

## Measurement of $d\sigma/dp_T (Z / \gamma^* \rightarrow e^+ e^-)$ at D0

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### **Outline of the Talk**



- Motivation
- D0 detector
- $Z / \gamma^* \rightarrow e^+ e^-$  event selection
- Single electron efficiencies
- Monte Carlo simulation
- Background
- Unfolding of detector smearing
- Systematic uncertainty estimate
- Preliminary result and conclusions



### **Motivation**



- Z boson production: QCD  $q\bar{q} \rightarrow Z$
- $p_T$  of Z boson : initial state gluon radiation of the colliding partons



- High p<sub>T</sub>, perturbative QCD

$$\frac{d^2\sigma}{dp_T^2 dy} \propto \frac{\alpha_w \alpha_s}{p_T^2} \ln\left(\frac{Q^2}{p_T^2}\right) \to \infty \ as \ p_T^2 \to 0$$

Low p<sub>T</sub>, soft and collinear gluon resummation



## **Motivation**



- **W(b,Q)** resumming the perturbative series  $p_T^{-2} \times [1 \text{ or } \ln(Q^2 / p_T^2)]$  to all orders in  $\alpha_{s}$
- Y(b,Q) diff between the fix order perturbative result and the singular part
- To extend to pT->0 region, where non-perturbative QCD ٠ dominate the production  $d\sigma/dp_T^*BR(Z \rightarrow ee) (pb/GeV)$

 $W(b,Q) \to W(b_*,Q)e^{-S_{NP}(b,Q)} \quad b_* \equiv b / \sqrt{1 + (b / b_{max})^2} \le b_{max}$ 

- **S<sub>NP</sub>: non-perturbative Sudakov form factor**
- Lindasky-Yuan form:  $S_{NP}^{LY}(b,Q^2) = g_1 b^2 + g_2 b^2 \ln(\frac{Q^2}{2Q^2}) + g_1 g_3 b \ln(100x_i x_j)$
- Provide sensitive test of the weak boson production formalism
- Help to reduce theory uncertainty of the precision W mass measurement



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DØ 1994-1996 ace (Ladinsky-Yuan)  $EQ4M g_1=0.11 \ GeV^2 g_2=0.58 \ GeV^2$  $g_2=-1.5 \ GeV^1$ 

g\_-space (Ellis-Veseli) MRSR1 a=0.1 GeV<sup>2</sup> q<sub>11im</sub>=4.0 GeV b-space (Davies-Webber-Stirling)

MRSA g ,=0.15 GeV2 g ,=0.4 GeV Fixed-order ( $O(\alpha_s^2)$ )

25 30 p<sub>r</sub>(GeV)

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**Problem:** resummation can not describe  $P_T$  in semi-inclusive DIS

Solution? Two large logs. Normal resummation handles

$$\frac{1}{\sigma} \frac{d\sigma}{dp_T^2} \simeq \frac{1}{p_T^2} [A_1 \alpha_s \ln(\frac{Q^2}{p_T^2}) + A_2 \alpha_s^2 \ln^3(\frac{Q^2}{p_T^2}) + \dots + A_n \alpha_s^n \ln^{2n-1}(\frac{Q^2}{p_T^2}) + \dots]$$

## Normal resumation does not handle $\alpha_s^n \ln^n(1/x)$ terms. Mock this in nonpert part



Small x broadening



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#### **D0 detector**



- Tracking system:
- Silicon microstrip tracker
- Scintillating fiber tracker
- Super conducting solenoid
- Uranium/liquid argon calorimeter

coverage:

Central: |η|<1.1

Endcaps: 1.5<|η|<4.2

Muon detector with toroid



#### **Data Sample and Event Selection**



- Data collected in D0 Run II between Oct 2002 and Jan 2006, bad runs removed.  $\int Ldt \sim 965 \pm 58 \, pb^{-1}$
- Two high p<sub>T</sub> electrons each satisfying:
  - p<sub>T</sub>>25 GeV
  - isolated
  - in good region of calorimeter
  - matches typical electron shower shape
  - have a track match
- Invariant mass between [70,110] GeV
- At least one electron fires a single electron trigger

CCCC	CCEC	ECEC	All
23957	30116	9583	63656

TABLE I: Results of event selection



#### **Single Electron Efficiencies**



• Tag and probe method

Tag	Tag electron: Passes all single electron requirements		CC	EC
Z/γ* Pass electreg		Electron identification	(99.5 ±0.1)%	(99.1 ±0.1)%
		Trigger	(98.5 ±0.1)%	(97.9 ±0.1)%
Probe electron:		Shower shape	(97.1 ±0.1)%	(96.9 ±0.1)%
Passes loose ele requirements a	s loose electron rements and	Track match	(90.5 ±0.1)%	(61.7 ±0.1)%
check for each selection requiremen		Ţ		

$$\varepsilon = \frac{N_{pass}^{probe}}{N_{All}^{probe}}$$

Their dependence on electron location and  $P_T$  are modeled in Monte carlo



#### **Monte Carlo simulation**



- Event generator: Resbos(Ladinsky+Yuan)+PHOTOS
- Resbos(hep-ph/9704258):initial state gluon resummation;g1,g2,g3;
- PHOTOS(hep-ph/0506026):final state QED radiation.





#### Detector response- Parameterized Monte Carlo Simulation



Parameterized Monte Carlo simulation(PMCS): smear electron energy, η and φ position;

Merge nearby photons to electrons;

Apply single electron efficiencies and acceptance cuts.

**Energy smearing:** 

$$E' = \alpha * E_{gen} + \beta$$

$$\frac{\sigma_E}{E} = \sqrt{C^2 + \frac{S^2}{E} + \frac{N^2}{E^2}}$$

$$E_{smear} = E' + x * \sigma_E$$

Parameters( $\alpha,\beta,C$ ) are determined from Z data by comparing invariant mass distribution.



### Backgrounds



- QCD backgrounds:
- Main source of backgrounds is QCD background with jets faking electrons. (Di-jet, EM+jet)
- Selection: inverse shower shape cut on the electrons.
- Non-QCD backgrounds: PYTHIA as generator, PMCS as detector smearing



	ε*A	σ*Br(pb)	N(0.968fb <sup>-1</sup> )
$Z \rightarrow \tau \tau \rightarrow ee$	0.0025	7	16.7
WW	0.059	0.1	6.2
WZ	0.17	0.009	15.5
Wγ	0.005	12.2	61
the non-QCD backgrounds are negligible			



#### **Background fraction and p<sub>T</sub>-shape**



• Minimum  $\chi^2$  fit: invariant mass of the Z candidates as linear sum of those of signal(MC) and background to get the background fraction



**QCD** background shape:

Region	$\chi^2/\mathrm{ndf}$	background fraction
CCCC	64/80	$1.30{\pm}0.14\%$
CCEC	106/80	$8.55 {\pm} 0.26\%$
ECEC	80/80	$4.71 \pm 0.30\%$
All	111/ 80	$4.70 {\pm} 0.13\%$





#### **Data/MC Comparisons**



# Put measured efficiency and tuned parameters in the PMCS





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## Efficiency\*Acceptance Z p<sub>T</sub> dependence



Data selection affect Z pT spectrum;

Jet activity affect the isolation of the electrons

**Dependence studied from Monte Carlo** 



Efficiency Z pT dependence: full Monte Carlo(GEANT-based); Signal overlaid on underlying events

smeared  $Z p_T(accep cuts, eff cuts)$ 

smeared  $Z p_T$  (accep cuts)





## Efficiency\*Acceptance Z p<sub>T</sub> dependence



Acceptance (electron pT, position and invariant mass) Z pT dependence: PMCS

smeared  $Z p_T$  distribution with accep cut(s)

smeared  $Z p_T$  distribution with no cuts



**Efficiency**\*Acceptance p<sub>T</sub> dependence

Measured Z  $\ensuremath{p_{\mathrm{T}}}$  spectrum need to correct for this dependence



## **Unfolding of Detector Smearing**



- *RUN*(*R*egularized *Un*folding)(hepex/0208022) by V.Blobel
- Maximum likelihood fit(data and trial result with parameters) with regularization term to smooth oscillations due to large correlations between adjacent bins.
- ► Input:
- $\checkmark$  Ntuples of Z pT data
- ✓ Ntuples of Z pT MC: generator, smeared(contains smearing information)
- ✓ Ntuples of pT of QCD background





#### **Systematic Uncertainty**

Energy resolution parameters:constant term and sampling term;Energy scale parameter;PDFs, background, unfolding method(input,parameters), data/full MC discrepancy of efficiency pT dependence





# Preliminary result, conclusions and further

- Preliminary result for 0.965 fb<sup>-1</sup> data
- All y region result consistent with CSS resummation calculation
- Next to perform tuning of Resbos g2
- Small x broadening test in progress for |y|>2







## Backup slides



#### **Tevatron for Run II**



#### New 120 GeV Main Injector CM increased form 1.8 TeV to 1.96 TeV



#### Collider Run II Peak Luminosity



#### Collider Run II Integrated Luminosity



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#### **D0 calorimeter**



#### Compensating sampling calorimeter:uranium/liquid argon



**D0 Run I** C = 1.0%**CC:**  $N = 0.38 \, GeV$  Eta coverage: Central: |η|<1.1 Endcaps: 1.5<|η|<4.2

 $(\sigma/E)^2 = C^2 + S^2/E + N^2/E^2$ 

Interpretation:

C = contant term (due to non-uniformities in the response; "intercalibration errors")

S = sampling term (due to sampling fluctuations; "intrinsic performance of the CAL")

N = noise term (due to "noise", e.g. from Uranium, the readout electronics, and, in collision data, the "underlying event")



### **Typical parameterization**





#### **Tracking matching efficiency for EC electrons versus vertex z position**



### **CC Energy resolutions**



Calorimeter suffers from the materials added in D0 Run II. Sampling is no longer constant, but depend on position.

**Constant Term tuned from Z(2.8%)** 



#### **Correct Z pT spectrum**



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