

Axiomatic Bounds on New Physics
or
What we WONT see at the LHC
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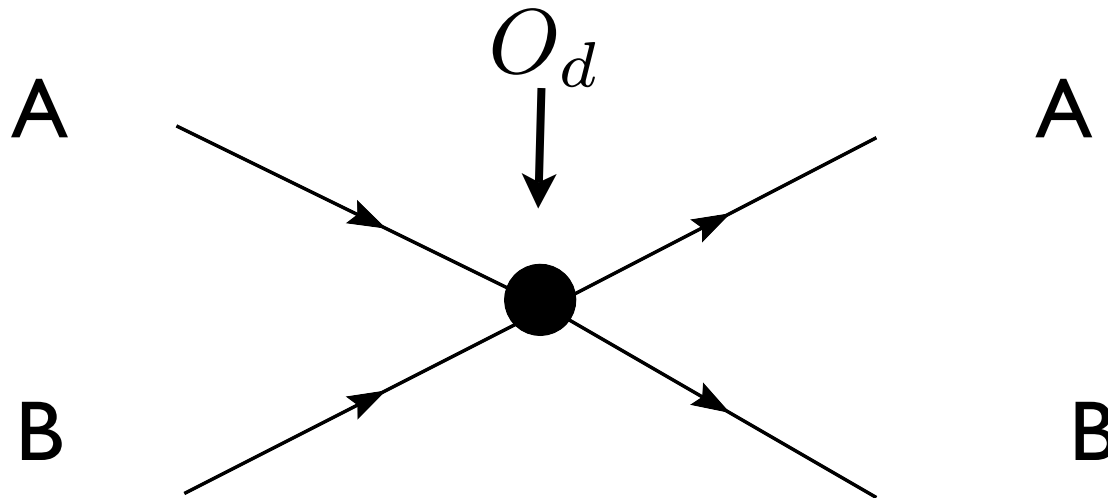
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EFT Generalities

$$L = L_0 + \sum_n C_n \frac{O_d}{\Lambda^{d-4}}$$

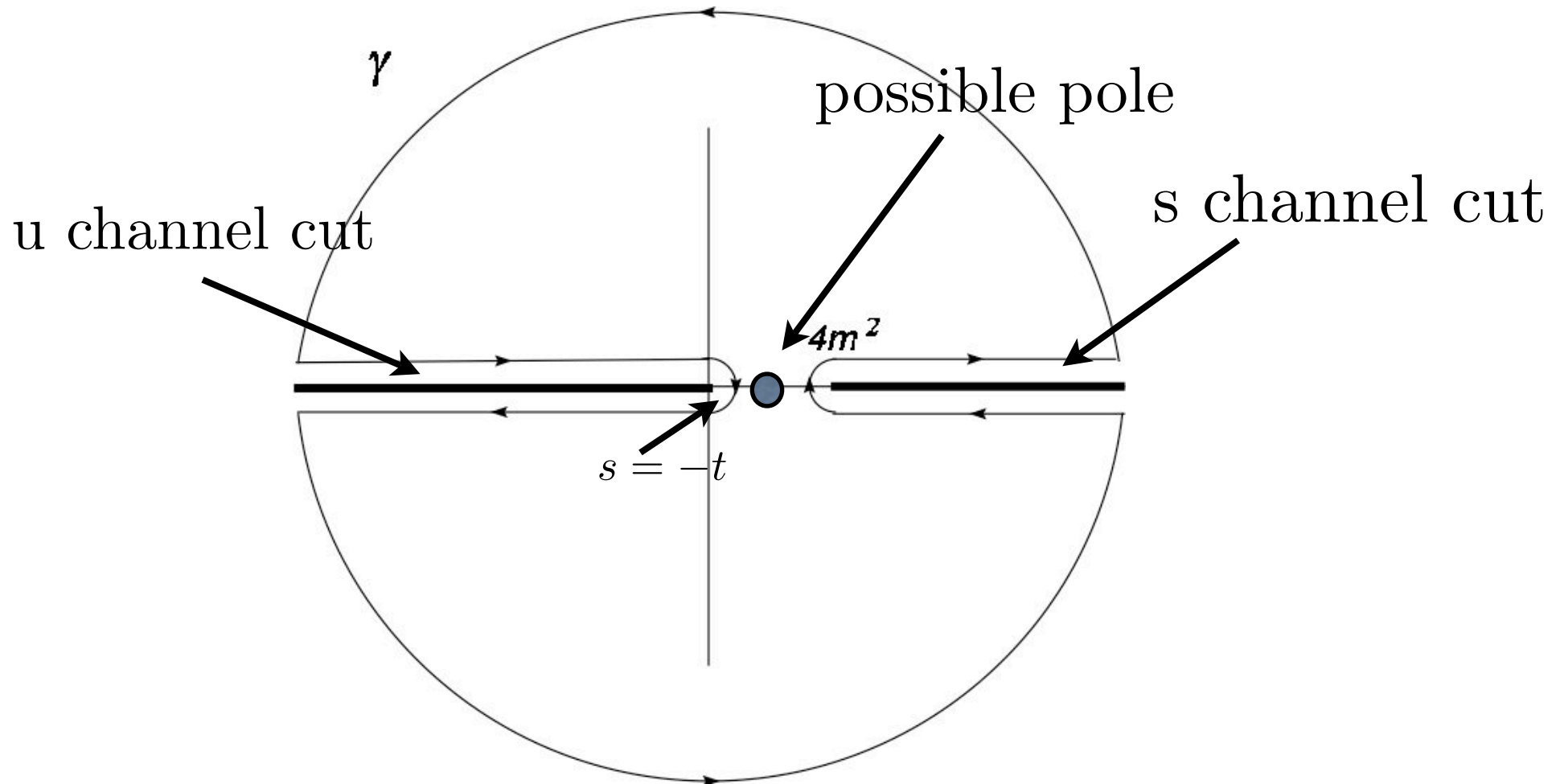
- C's are fixed by a matching procedure
- Can they take on arbitrary values?
- Naturalness: Given scale of new physics, C's should be order one.

Consider an elastic scattering process to which the operator of interest contributes
Defining the s-channel as:



$$iM \propto C_o \frac{F(s, t, u)}{\Lambda^a}$$

Fixed t : Dispersion Relation



- Assuming cut structure dictated by unitarity

Assumes unitarity and Lorentz invariance at all scales

$$\frac{\partial^2 T(s, t)}{\partial s^2} = \frac{2}{\pi} \int_t^\infty \frac{dx \operatorname{Im} T(-x + i\epsilon, t)}{(x + s)^3} + \frac{2}{\pi} \int_{4M_\pi^2}^\infty \frac{dx \operatorname{Im} T(x + i\epsilon, t)}{(x - s)^3} \\ + \sum_{s_0^i} \frac{\operatorname{res}(s_0^i)}{(s - s_0^i)^2}$$

- Twice subtracted for convergence at infinity.
- Froissart Bound follows from unitarity

$$\lim_{s \rightarrow \infty} \sigma(s) < s \ln^2 s$$

Also no long range forces

$$\frac{\partial^2 T(s, t)}{\partial s^2} = \frac{2}{\pi} \int_t^\infty \frac{dx \operatorname{Im} T(-x + i\epsilon, t)}{(x + s)^3} + \frac{2}{\pi} \int_{4M_\pi^2}^\infty \frac{dx \operatorname{Im} T(x + i\epsilon, t)}{(x - s)^3} + \sum_{s_0^i} \frac{\operatorname{res}(s_0^i)}{(s - s_0^i)^2}$$

If residue contribution is pos. def. and if we choose $s > 4m^2$
 $t = 0$

Then RHS is positive definite:

$$\text{LHS} = iM \propto C_o \frac{F(s, t, u)}{\Lambda^a} + \text{low energy known physics cont.}$$

Leads to a bound on the coupling: only
 assuming: Unitarity L.I. and Analyticity

Not all operators are boundable: Naively seems to need AT
LEAST two derivatives.

Gauged Non-linear sigma model (Heavy Higgs)

$$L = -\frac{v^2}{4} \text{Tr}(V_\mu V^\mu) + l_1 (\text{Tr}(V_\mu V^\mu))^2 + l_2 (\text{Tr}(V_\mu V^\mu))^2 + \dots$$

$$V_\mu = (D_\mu U)U^\dagger$$

Other ops well constrained by EWPO's

Consider pi-pi scattering: $(\pi_0\pi_0, \pi^+\pi_0)$

$$iM = \text{[Loop Diagram]} + \text{[Cross Diagram]} + \text{Tree level Poles}$$

l_1, l_2

Poles are not an obstruction: $\frac{(s, m^2, s^2/m^2)}{s - m^2}$

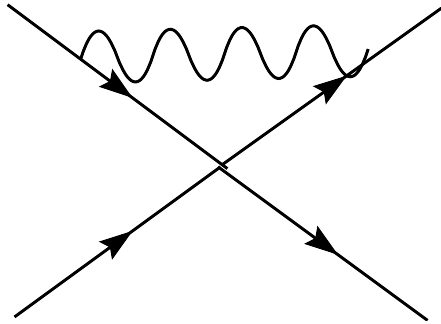
General
structure:

$$\frac{\partial^2 T(s, t)}{\partial s^2} = \frac{2}{\pi} \int_t^\infty \frac{dx \operatorname{Im} T(-x + i\epsilon, t)}{(x + s)^3} + \frac{2}{\pi} \int_{4M_\pi^2}^\infty \frac{dx \operatorname{Im} T(x + i\epsilon, t)}{(x - s)^3} + \sum_{s_0^i} \frac{\operatorname{res}(s_0^i)}{(s - s_0^i)^2}$$

Pole contributions cancel:
Would not be true for s^3
terms. Which can not be
there.

Proof: higgs unitarizes the theory and its
contributions go like s^2

Power counting:



$$\frac{\partial^2 T}{\partial s^2} \propto \frac{g^2}{16\pi^2} \left(\frac{1}{v^4}, \frac{1}{sv^2} \right)$$

Compare to the contribution
from our operators

$$\frac{l_i}{v^4}$$
$$l_i \propto \frac{1}{16\pi^2} \quad (NDA)$$

Working below
threshold:

$$s < 4m^2 \propto g^2 v^2$$

To get bounds would need to include
EW corrections

Working Above Threshold: Pion Loops dominate

$$\frac{d^2\hat{T}(s)}{ds^2} = 2! \int_{4m^2}^{\infty} \frac{dx}{\pi} \sqrt{x(x-4m^2)} \times \left(\frac{\sigma(x)}{(x-s)^3} + \frac{\sigma_u(x)}{(x-4m^2+s)^3} \right)$$

S-channel contribution no longer positive definite

However, for $s < 4\pi v^2$ we can calculate using the EFT

At $t=0$: split integral up into two intervals

$$4m^2 < s < kv^2 \qquad kv^2 < s < \infty$$

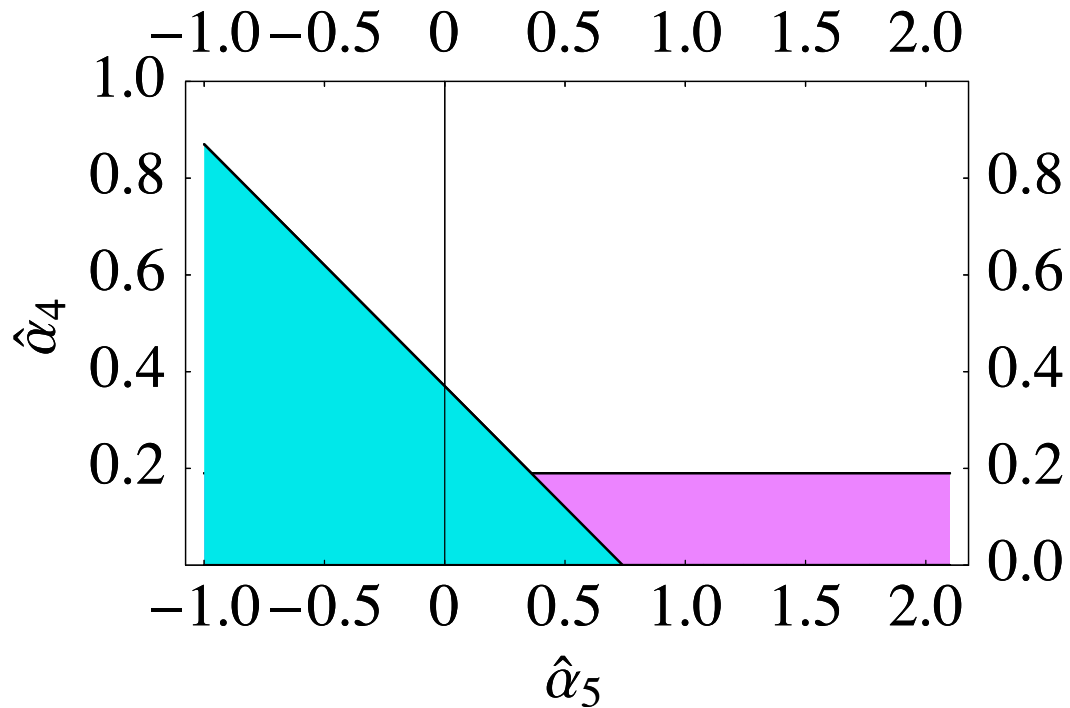
Take $s \sim v^2$ and use **Eq.Thm**

S matrix elements of long. GB's can be reproduced by S matrix elements of GB's up to corrections of order m^2/v^2 .

Use ET to calculate LHS and the first integral on the RHS

$$\text{Re} \left(\frac{d^2 \hat{T}(s)}{ds^2} \right) = 2! \left(\int_{4m^2}^{kv^2} + \int_{kv^2}^{\infty} \right) \frac{dx}{\pi} \sqrt{x(x-4m^2)} \times \left(\frac{\sigma(x)}{(x-s)^3} + \frac{\sigma_u(x)}{(x-4m^2+s)^3} \right)$$

Since RHS grows with s (near threshold), choose s to be as large as possible within **errors (chPT should still converge)**



Dominant source of errors on
bound come from LHS

$$\delta l \sim \frac{g^2}{k} \sim \%20$$

- Note: can't use this technique (reliably in chiral Lagrangian (QCD) since in that case

$$f_\pi < m_\pi$$

So there is **very little window to work with**

$$4m_\pi^2 < s < 4\pi f_\pi$$

Suppose bounds were
violated:

1) Underlying theory does not obey usual axioms of QFT.
NOT string theory (at least in form we build models with).

2) There exists light resonances below $4\pi v$

e.g. 5-d theory in Ads dual to large N

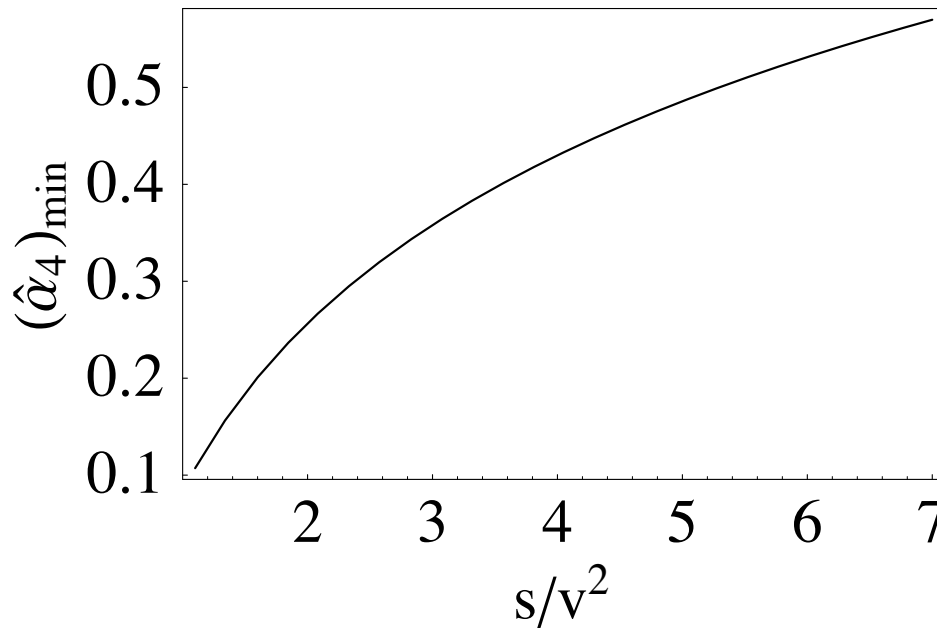
ERRORS to RHS

$$\delta_{ew}T \propto O\left(\frac{g^2 s}{(4\pi v)^2} \ln(s/\mu^2)\right)$$

EW loops

$$\delta_{\chi}T \propto O\left(\frac{s^3}{(4\pi)^4 v^6} \ln(s/\mu^2)\right).$$

Chiral loops+CT's



Choose s large while keeping these errors order %20