

Observation of Exclusive Hadronic Processes at 10.6 GeV: $e^+e^- \rightarrow \rho^0\rho^0, \phi\rho^0, \gamma\eta, \gamma\eta', \phi\eta$

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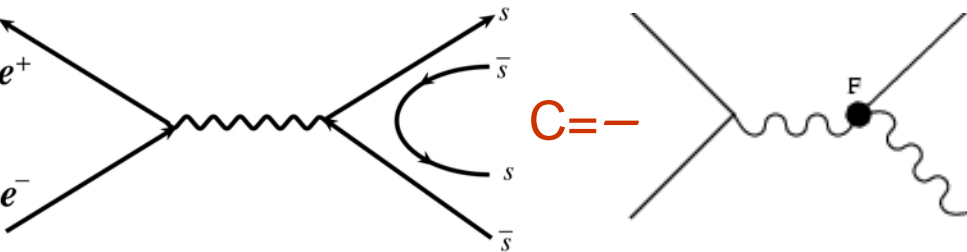
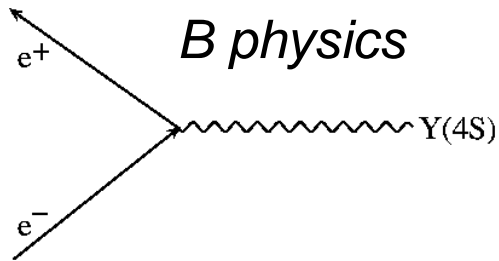


Outline

- *Introduction*
- *First Observation of **positive C parity** hadronic final states in the reactions $e^+e^- \rightarrow \rho^0\rho^0, \phi\rho^0$*
- *Observation of $e^+e^- \rightarrow \gamma\eta, \gamma\eta'$ ($C=-1$) and the time-like transition form factor at 10.6 GeV*
- *Observation of $e^+e^- \rightarrow \phi\eta$ ($C=-1$ final state)*
- *Summary and Outlook*

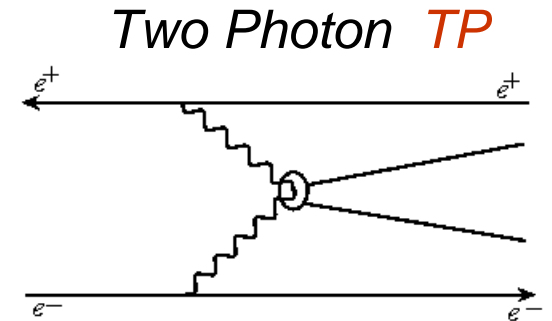
Introduction—possible processes at B factories

One photon related process

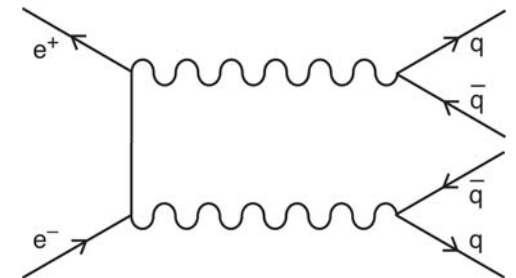


A test ground for QCD

Two photon interaction processes



$C=+$



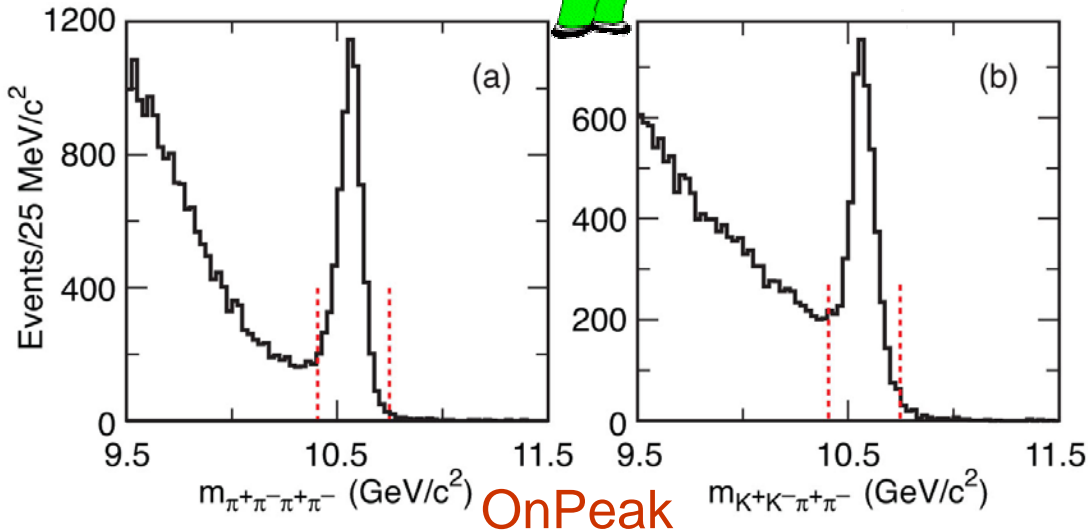
Two-Virtual-Photon-Annihilation, TVPA; new observation at BaBar

Today's topic

Introduction

- B factories are desired to study the following process:
 $e^+e^- \rightarrow \gamma^* \rightarrow Y(4S) \rightarrow B \bar{B}$
- The high integrated luminosity at B factories also provides a **new window** on other e^+e^- physics :
 - I. Two-virtual-photon-annihilation (TVPA) yielding hadronic final states*
 - Different from well-known TP
 - Speculated to be significant but not observed for double charmonium production
 - First time observation
 - II. one photon to **exclusive** hadronic final states*
 - $e^+e^- \rightarrow$ hadronic states are studied everywhere, exclusive is special
cross section very small @10.6 GeV, test s dependence
 - III. Interference between these processes*
- Very large integrated luminosity at interesting s values!
~205 fb⁻¹ @10.58 GeV (OnPeak), 20 fb⁻¹ @10.54 GeV (OffPeak)
Can explore new opportunities; much more to be analyzed!

Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ ($C=+1$)

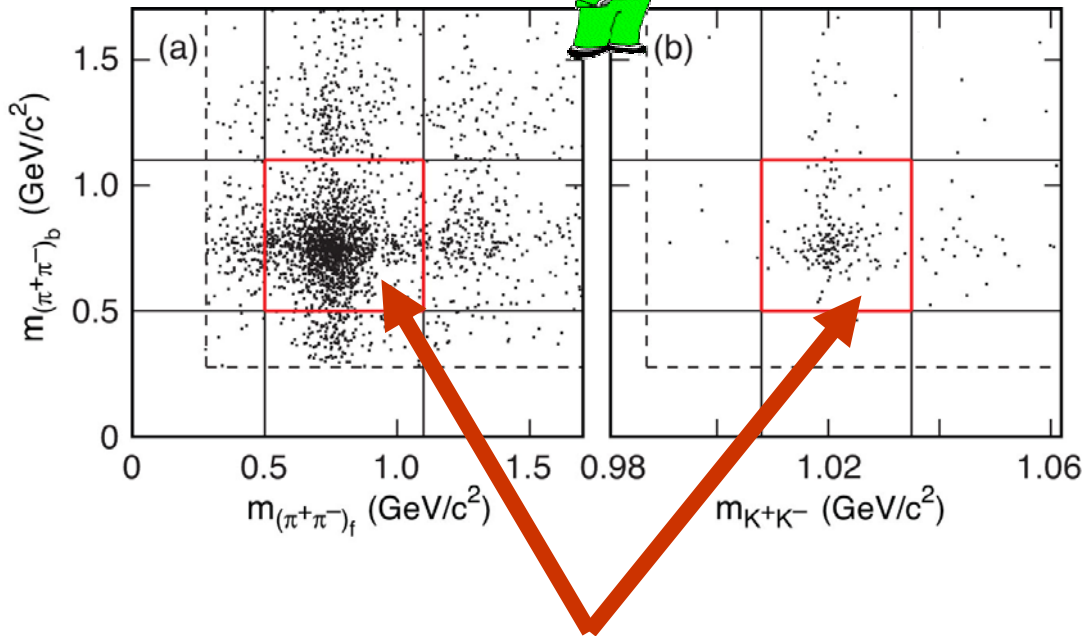
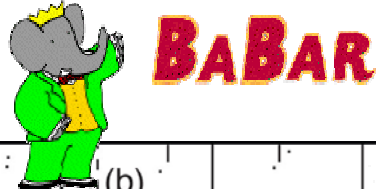


- 4 well-reconstructed charged tracks ($K^+K^-\pi^+\pi^-/\pi^+\pi^-\pi^+\pi^-$)
- identified kaons or pions
- Vertex Constraint $p(\chi^2) > 0.1\%$
- $p_K > 800$ MeV/c, $p_\pi > 600$ MeV/c to reduce QED background

Select events of invariant mass of 4 track system ($K^+K^-\pi^+\pi^-/\pi^+\pi^-\pi^+\pi^-$) within 170 MeV of c.m. energy.

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Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ ($C=+1$)



Scatter plots of :

m_{KK} vs $m_{\pi\pi}$ ($KK\pi\pi$)

$m_{(\pi\pi)_f}$ vs $m_{(\pi\pi)_b}$ ($\pi\pi\pi\pi$)

Only *one entry* in the $\pi^+\pi^-\pi^+\pi^-$ scatter plot out of two possible combinations
(no ambiguity)

Use *binned log-likelihood fit* over 9 tiles to extract signal

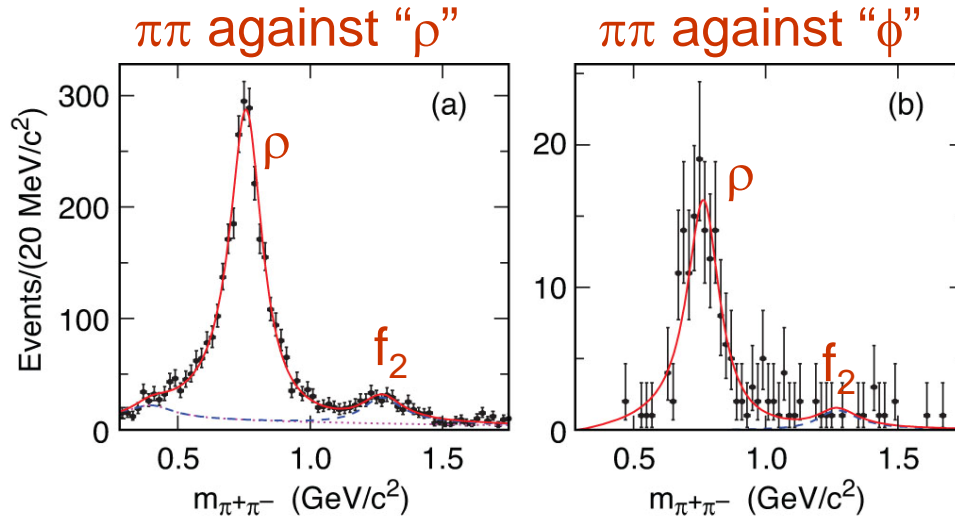
Define *signal tiles*:

$$0.5 < m_{\pi\pi} < 1.1 \text{ GeV}/c^2 \quad \& \quad 1.008 < m_{KK} < 1.035 \text{ GeV}/c^2$$

“ ρ ”
“ ϕ ”

Use adjacent tiles to subtract background

Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ (C=+1)



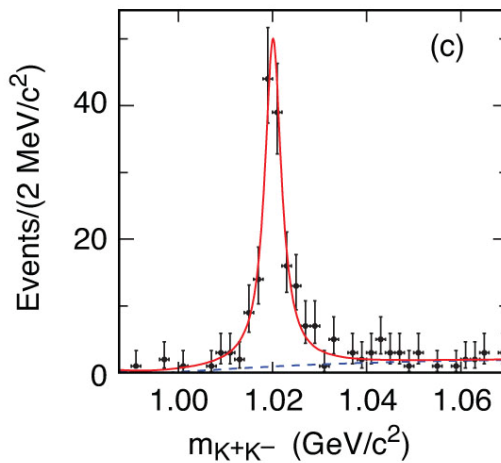
*Mass projections for
selected events*

Clean ϕ and ρ^0 signals

*Relativistic Breit-Wigner
lineshapes used in the fit*



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KK against “ ρ ”

Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ ($C=+1$)

The extracted signals:

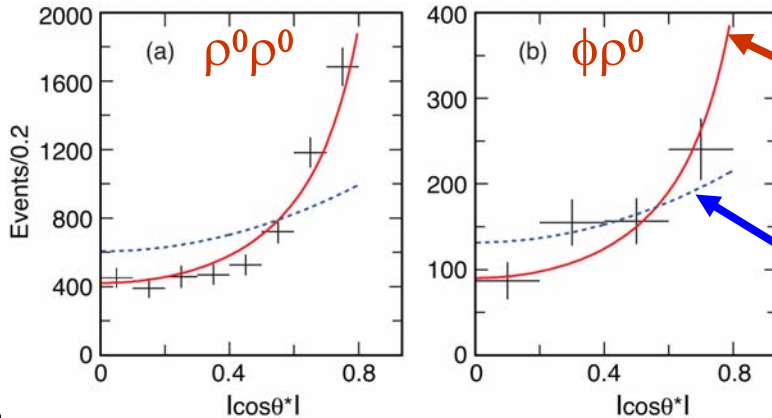
	On+OffPeak		OnPeak Yield	OffPeak Yield	OffPeak expected (continuum)
	Yield	significance			
$\rho\rho$	1243 ± 43	$\gg 5\sigma$	1138 ± 42	104 ± 14	112 ± 4
$\phi\rho$	147 ± 13	$\gg 5\sigma$	135 ± 13	14 ± 4	13 ± 1

 consistent

$Y(4S)$ production does not contribute because of C parity violation.
The yields are consistent with **continuum production**
Signals large enough that can investigate the angular distributions

Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ ($C=+1$)

Efficiency corrected distributions



Production angle θ^* – polar angle of ϕ or ρ_{forward} in CM

Expectation for TVPA:

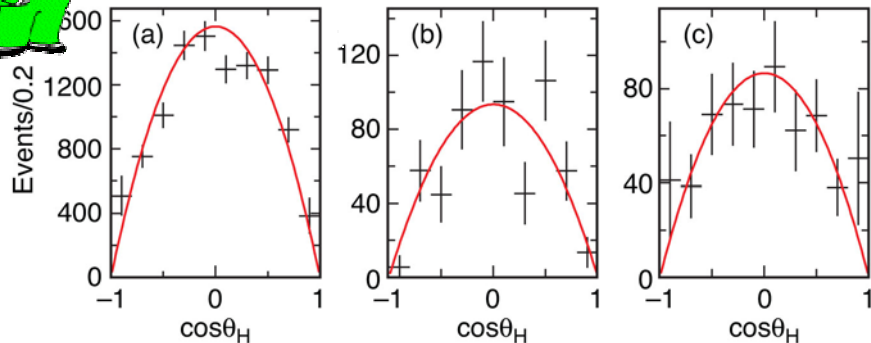
$$\frac{d\sigma}{d\cos\theta^*} \propto \frac{1 + \cos^2\theta^*}{k^2 - \cos^2\theta^*} \quad \text{with } k^2 \sim 1$$

where k depends on the vector meson masses and \sqrt{s} :

$1 + \cos^2\theta^*$ for one photon annihilation (for comparison)



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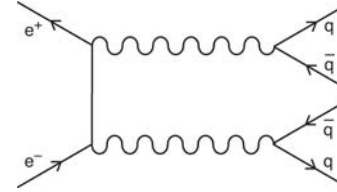


$\sin^2\theta_H$ distributions

All angular distributions are consistent with TVPA expectation

Observation of $e^+e^- \rightarrow \rho^0\rho^0/\phi\rho^0$ ($C=+1$)

These are the *first observations* of $e^+e^- \rightarrow C=+1$ exclusive hadronic final states
Arguments to support TVPA production:



- C parity
- Consistency of the yield between OnPeak and OffPeak
- Observed angular distributions

Define $1.008 < m_\phi < 1.035$ GeV and $0.5 < m_\rho < 1.1$ GeV (reduce uncertainty);
the cross sections for these mass regions and $|\cos\theta^*| < 0.8$ are:

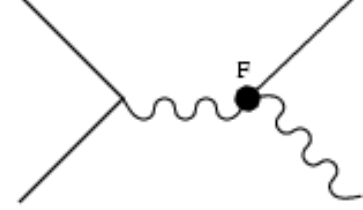
$$\sigma(\rho\rho) = 20.7 \pm 0.7(\text{stat}) \pm 2.7(\text{syst}) \text{ fb} \quad 21.4 \pm 0.7 / 17.7 \pm 0.6 \text{ fb (Theory)}$$

$$\sigma(\phi\rho) = 5.7 \pm 0.5(\text{stat}) \pm 0.8(\text{syst}) \text{ fb} \quad 6.0 \pm 0.1 / 5.6 \pm 0.2 \text{ fb (Theory)}$$

[note: $\sigma(e^+e^- \rightarrow \text{hadrons at } E_{\text{cm}}=10 \text{ GeV}) \sim 3 \text{ nb}$]

The corresponding contribution to muon g-2 from TVPA is small; this removes a possible uncertainty source in the g-2 calculation

$e^+e^- \rightarrow \eta\gamma, e^+e^- \rightarrow \eta'\gamma$



- The cross section of $e^+e^- \rightarrow \gamma^* \rightarrow P\gamma$ ($P = \eta$ or η'):

$$\sigma(e^+e^- \rightarrow P\gamma) = \frac{2}{3} \pi^2 \alpha^3 |F_P(q^2)|^2$$



transition form factor

$F_P(q^2)$ can be calculated using pQCD:

$$-q^2 F_P(q^2) = \sqrt{2} f_P \left(1 - \frac{5}{3} \frac{\alpha_s(q^2)}{\pi} \right) \text{ for } q^2 \gg 1 \text{ GeV}^2$$



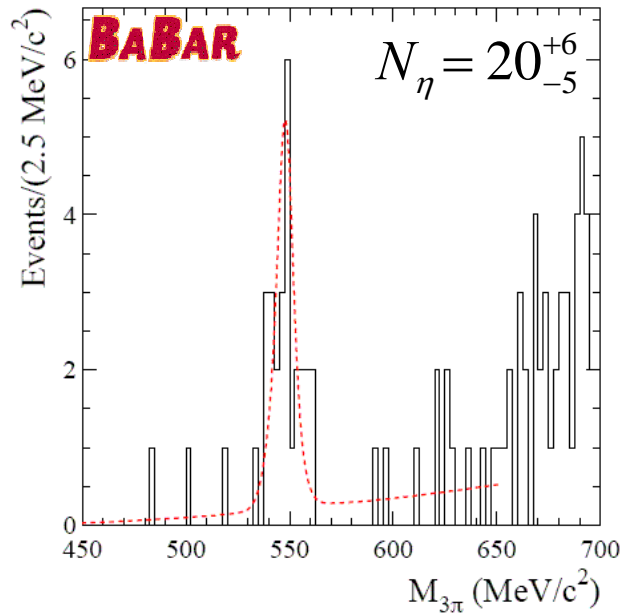
effective decay constant

- Measurement of form factors at high q^2 is needed to test phenomenological models.
 - Previous measurements from $\gamma^*\gamma^* \rightarrow \eta$ (') measure up to $q^2 \approx 20 \text{ GeV}^2$.
 - We measure the transition form factor at $q^2 = 112 \text{ GeV}^2$, i.e. much closer to the asymptotic region.

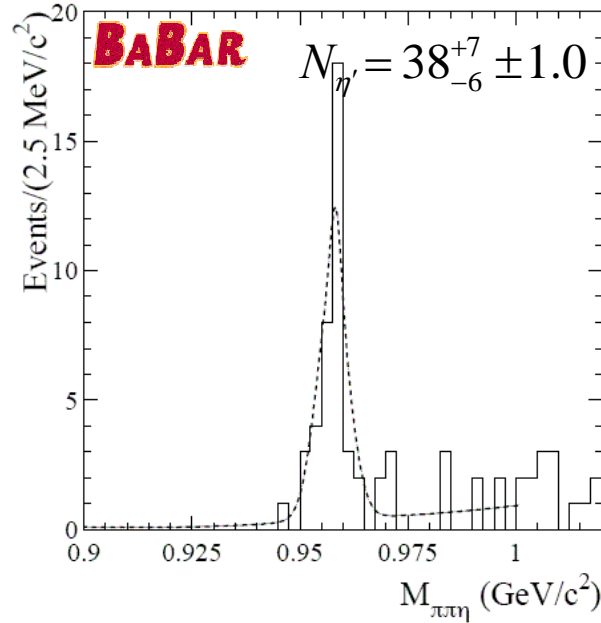
$e^+e^- \rightarrow \eta\gamma, e^+e^- \rightarrow \eta'\gamma$

- Very rare but clean two-body final states:

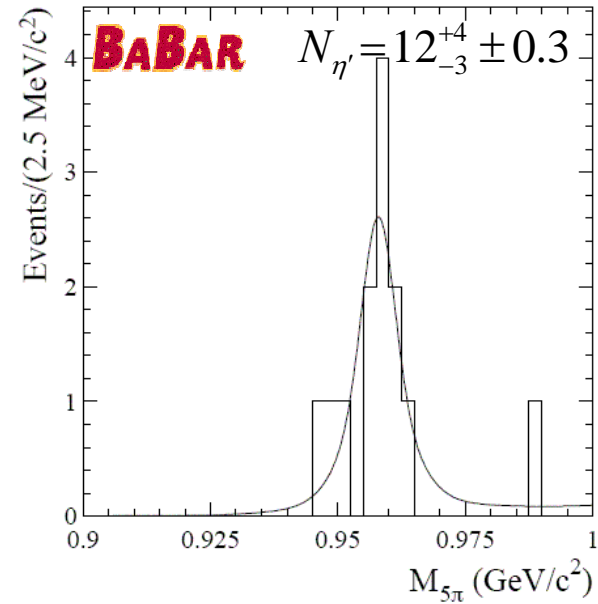
$e^+e^- \rightarrow \eta\gamma \rightarrow 3\pi\gamma$



$e^+e^- \rightarrow \eta'\gamma \rightarrow 2\pi 3\pi\gamma$



$e^+e^- \rightarrow \eta'\gamma \rightarrow 5\pi\gamma$



$$\sigma(e^+e^- \rightarrow \eta\gamma) = 4.5^{+1.2}_{-1.1} \pm 0.3 \text{ fb}$$

$$\sigma(e^+e^- \rightarrow \eta'\gamma) = 5.4 \pm 0.8 \pm 0.3 \text{ fb}$$

PRD 74:012002 (2006)

$e^+e^- \rightarrow \eta\gamma, e^+e^- \rightarrow \eta'\gamma$

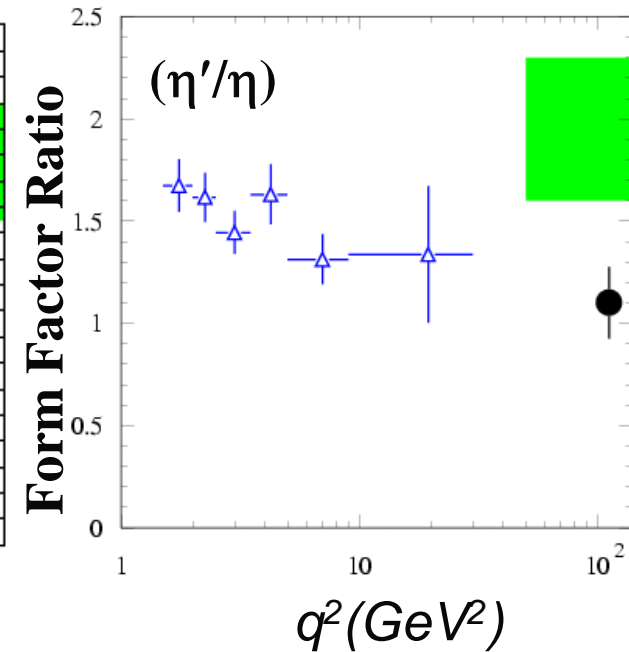
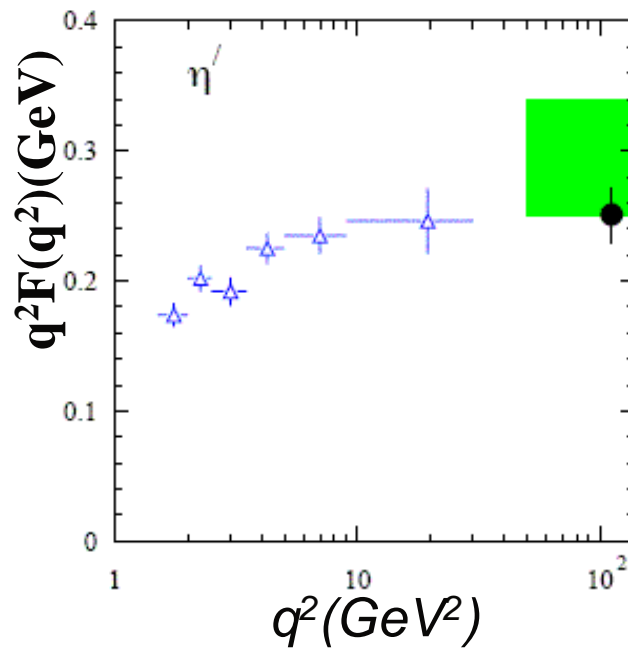
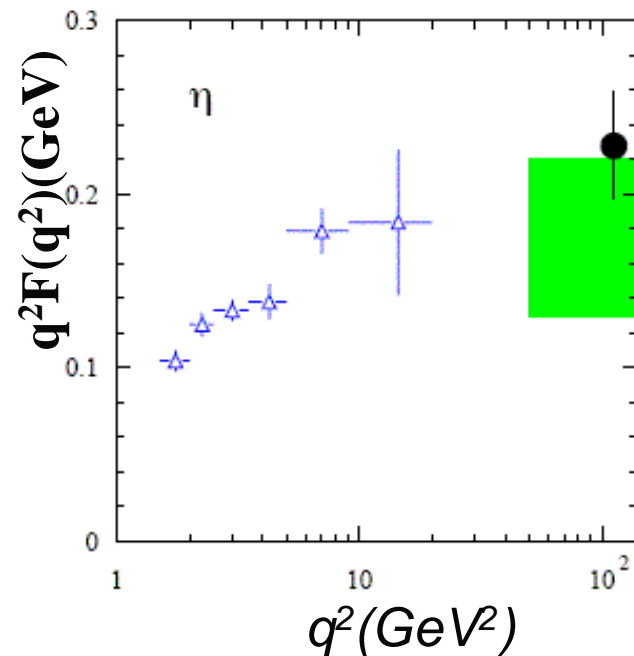
$q^2=112 \text{ GeV}^2$	<i>BaBar Data</i>	<i>Predictions</i>
$q^2F(\eta')$	$0.251 \pm 0.019 \pm 0.008$	$0.25 - 0.34$
$q^2F(\eta)$	$0.229 \pm 0.030 \pm 0.008$	$0.11 - 0.22$
<i>Ratio</i> (η'/η)	1.10 ± 0.17	$1.56 - 2.27$

The Transition Form Factors

• **BaBar**

△ **CLEO**

Theoretical predictions



Observation of $e^+e^- \rightarrow \phi\eta$ at ~ 10.6 GeV

To be submitted to PRD

- $e^+e^- \rightarrow \phi\eta$ is analogous to $e^+e^- \rightarrow J/\psi\eta_c$ in the s quark sector
need relativistic treatment, η --not pure s -sbar pair
- *Provide information on s -dependence by combining with a CLEO measurement at lower energy*

Basic requirements (similar to $\phi\rho$:

--2 well-reconstructed tracks

-- at least 2 good photons ($E_\gamma > 500$ MeV)

--identified kaons

--Vertex constraint, $p(\chi^2) > 0.001$

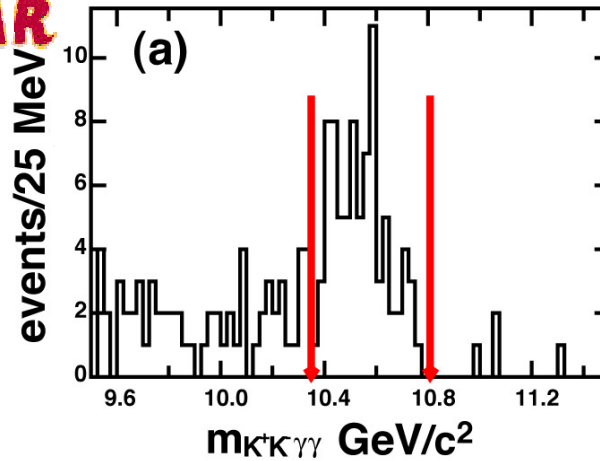
-- $m_{KK} < 1.1$ GeV/ c^2 , $0.4 < m_{\gamma\gamma} < 0.6$ GeV/ c^2

*--Invariant mass of 4 particle ($K^+K^-\gamma\gamma$)
within 230 MeV of c.m. energy.*

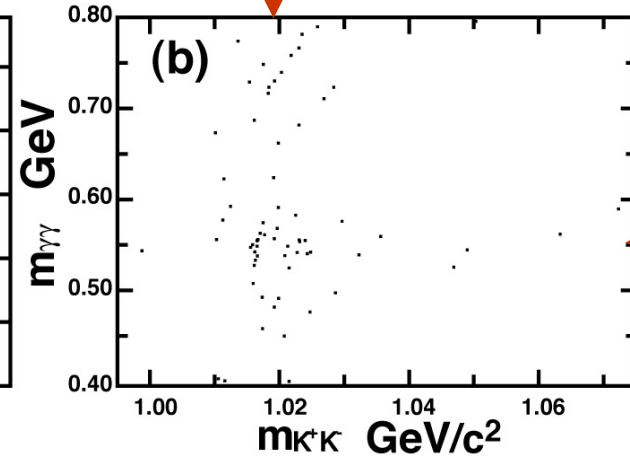
Observation of $e^+e^- \rightarrow \phi\eta$ at ~ 10.6 GeV



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Y(4S) data



Y(4S)+Off peak data

We see $\phi\eta$ correlation. Define $1.008 < m_\phi < 1.035$ GeV/c²

use a two-dimensional log-likelihood fit to extract signals

P -wave relativistic Breit-Wigner for ϕ

Gaussian resolution function for η

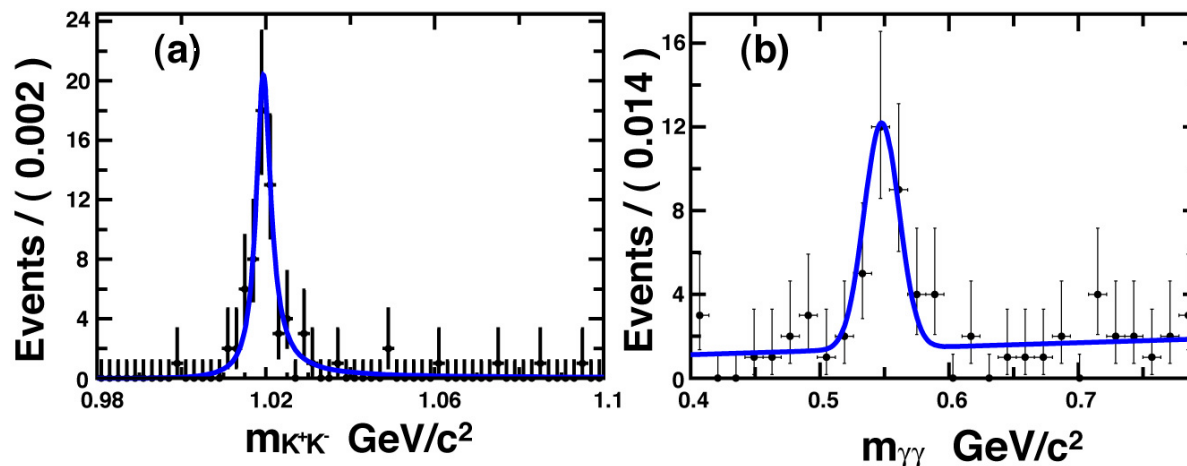
fix the mass and width of the ϕ and the mass of the η to PDG 2004 values

fix the r.m.s value of the η Gaussian to the MC resolution value (13.6 MeV/c²).

Observation of $e^+e^- \rightarrow \phi\eta$ at ~ 10.6 GeV



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Mass projections

Yield in the defined mass window:

On+Off: 24 ± 5 On: 20 ± 5 Off: 3 ± 2

Significance: 6.5σ

U.L. for $Y(4S)$ decay @90% CL based on -10 ± 21 events: 2.5×10^{-6}

Observation of $e^+e^- \rightarrow \phi\eta$ at ~ 10.6 GeV

The full angular distribution, assuming one-photon production, is given by:

$$\frac{dN}{d \cos \theta^* d \cos \theta_\phi d\varphi_\phi} \propto \sin^2 \theta_\phi (1 + \cos^2 \theta^* + \cos 2\varphi_\phi \sin^2 \theta^*)$$

Projections:

θ^* : polar angle in c.m.

$$\frac{dN}{d \cos \theta^*} \propto (1 + \cos^2 \theta^*)$$

θ_ϕ : ϕ helicity angle

$$\frac{dN}{d \cos \theta_\phi} \propto \sin^2 \theta_\phi$$

φ_ϕ : azimuthal angle of the decay plane ϕ w.r.t. production plane normal in the c.m. frame

$$\frac{dN}{d\varphi_\phi} \propto (2 + \cos 2\varphi_\phi)$$

Distributions in data are consistent with predictions with limited statistics

We *re-weight* MC angular distributions as above for efficiency estimation.

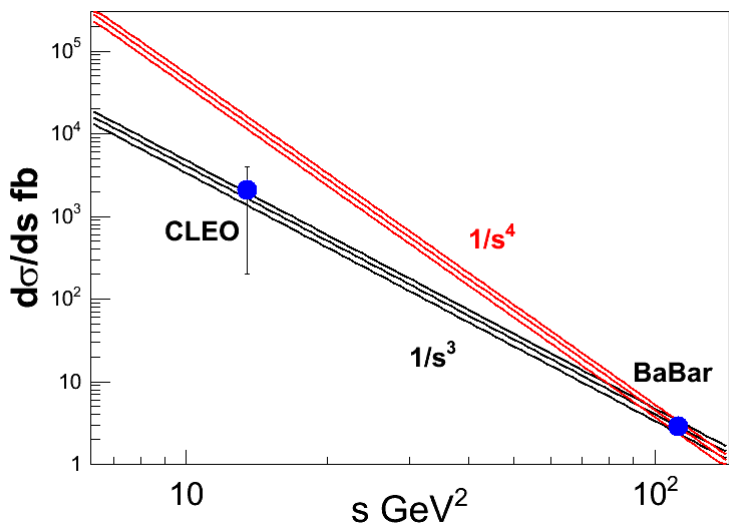
Observation of $e^+e^- \rightarrow \phi\eta$ at ~ 10.6 GeV

Define $1.008 < m_\phi < 1.035$ GeV/c² and $0.5 < m_\rho < 1.1$ GeV/c². After radiative corrections, the restricted cross section ($|\cos\theta^*| < 0.8$) are measured as:

$$\sigma(\phi\eta) = 2.1 \pm 0.4(\text{stat}) \pm 0.1(\text{syst}) \text{ fb}$$

Extend to $|\cos\theta^*| \leq 1.0$ by assuming $1 + \cos^2\theta^*$ distribution:

$$\sigma(\phi\eta) = 2.9 \pm 0.5(\text{stat}) \pm 0.1(\text{syst}) \text{ fb}$$



Combine with CLEO measurement at lower energy; $1/s^3$ energy dependence favored over $1/s^4$, assuming continuum production

A direct comparison with theoretical calculation is interesting. Another puzzle?

Analysis in Initial State Radiation data is ongoing, will provide more information on the s dependence

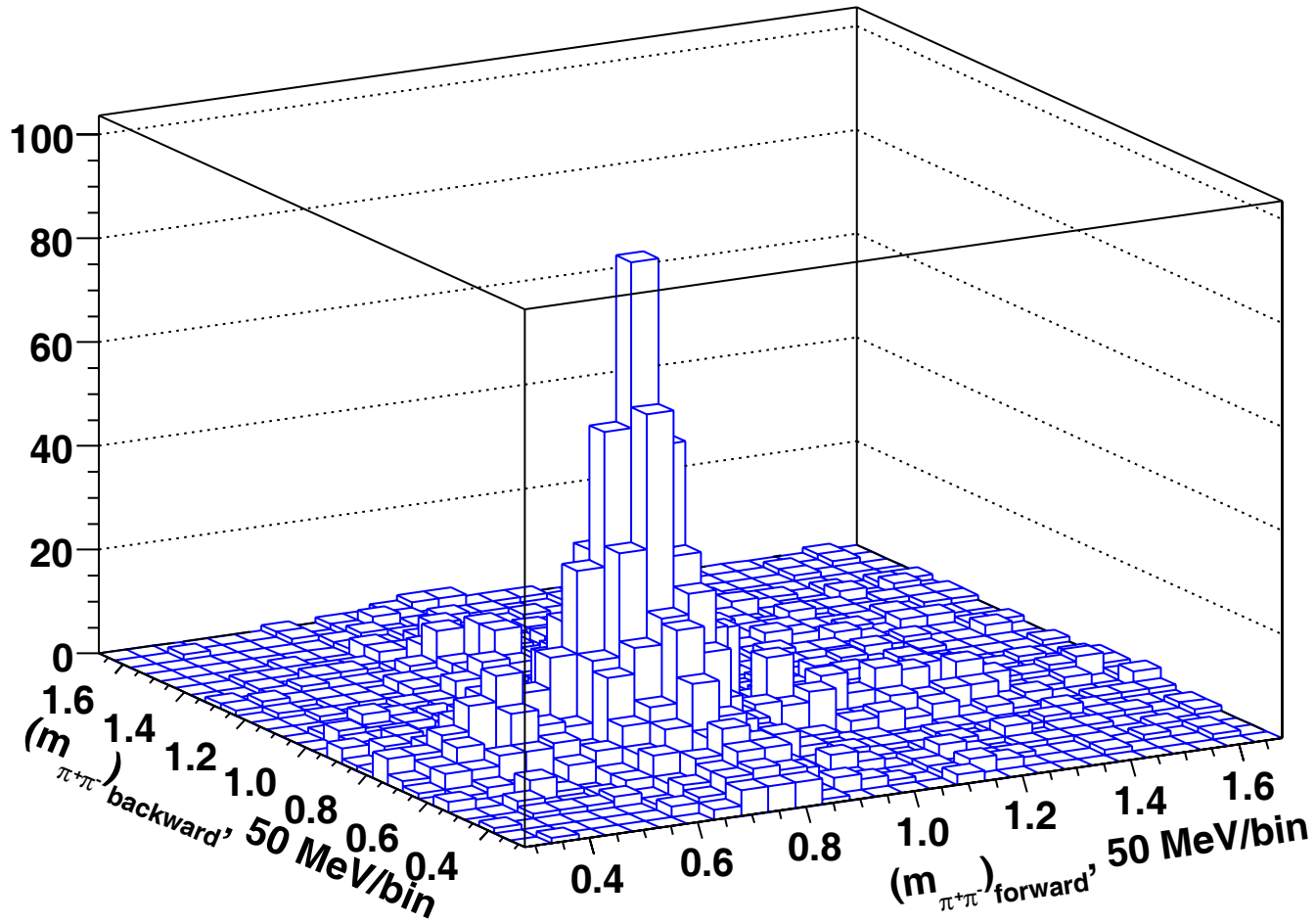
Summary and Outlook

At BaBar we have:

- Made the first observations of the TVPA hadronic processes:
 $e^+e^- \rightarrow \rho^0\rho^0$ **and** $e^+e^- \rightarrow \phi\rho^0$
- Measured the η and η' transition form factors at high q^2 ($\sim 112 \text{ GeV}^2$)
- Observed the process $e^+e^- \rightarrow \phi\eta$ and measured the cross section; this provides an interesting test of QCD prediction of the energy dependence, assuming continuum production.

We plan to use the high integrated luminosity at BaBar to investigate other interesting low-multiplicity exclusive final states in the near future!

Backup plots



Nine Tile likelihood fit Method

A binned likelihood fit based on nine-tiles is designed to extract signal yield in order to handle possible nonlinear background in $e^+e^- \rightarrow \phi\rho^0$ and $e^+e^- \rightarrow \rho^0\rho^0$ which is described in Chapter 7. The likelihood of obtaining the events observed in each tile n_i is given by

$$L = \prod_i e_i^{-m_i} \frac{m_i^{n_i}}{n_i!} \quad (2)$$

where m_i is the number of events expected for a given choice of the parameters, and n_i is the number seen. For purposes of fitting it is convenient to take the log and drop terms that do not depend on the parameters through the m_i 's. This yields

$$\ln L = \sum (-m_i + n_i \ln m_i) \quad (3)$$

The parameters are the number of signal events S , the number of ϕ -only background events B_ϕ , the number of ρ -only background events B_ρ and B the number of events with neither a ϕ nor ρ .

By combining Monte Carlo, mass projection fit, and other assumptions, we can specify the fraction of events of each type that contribute to each tile – f_i^X . Thus the contributions to each tile are given by

$$m_i = f_i^S S + f_i^\phi B_\phi + f_i^\rho B_\rho + f_i^B B \quad (4)$$

where the subscript (or lack thereof) indicates the component.

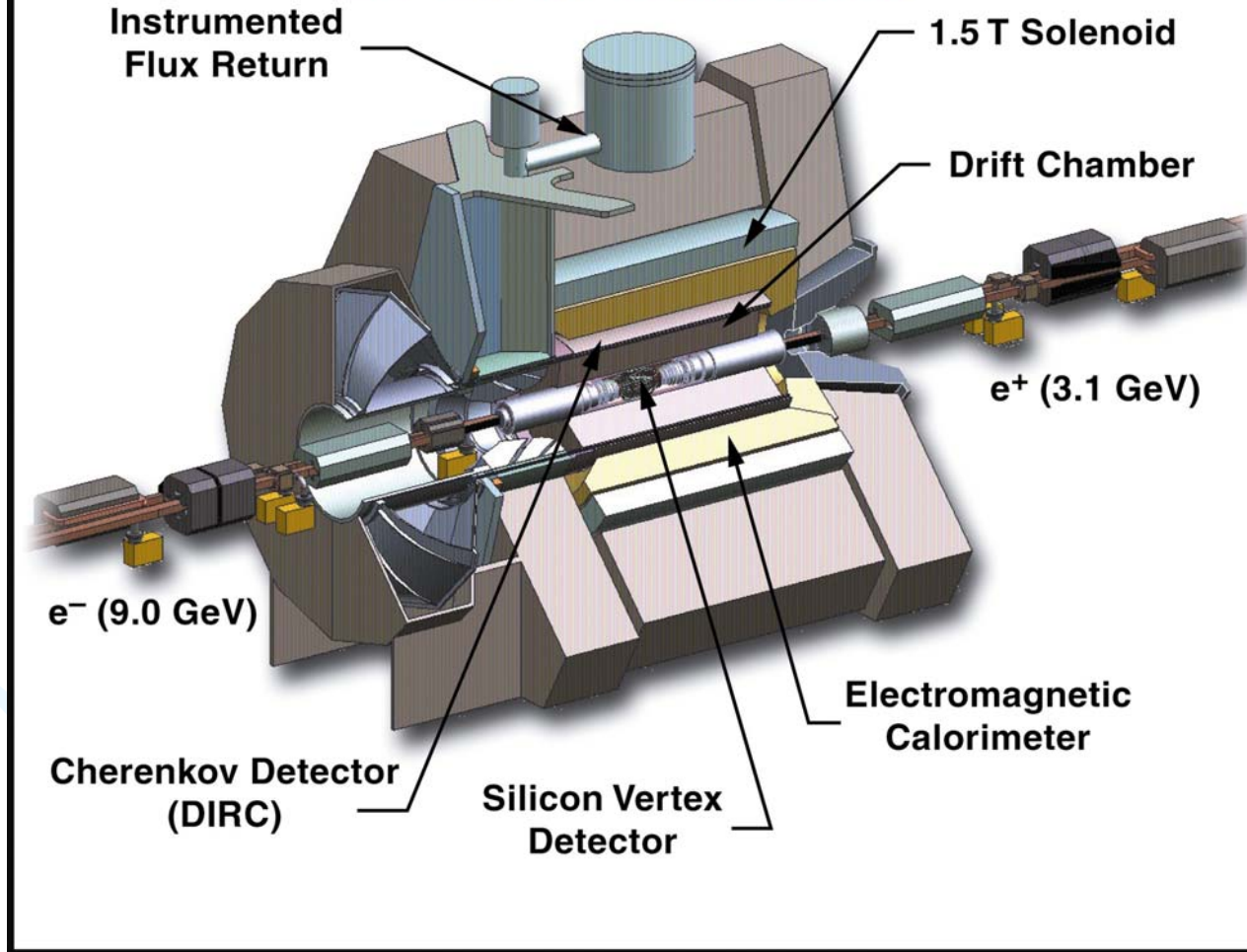
N_1	N_2	N_3
N_4	N_5	N_6
N_7	N_8	N_9

f_i^S —from Monte Carlo

f_i^ρ, f_i^ϕ —from mass projection fit

f_i^B —float slopes assuming linear

BABAR Detector



Radiative Corrections

Assuming ϕ and η decaying branching ratios to be 1 and ignoring other correction factors first, the visible cross section and the Born cross section have the following relation due to the radiation of soft photons³:

$$\sigma_{visible} = \frac{N_{observe}}{\mathcal{L}} = \int_0^{x_{max}} \sigma_B(s(1-x)) * \epsilon(x) * W(s, x) dx \quad (12)$$

where $N_{observe}$ is the number of observed signal events, \mathcal{L} is the integrated luminosity, $s = 10.58^2 \text{ GeV}^2$, $x = E_\gamma / (\sqrt{s}/2)$ is the normalized radiated photon energy, $\epsilon(x)$ is the detection efficiency for certain x , $W(s, x)$ (See Appendix B) is the QED radiator function [20], and $\sigma_B(s(1-x))$ is the Born cross section at $q^2 = s(1-x)$. Assuming σ_B has a $1/s^n$ dependence, then $\sigma_B(s(1-x)) = \sigma_B(s) \left(\frac{s}{s(1-x)}\right)^n = \sigma_B(s) * \frac{1}{(1-x)^n}$ which has a $1/s^n$ dependence, equation 12 can be rewritten as:

$$\sigma_{visible} = \sigma_B(s) * \epsilon(x=0) * (1 + \delta) \quad (13)$$

where δ is the radiative correction, and

$$(1 + \delta) = \int_0^{x_{max}} \left(\frac{1}{1-x}\right)^n * \frac{\epsilon(x)}{\epsilon(x=0)} * W(s, x) dx \quad (14)$$

The $\epsilon(x)$ is obtained from MC simulation, and the angular distributions are re-weighted by equa-