin^2 2\theta_{13} = 0.1, \sin^2 \theta_{23} = 0.64$$



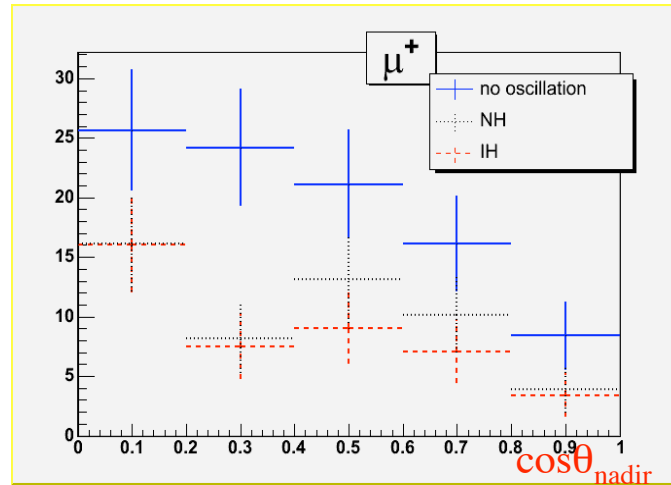
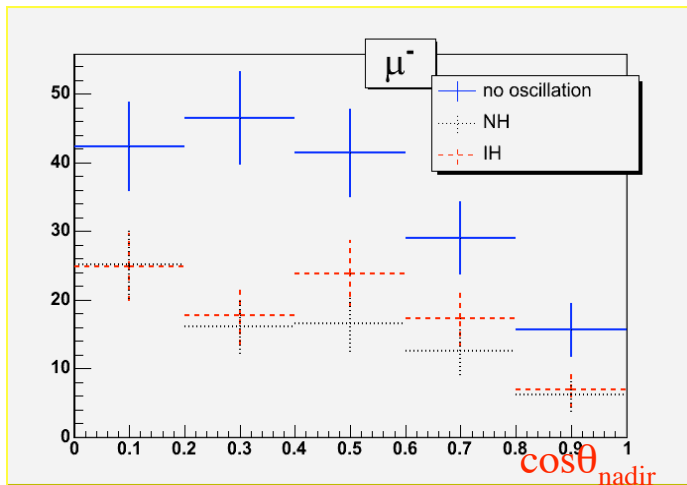
Good agreement between MC charge asymmetry and theory prediction

Event Rates - 100 yrs

$N(\mu^-)$ and $N(\mu^+)$ after 100 years, contained vertex ν , no oscillation:

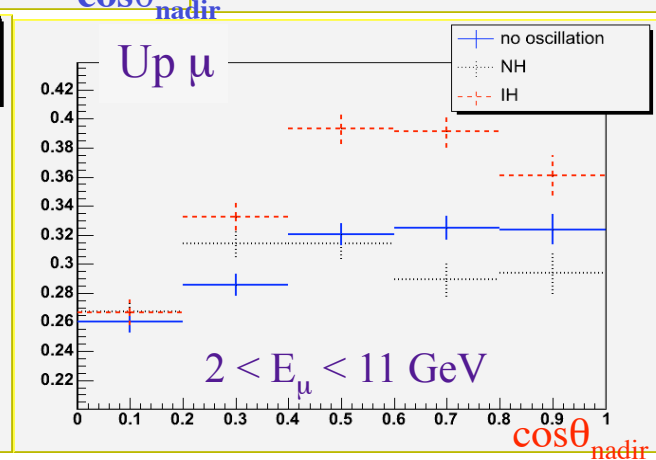
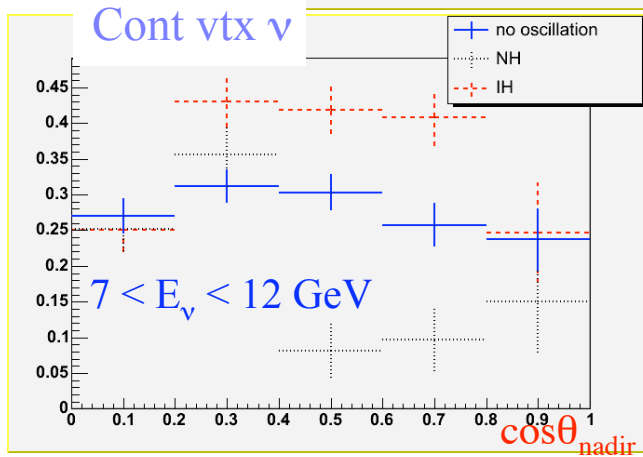
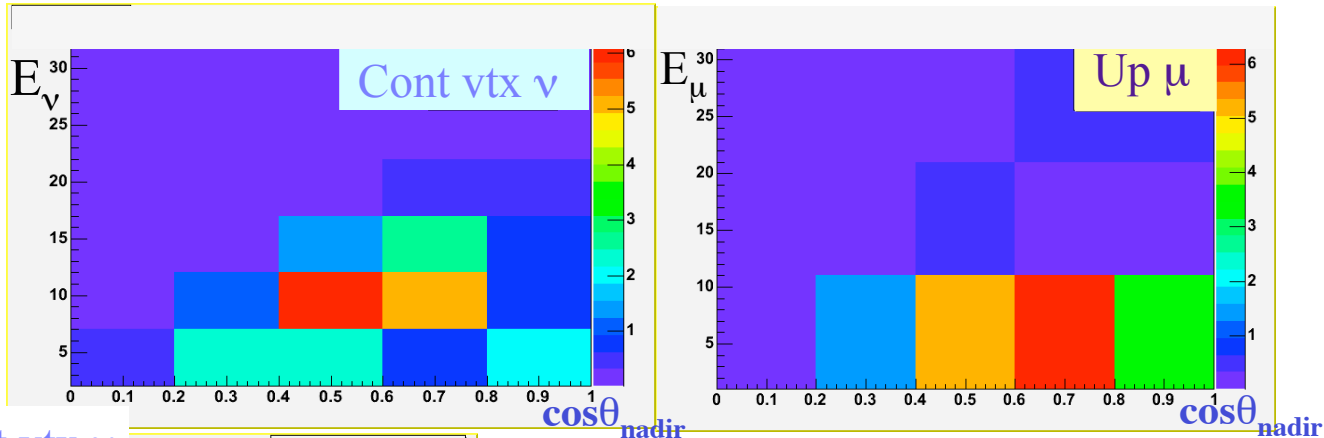


$N(\mu^-)$ and $N(\mu^+)$ vs $\cos\theta_{\text{nadir}}$ for $7 \text{ GeV} < E_\nu < 12 \text{ GeV}$:



Statistical Significance

➔ Statistical significance $|A_{IH} - A_{NH}|/\sigma$ after **2,500 years**:



For other exposures:

$\frac{(A_{IH} - A_{NH})}{\sigma}$	# of years
1σ	52
1.4σ	100
2σ	203
3σ	456
5σ	1266

➔ Comparable discrimination can be achieved for the two topologies

Likelihood Study

- Build likelihood function L to discriminate between normal and inverted hierarchy.
- For a given set of data, the ratio $R = L_{\text{NH}}/L_{\text{IH}}$ measures whether the data is more consistent with the NH or IH hypothesis.
- Generate large number of pseudoexperiments for a given exposure assuming NH (IH) and evaluate R_{NH} (R_{IH}).
- In practice, calculate $L = -2\ln(L)$ and determine $\Delta L = L_{\text{NH}} - L_{\text{IH}}$ (the test statistics)
- The larger the separation between the two distributions (ΔL_{NH} and ΔL_{IH}), the better the two mass hierarchies can be discriminated.
- Determine the 95% C.L. sensitivity for rejecting IH(NH) by calculating the fraction of pseudoexperiments generated assuming NH (IH) with less than 5% IH(NH) contamination.
- Repeat for different exposures and values of $\sin^2 2\theta_{13}$ and $\sin^2 \theta_{23}$

Likelihood Function

- Unbinned likelihood for contained vtx ν and upward going μ :

Poisson constraints on total numbers of events

$$L = \frac{(\lambda^-)^{N^-} e^{-\lambda^-}}{N^-!} \times \prod_{i=1}^{N^-} \underbrace{f^-(E_i, \cos \theta_{nadir,i})}_{\text{Shape constraints}} \times \frac{(\lambda^+)^{N^+} e^{-\lambda^+}}{N^+!} \times \prod_{i=1}^{N^+} \underbrace{f^+(E_i, \cos \theta_{nadir,i})}_{\text{Shape constraints}}$$

N^\pm : observed number of μ^\pm .

λ^\pm : expected number of μ^\pm . Depend on oscillation parameters.

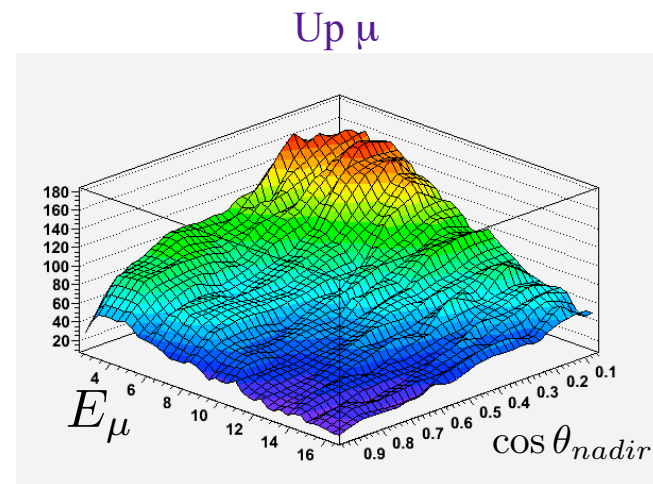
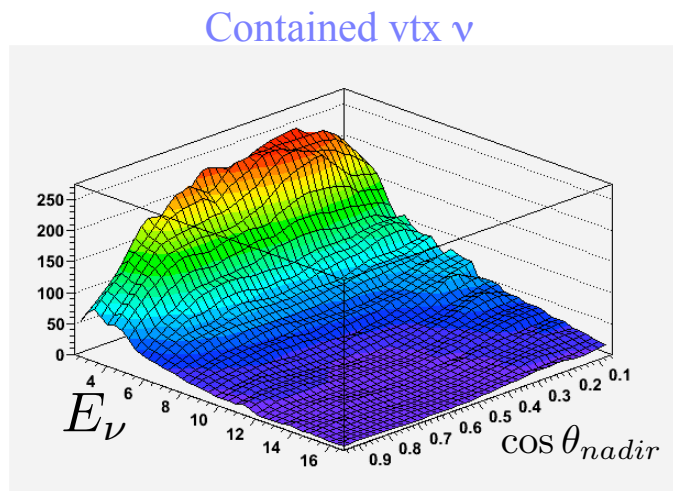
f^\pm : normalized shapes of μ^\pm rates as a function of $\cos \theta_{nadir}$ vs E distributions.
Depend on oscillation parameters.

- Fix all oscillation parameters, except for the sign of Δm^2_{31} and consider the following values of:

$$\sin^2 2\theta_{13} = 0.05, 0.1, 0.15 \quad \sin^2 \theta_{23} = 0.5, 0.64$$

Pseudoexperiments

- Generate number of events (μ^+, μ^- , for cont vtx ν and up μ) according to Poisson statistics for a given exposure.
- Interpolate $\cos\theta_{nadir}$ vs energy MC distributions and randomly generate $\cos\theta_{nadir}$, E for each event according to these shapes and both hierarchies.



$$\sin^2 2\theta_{13} = 0.1, \sin^2 \theta_{13} = 0.64, \mu^-, NH$$

- Determine ΔL_{NH} and ΔL_{IH} for each pseudoexperiment
- Repeat for several exposures and values of the mixing angles

ΔL for Cont vtx ν -10 yrs

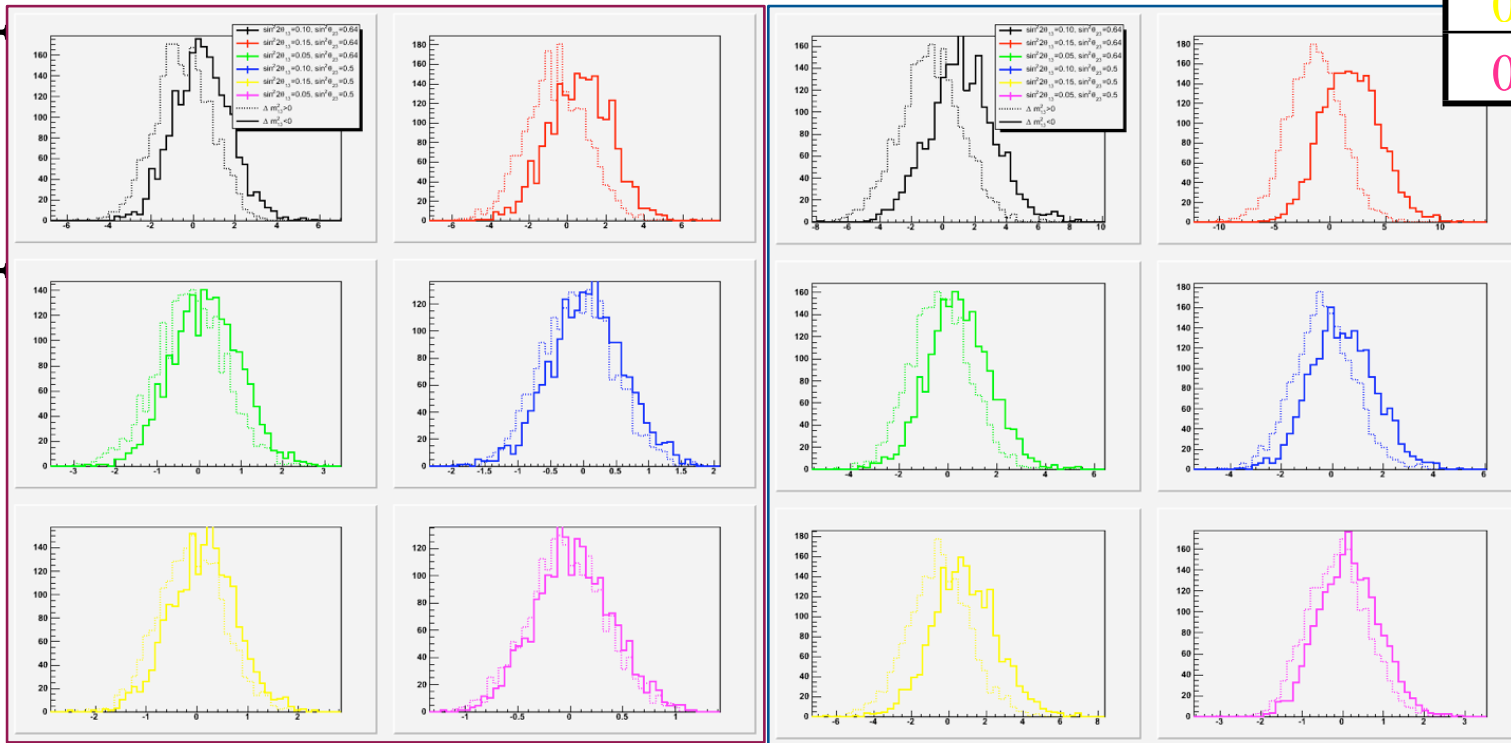
* Compare ΔL for NH (dashed) and IH (solid) hypotheses

$\sin^2 2\theta_{13}$	$\sin^2 \theta_{23}$
0.10	0.64
0.15	0.64
0.05	0.64
0.10	0.50
0.15	0.50
0.05	0.50

Poisson constraint only

Poisson+Shape constraint

of pseudoexpts



— $\Delta m_{31}^2 < 0$
 $\Delta m_{31}^2 > 0$

ΔL

➔ Better discrimination achieved by making use of shape information in the likelihood

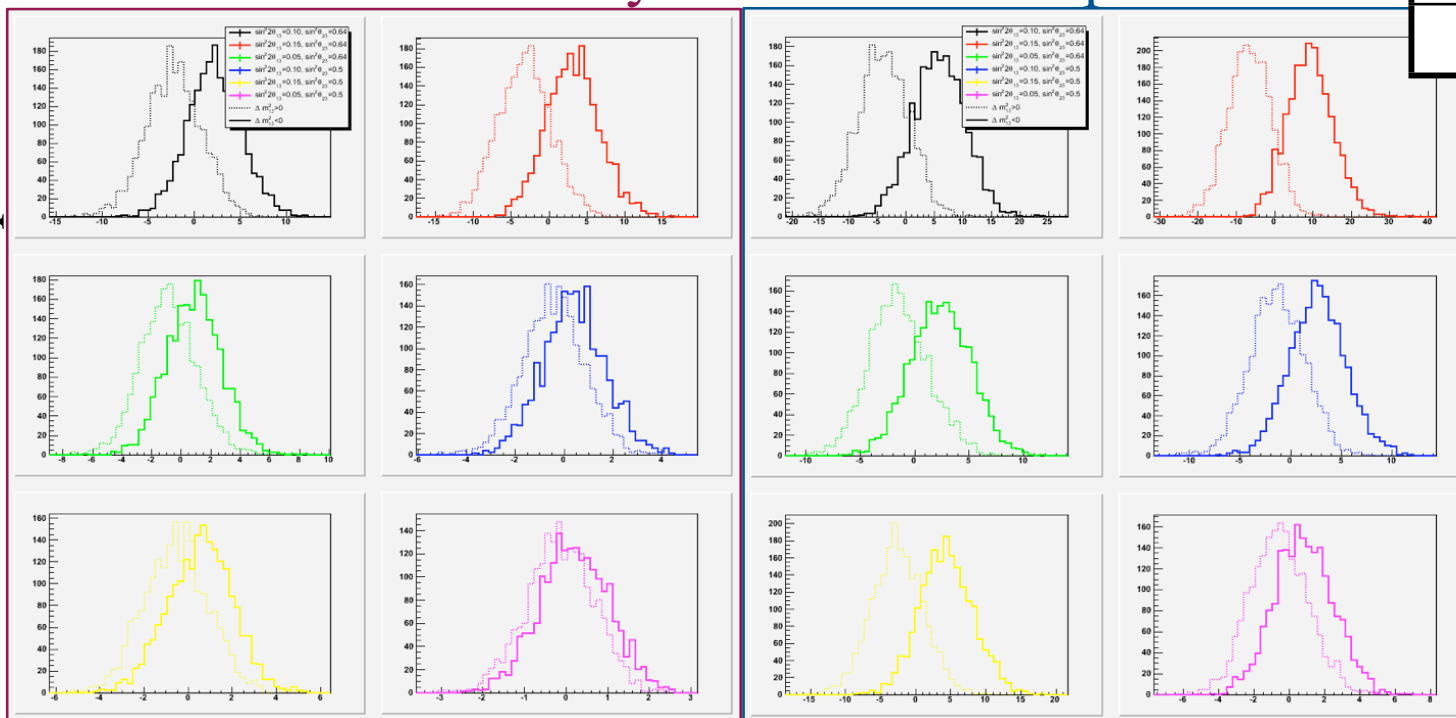
ΔL for Cont vtx ν -50 yrs

$\sin^2 2\theta_{13}$	$\sin^2 \theta_{23}$
0.10	0.64
0.15	0.64
0.05	0.64
0.10	0.50
0.15	0.50
0.05	0.50

of pseudoexpts

Poisson constraint only

Poisson+Shape constraint



— $\Delta m_{31}^2 < 0$
 $\Delta m_{31}^2 > 0$

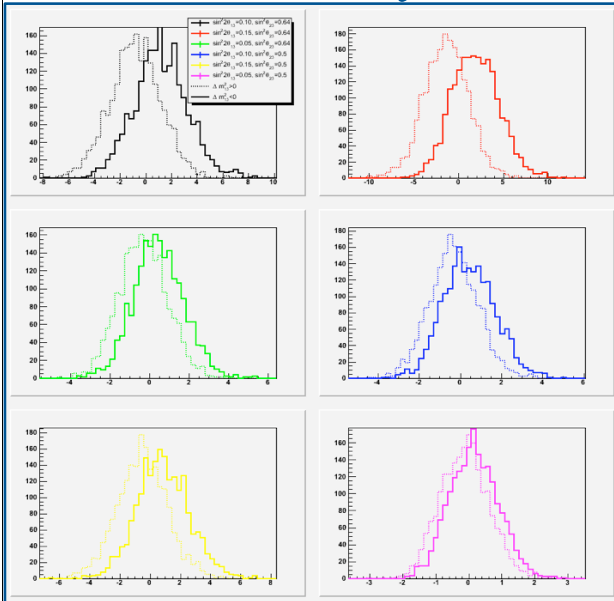
ΔL

➔ As expected, discrimination improves for larger $\sin^2 2\theta_{13}$

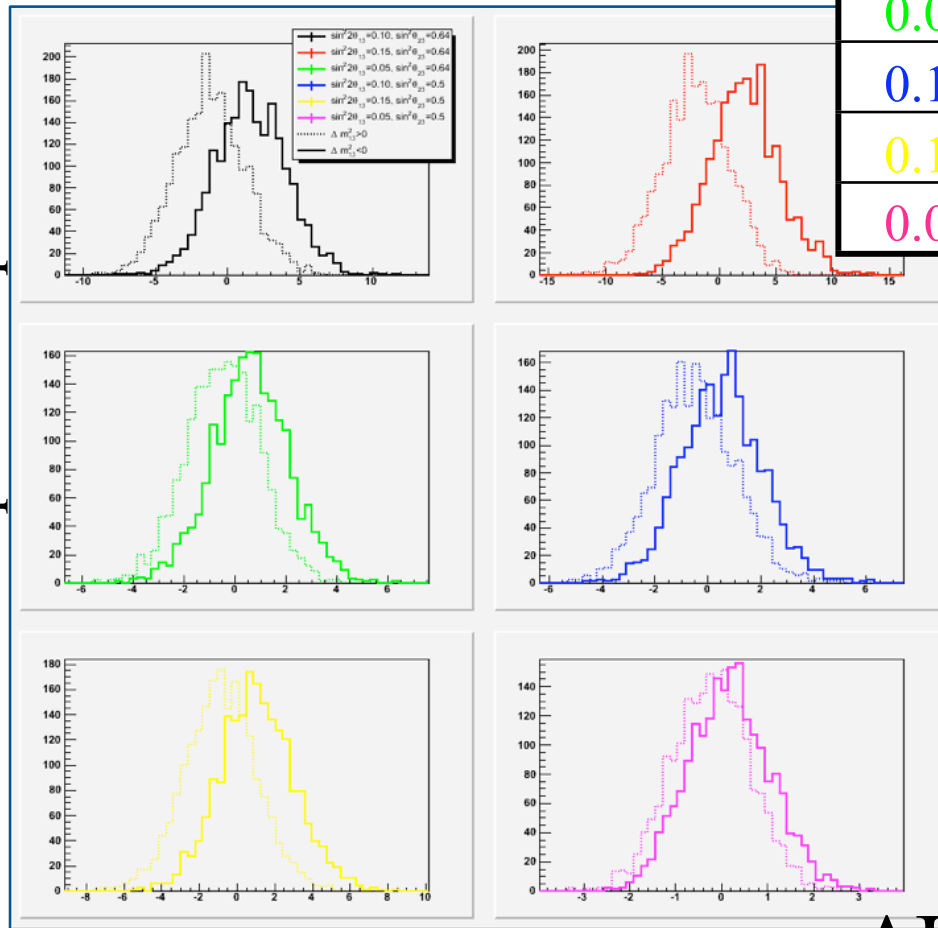
Combined ΔL -10 yrs

$\sin^2 2\theta_{13}$	$\sin^2 \theta_{23}$
0.10	0.64
0.15	0.64
0.05	0.64
0.10	0.50
0.15	0.50
0.05	0.50

Cont vtx ν only



Combined



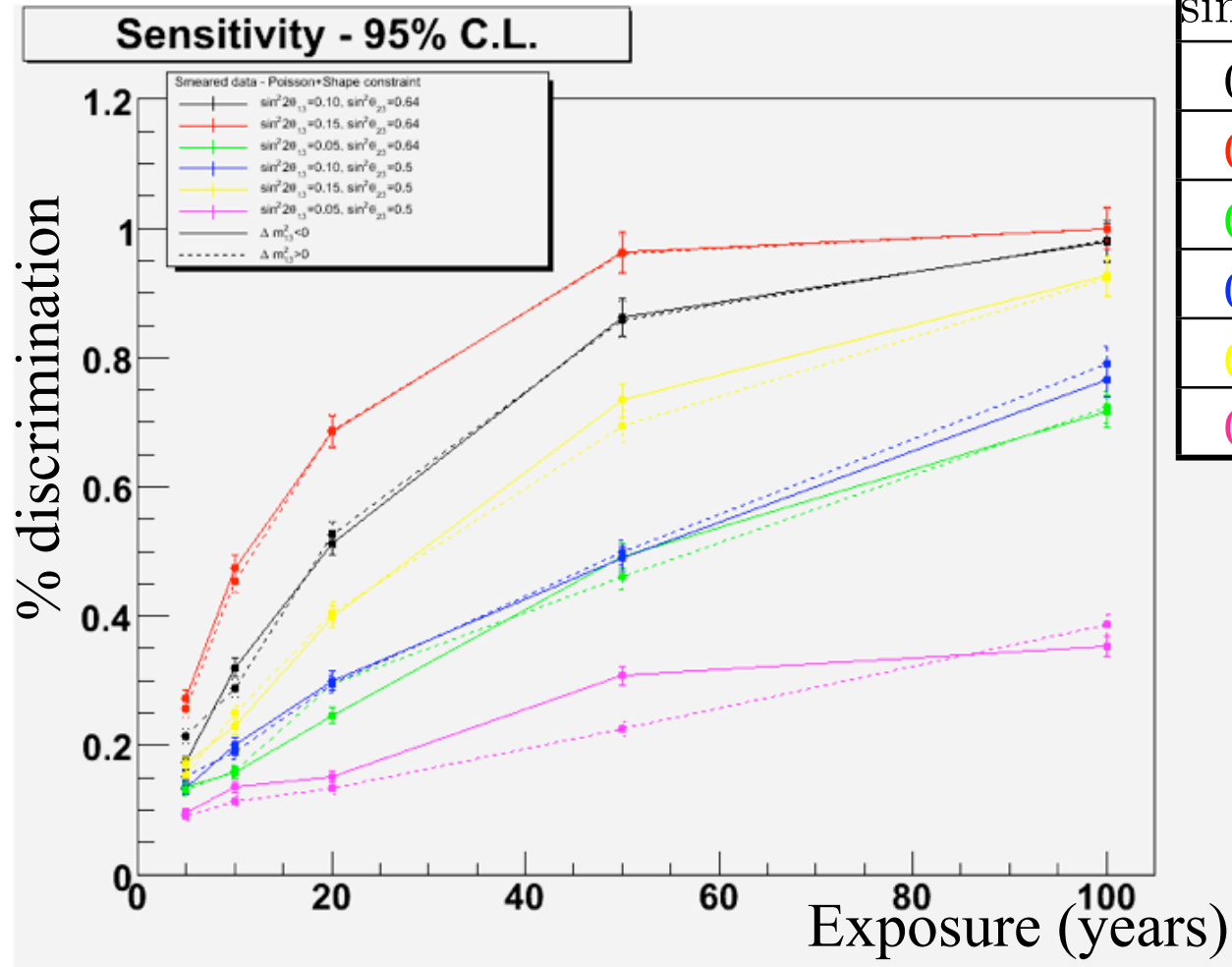
of pseudoexpts

— $\Delta m_{31}^2 < 0$
 $\Delta m_{31}^2 > 0$

ΔL

➔ Further improvement is obtained by combining the contained vertex ν sample with the upgoing μ sample

Sensitivity-Combined



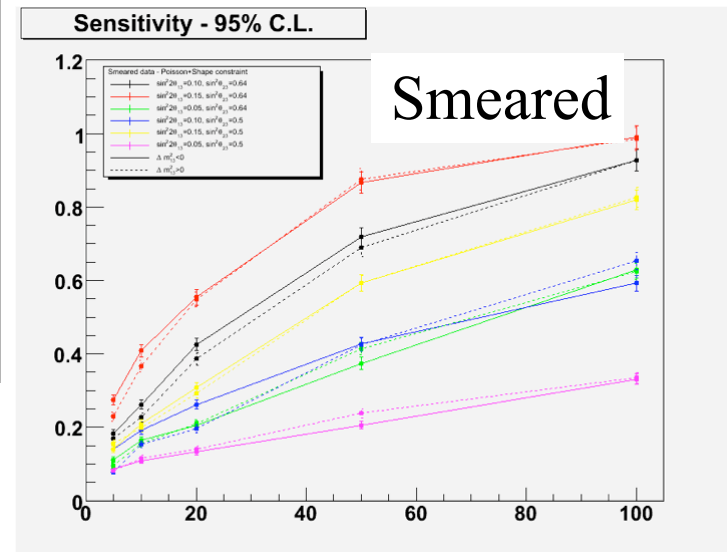
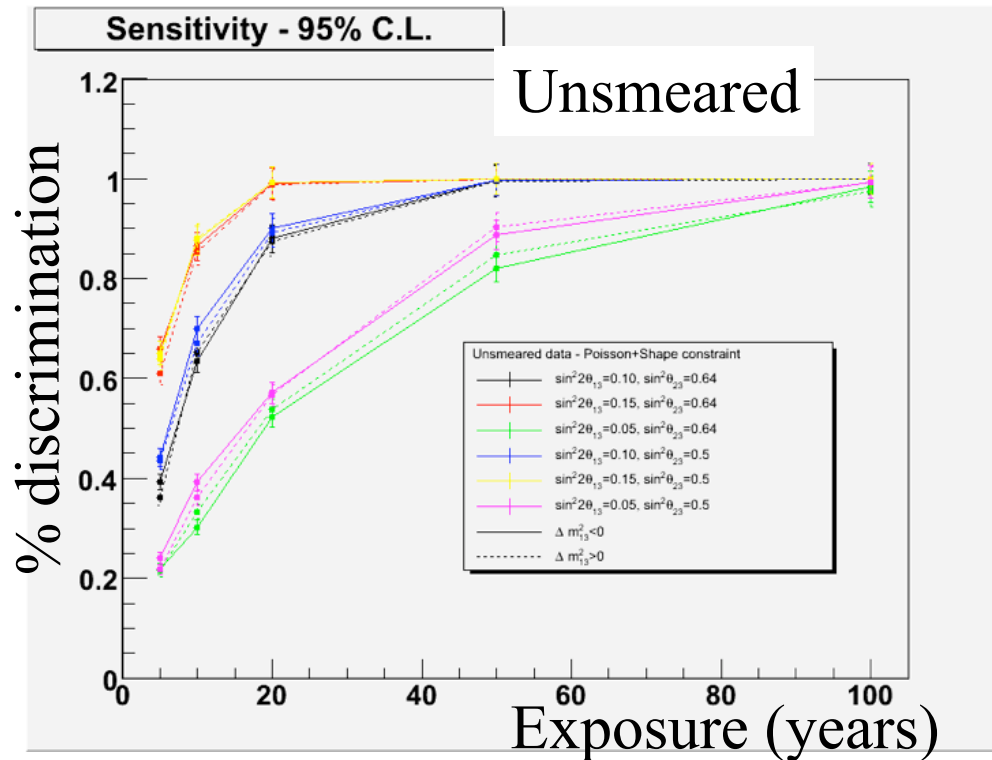
- ➔ Similar sensitivity achieved for NH and IH
- ➔ Need over 10 years of running, if oscillation parameters are favorable, for a >50% chance to determine the mass hierarchy at 95% C.L.

Sensitivity-Unsmeared

Contained vtx ν only

$\sin^2 2\theta_{13}$	$\sin^2 \theta_{23}$
0.10	0.64
0.15	0.64
0.05	0.64
0.10	0.50
0.15	0.50
0.05	0.50

— $\Delta m_{31}^2 < 0$
 $\Delta m_{31}^2 > 0$



Summary

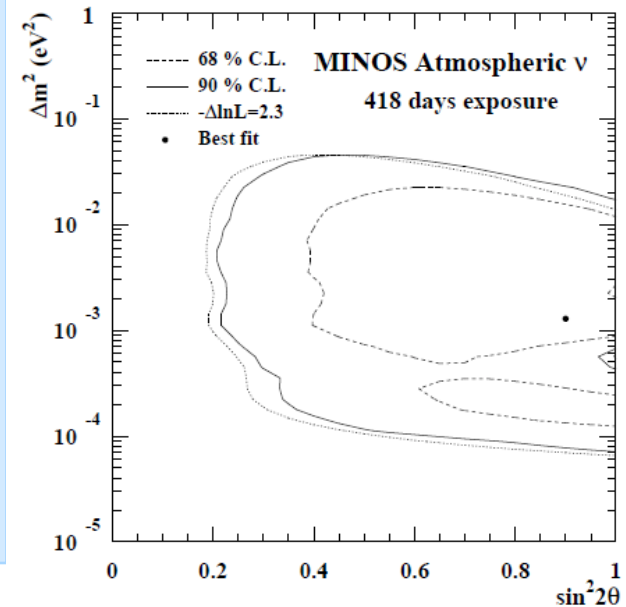
- We have studied the sensitivity of the MINOS far detector to discriminate the neutrino mass hierarchy using simulated atmospheric neutrinos (contained vertex ν events and upward going μ) reconstructed at the far detector
- We have determined the sensitivity as a function of the exposure, $\sin^2 2\theta_{13}$, $\sin^2 \theta_{23}$ and the sign of Δm_{31}^2
- We find that the sensitivity is comparable for NH and IH, and it improves significantly if $\sin^2 2\theta_{13}$ is large
- The analysis reveals that >10 years of exposure at the MINOS far detector would be needed to discriminate between mass hierarchies at 95% C.L.

BACKUP

Far Detector : Atmospheric Neutrino Oscillations

Phys. Rev. D73, 072002 (2006)

- The MINOS far detector is the first large magnetized underground detector. Installation completed in July 2003.
- Veto shield used to reject cosmic rays for atmospheric ν analysis
- L/E analysis excludes no oscillations at 98% CL
- First measurement of atmospheric ν_μ vs. $\bar{\nu}_\mu$ oscillations



$$\frac{R_{up/down}^{data}}{R_{up/down}^{MC}} = 0.62^{+0.19}_{-0.14}(\text{stat}) \pm 0.02(\text{syst})$$

$$\frac{R_{\bar{\nu}_\mu/\nu_\mu}^{data}}{R_{\bar{\nu}_\mu/\nu_\mu}^{MC}} = 0.96^{+0.38}_{-0.27}(\text{stat}) \pm 0.15(\text{syst})$$



MC assumes ν and $\bar{\nu}$ oscillate in the same way

