

An NLO QCD and EW analysis of the ZEUS inclusive DIS and jet cross sections

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- Proton structure
- QCD only analysis
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- Summary

The world's only e-p collider: HERA

p 920 GeV

e^+e^- 27.5 GeV

$$\rightarrow \sqrt{s} = 318 \text{ GeV}$$

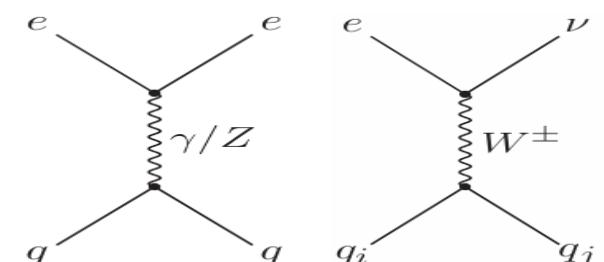
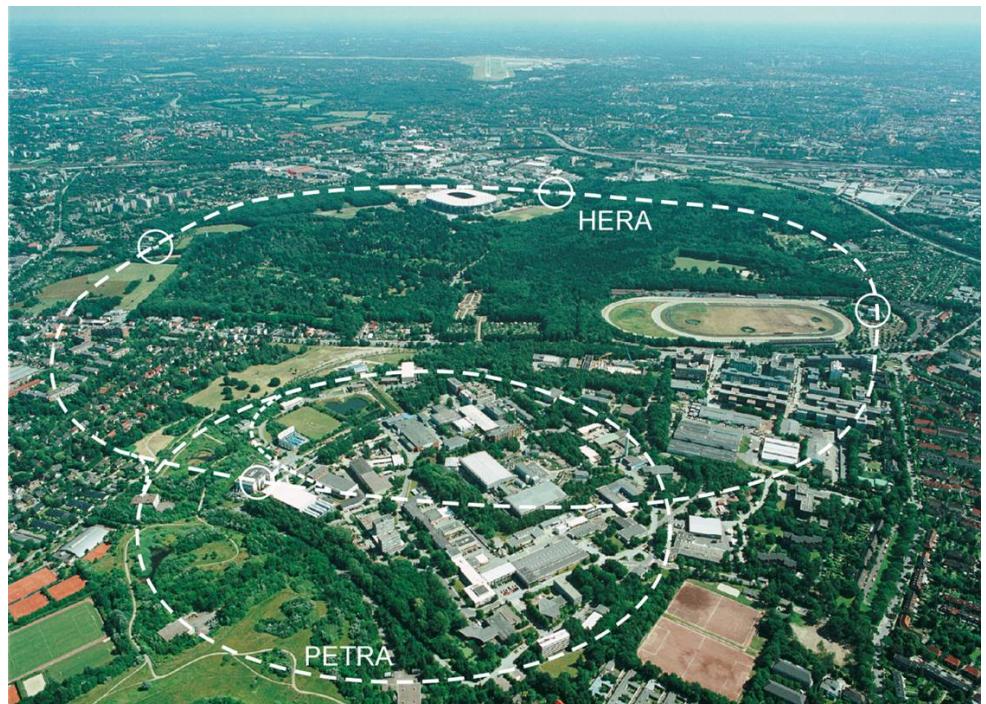
- ◆ Electrons can probe the inside of a proton.

i.e. Electrons interact with quarks in a proton.

= Deep inelastic scattering (DIS)

- ◆ eq collision @ EW energy scale.

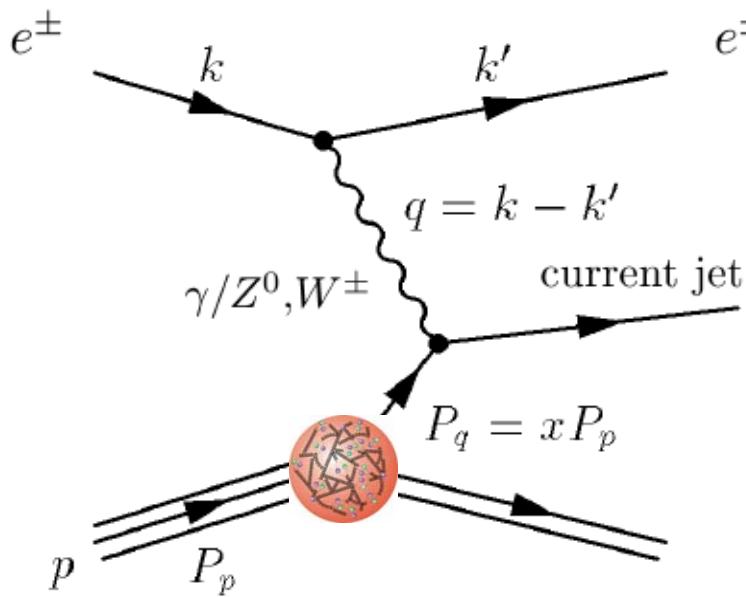
Z^0, W^\pm are exchanged in a space-like process.



High energy e-p collision at HERA →

- proton structure
- EW physics

Deep Inelastic Scattering



- ♦ DIS cross section can be described by
 - ♦ Proton has a structure.
- DIS cross section can be written with **Structure functions (SFs)**.

$$\text{e.g.) } \gamma\text{-exchange only; } \frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4} (1 + (1 - y)^2) F_2(x, Q^2)$$

$F_2=1$, if proton is point-like with charge=1

- ♦ Q^2 : Virtuality
→ probing power $\lambda \sim \frac{1}{Q^2}$
- ♦ x : Bjorken scaling variable
→ momentum fraction of struck quark
- ♦ y : Inelasticity

$$Q^2 = -q^2 = -(k - k')^2$$

$$x = \frac{Q^2}{2p \cdot q}$$

$$y = \frac{p \cdot q}{p \cdot k} = \frac{Q^2}{sx}$$

SFs parameterize how the proton differs from point-like.

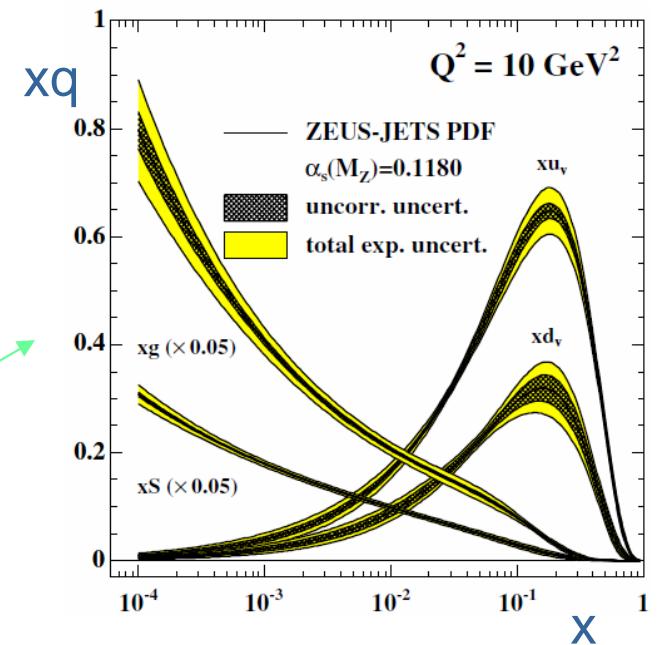
Proton structure

- ♦ In QCD, proton structure is described by Parton Density Functions (PDFs).

$q(x, Q^2)$ = number density of parton q with momentum x at Q^2 .

→ Essential to understand any physics process involving proton.

ZEUS published PDFs
(ZEUS-JETS, *Eur. Phys. J. C42*, 1-16 (2005))
 xu_v , xd_v , $xSea$, xg



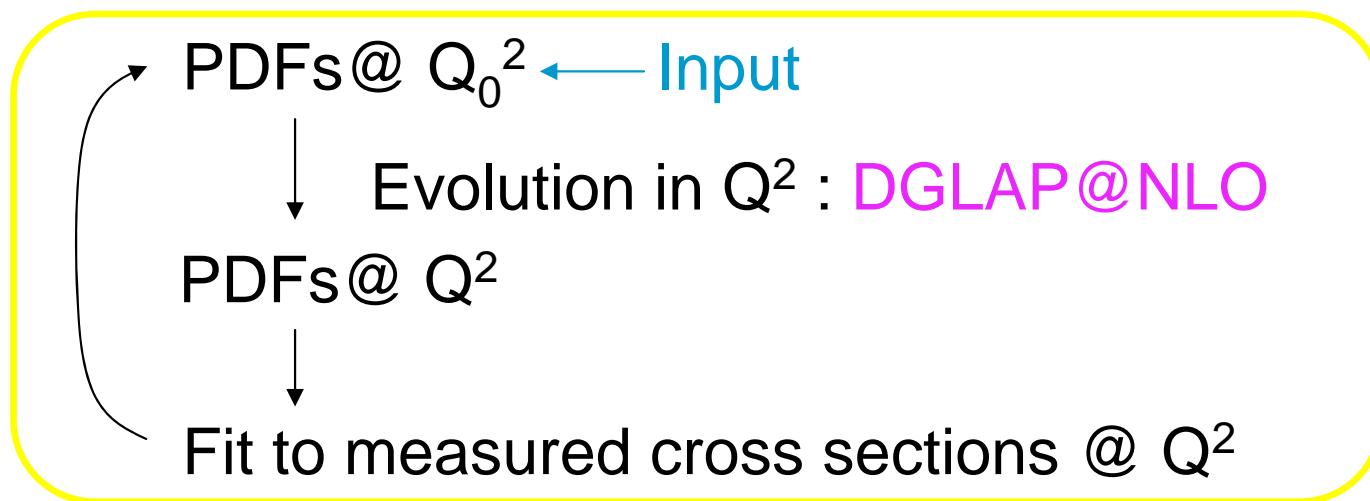
- ♦ Structure functions can be written with PDFs.

Structure functions (SFs) = $\sum_{\text{quark}} [\text{coupling}] \times [\text{PDF}]$

$$\begin{aligned} F_2 &= \sum A_q x(q + \bar{q}) && \leftarrow \text{parity conservative term} \\ xF_3 &= \sum B_q x(q - \bar{q}) && \leftarrow \text{parity violating term} \end{aligned}$$

Extraction of PDFs

- ◆ Q^2 evolution of PDFs can be predicted by perturbative QCD, i.e. by DGLAP equation.
- ◆ x -dependence of PDFs can be extracted from fits to measured cross sections.



ZEUS

PDFs are parameterized @ $Q_0^2 = 7 \text{ GeV}^2$

$$x f(x) = A x^b (1-x)^c (1+dx) \quad \text{for } xu_v, xd_v, xS, xg, x\Delta (= x\bar{d} - x\bar{u})$$

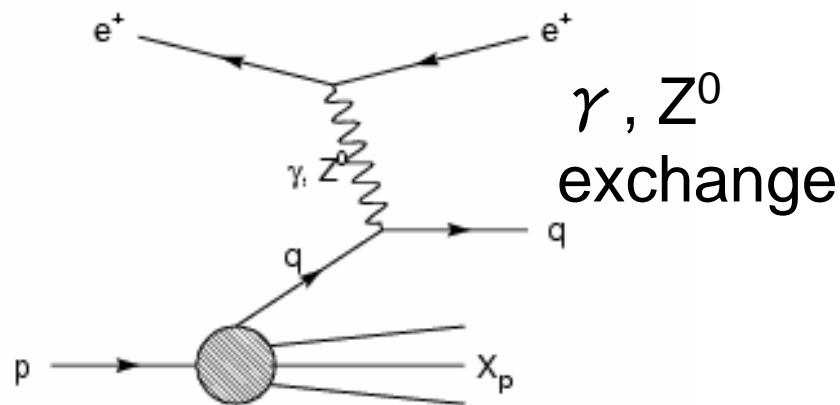
A: Normalization, b: Low x, c: High x, d: smoothing for middle x

Constraints from momentum and number sum rule, etc.

→ 11 free parameters

Cross sec. sensitive to PDF at HERA

- Neutral current DIS (NC)



$\gamma, Z^0 \rightarrow F_2 \propto \sum x(q + \bar{q})$
Sea + valence quark

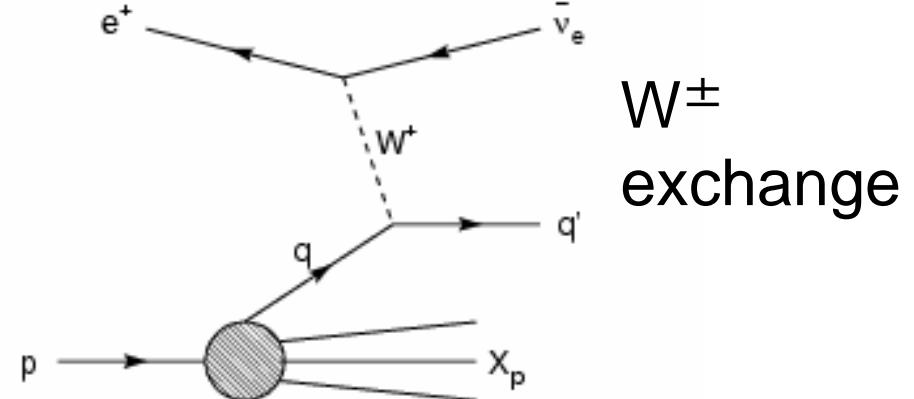
$$\frac{\partial F_2}{\partial \ln Q^2} \propto xg \quad \text{gluon}$$

Z^0 introduces parity violation.

$$\rightarrow xF_3 \propto \sum x(q - \bar{q})$$

valence quark

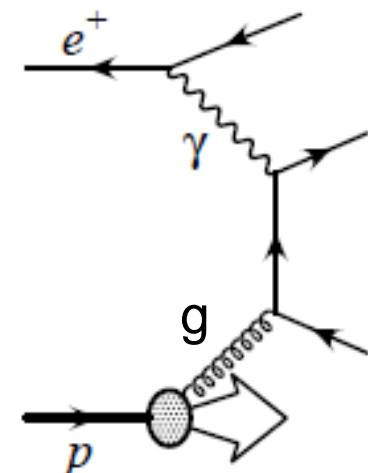
- Charged current DIS (CC)



Charge selective interaction
 e^- : u quark e^+ : d quark

- Jet process

Directly
sensitive to
gluon density

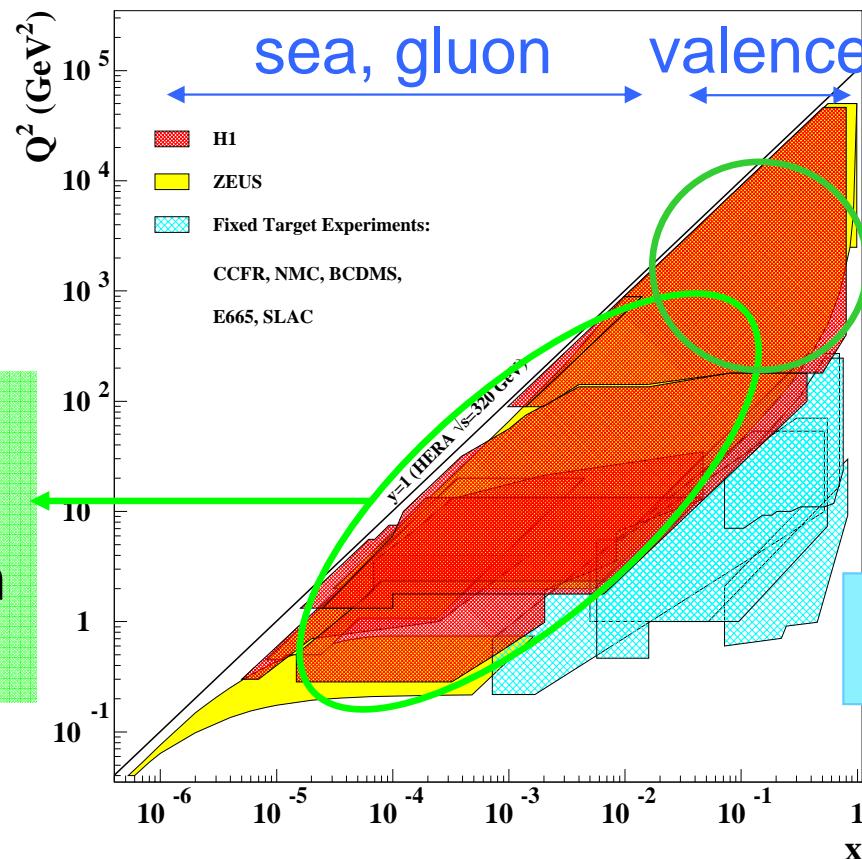


PDF extraction at ZEUS

A single experiment can determine PDFs!

Jets cross sections
→ gluon

γ exchange
→ sea
scaling violation
→ gluon



QCD + EW physics

Z^0 exchange
→ • sea + valence
• valence only

W^\pm exchange
→ u or d quark

- ♦ Pure proton target → Free from target correction, nuclear effect.
- ♦ Single experiment → systematic uncertainties are well understood.

Furthermore, based on own knowledge of PDFs, EW physics can be studied.

ZEUS fits

DIS is a convoluted phenomenon of electron-quark scattering and proton structure.

$$\sigma(ep) \propto \sum \sigma(eq) \otimes (\text{PDF})$$

↓ ↓
EW QCD

A fit on $\sigma(ep)$ is a combined analysis of QCD and EW.

- ◆ Data (In total, ~800 data points)
 - 94-00 inclusive e^-p/e^+p NC/CC cross sections
 - 96-97 Jet cross sections in DIS and PHP
 - 04-05 polarized e^-p NC/CC inclusive cross sections (prel.)
- ◆ QCD only analysis : ZEUS-pol fit
Only PDFs are free to see the impact of the new data.
- ◆ Combined QCD and EW analysis
EW parameters and PDFs are determined simultaneously to exploit HERA sensitivity fully.

Published
PDFs

→ New data (See next slide)

QCD only analysis: ZEUS-pol fit

HERA is now running as **HERA II** since 2003 with upgrades;

- Large luminosity
- Polarized e^\pm beam

ZEUS has measured e^-p NC/CC inclusive cross sections in HERA II.

→ **Much statistics at High Q^2 with Polarized electrons**

NC/CC electron data

HERA I	16pb^{-1}
HERA II	121.5pb^{-1}

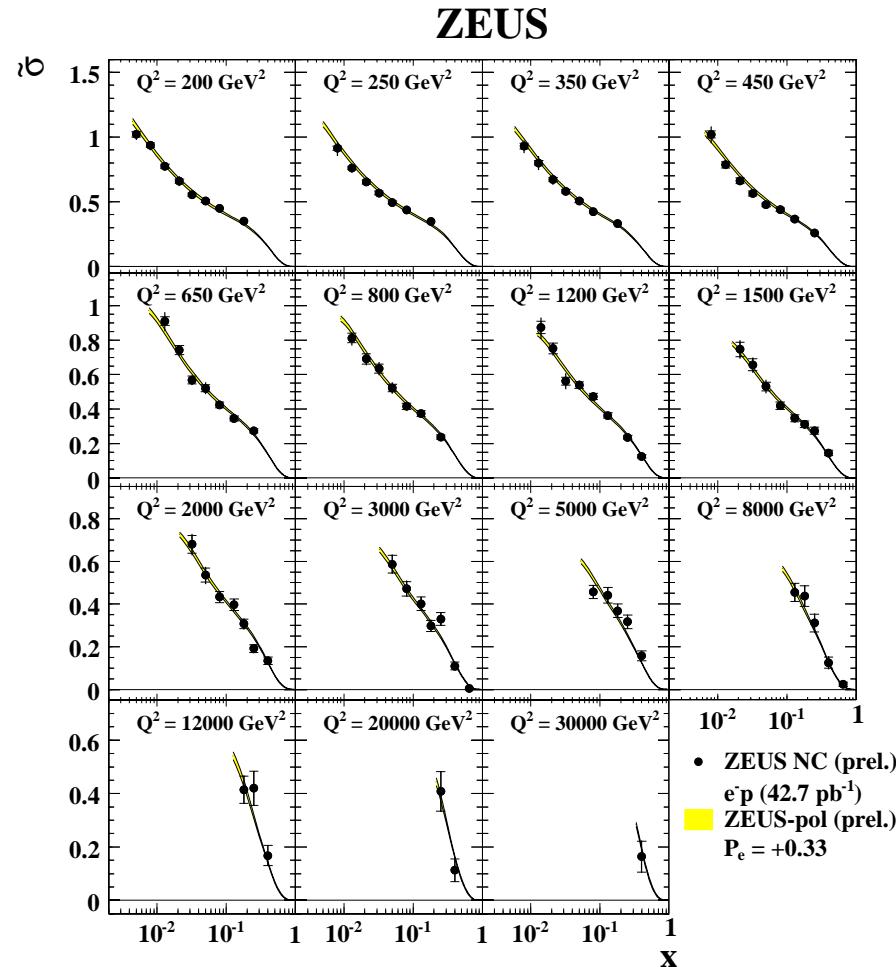
- better determination of PDFs at high x (\leftarrow high Q^2).
- better sensitivity to EW

ZEUS performed the first fit including the HERA II cross sections.

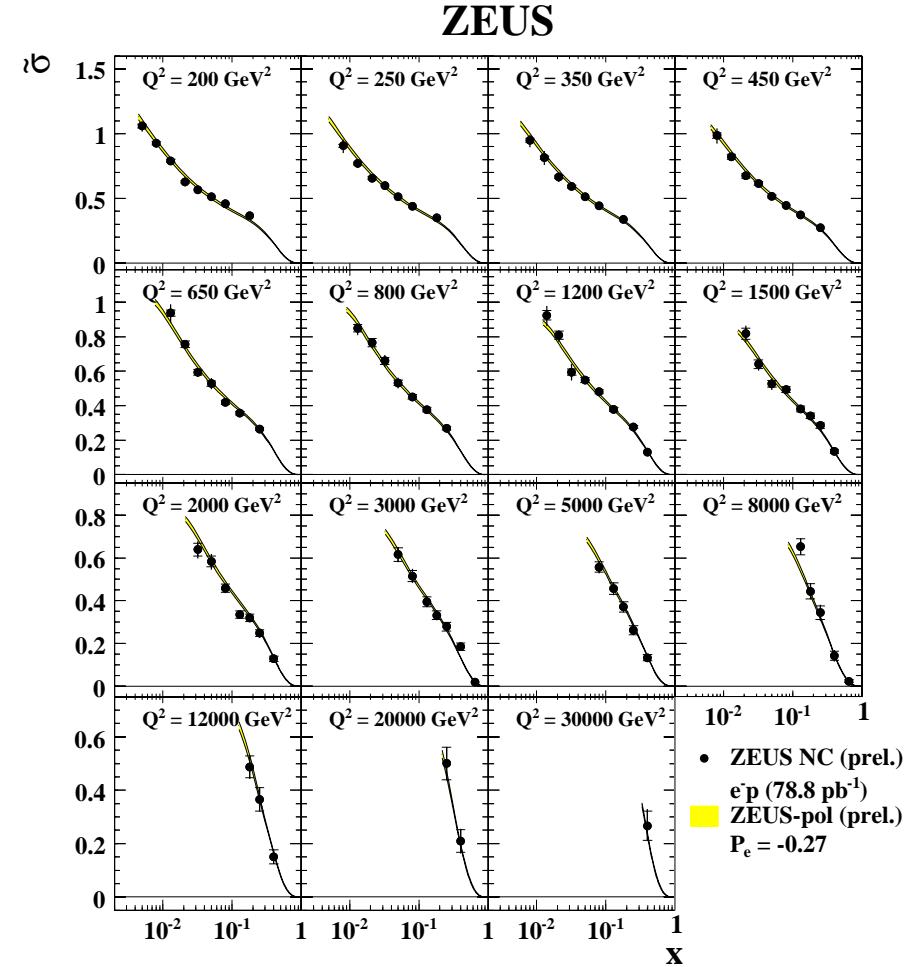
ZEUS-pol fit (prel.)

- ◆ The first fit including polarized electron DIS cross sections.
- ◆ The fit to see any effect on PDFs by including these new data.
- ◆ All EW parameters are fixed to SM values.

Polarized NC cross sections



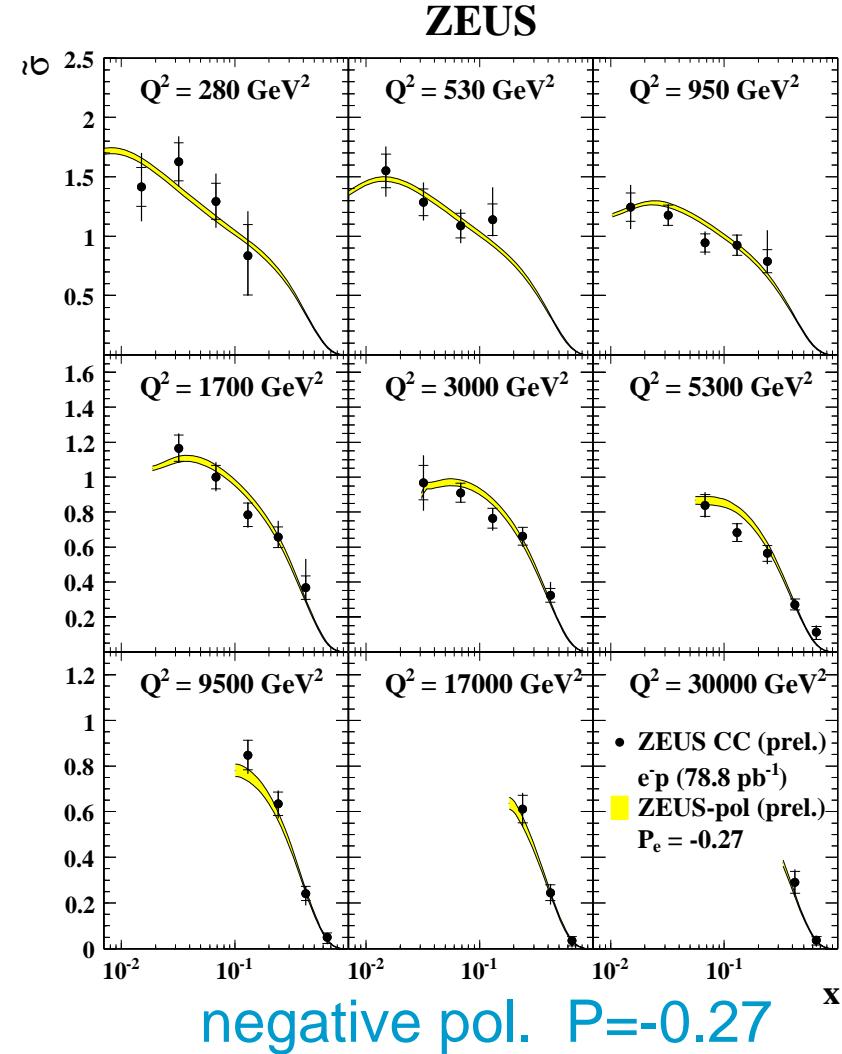
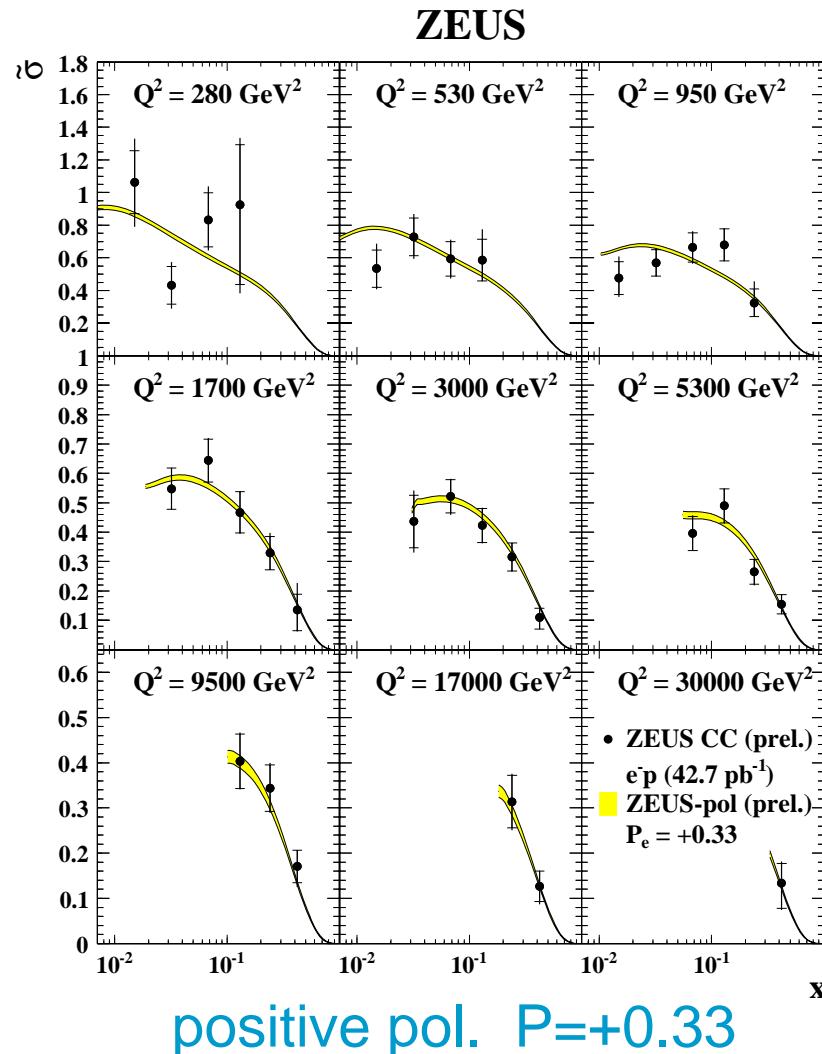
positive pol. $P=+0.33$



negative pol. $P=-0.27$

Data is well described by ZEUS-pol Fit.
 The polarized cross sections from HERA-II were
 successfully fitted for the first time.

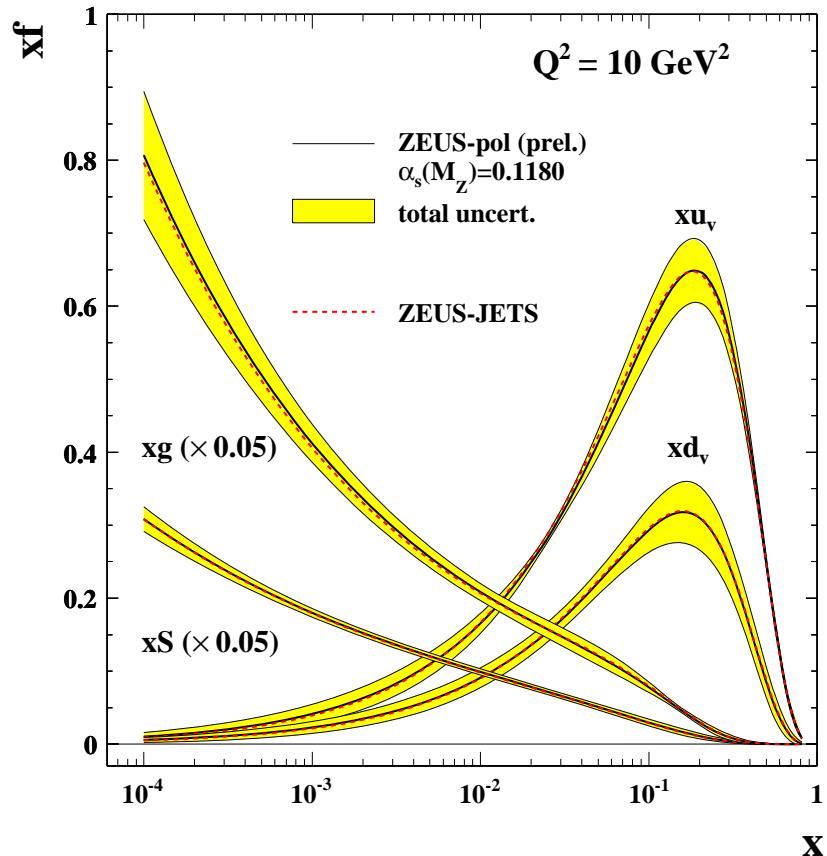
Polarized CC cross sections



Data is well described by ZEUS-pol Fit.
 The polarized cross sections from HERA-II were successfully fitted for the first time.

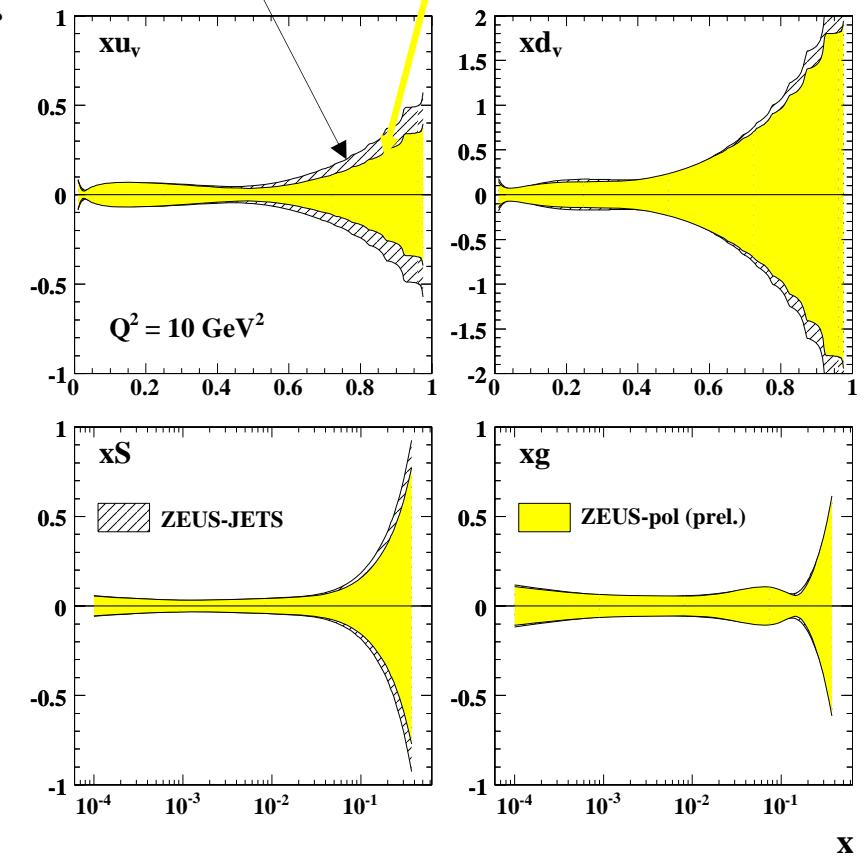
PDFs

ZEUS



ZEUS-JETS
(without HERA II)

ZEUS-pol
(with HERA II)



- ◆ Central values of PDFs are almost unchanged by addition of HERA II electron data.
- ◆ Uncertainties are reduced. – high- x and particularly on xu_v

$$e^-p: \quad e_u = \frac{2}{3}e, \quad e_d = -\frac{1}{3}e \rightarrow \sigma_{NC} \propto (4u+d), \quad \sigma_{CC} \propto u$$

Combined QCD and EW analysis

- ◆ Precise understanding of PDFs at HERA allows to see physics in eq scattering at EW scale ($Q^2 \sim M_W^2, M_Z^2$).
- ◆ **HERA II data:**
In addition to the **large statistics** at EW energy scale, **polarization** gives direct sensitivity to EW.
→ Let's exploit the sensitivity to determine EW parameters!

A combined QCD + EW analysis

EW parameters and PDFs are determined simultaneously.

← The correlation between them is taken into account automatically in the fit.

1. Extraction of M_W

(ZEUS-pol- M_W , ZEUS-pol-g- M_W) ←CC cross sections

2. Extraction of quark couplings to Z

(ZEUS-pol- a_u - v_u , ZEUS-pol- a_d - v_d etc.) ←NC cross sections³

Extraction of M_W

- ◆ CC cross section

M_W is space-like.

$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2} = \frac{G_F^2}{4\pi x} \frac{M_W^4}{(Q^2 + M_W^2)^2} [Y_+ F_2^{CC}(x, Q^2) - y^2 F_L^{CC} \mp Y_- x F_3^{CC}(x, Q^2)]$$

- ◆ M_W and PDF parameters are free:

(M_W contributes also to normalization due to fixed G_F .)

$M_W = 79.1 \pm 0.77 \text{ (stat+uncorr)} \pm 0.9 \text{ (corr.sys)} \text{ [GeV]}$ prel.

HERA I result:

$M_W = 78.9 \pm 2.0 \text{ (stat)} \pm 1.8 \text{ (sys)} {}^{+2.2}_{-1.8} \text{ (PDF)} \text{ [GeV]}$

- ◆ Determination of M_W as general ‘propagator mass’ with general coupling g ($= G_F M_W^2$)

$g = 0.0772 \pm 0.0021 \pm 0.0019$

prel.

$M_W = 82.8 \pm 1.5 \pm 1.3 \text{ [GeV]}$

$$\frac{1}{4\pi x} \frac{g^2}{(Q^2 + M_W^2)^2}$$

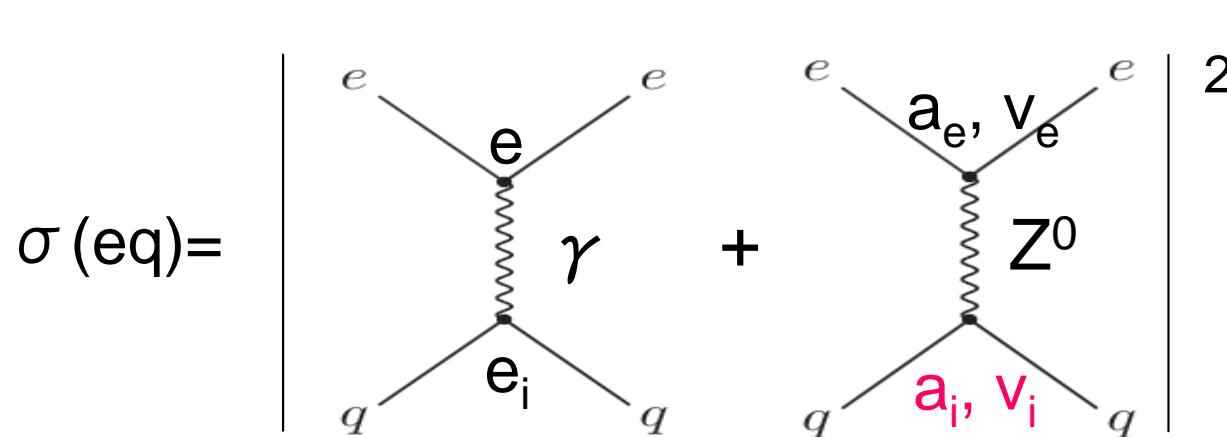
The combined QCD+EW analysis on HERA I + II improves M_W .
Complementary and consistent with time-like M_W at LEP/Tevatron.

NC cross sections

At HERA, NC interaction occurs by γ and Z^0 exchange.

→ electron-quark cross sections are sum of three terms;

γ -term, γZ interference-term, Z -term



a : axial coupling
v : vector coupling

In SM formalism,

$$a_i = T_i^3$$

$$v_i = T_i^3 - 2e_i \sin^2 \theta_W$$

- ◆ DIS cross sections are sensitive to **quark couplings to Z** (a_i, v_i).
- ◆ They can be determined together with PDFs.

Note: a, v are parameterizations of the couplings in the most general way. i.e. less SM formalism.

Polarized NC cross sections

$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4} \left[\underbrace{(Y_+ F_2^0 \mp Y_- x F_3^0)}_{\text{unpol.}} \mp P \underbrace{(Y_+ F_2^P \mp Y_- x F_3^P)}_{\text{pol.}} \right] \text{ with polarization } P$$

Structure functions: $F_2^{0,P} = \sum_i A_i^{0,P}(Q^2) [xq_i(x, Q^2) + x\bar{q}_i(x, Q^2)]$
 $xF_3^{0,P} = \sum_i B_i^{0,P}(Q^2) [xq_i(x, Q^2) - x\bar{q}_i(x, Q^2)]$

unpolarized coefficients

$$A_i^0(Q^2) = \frac{e_i^2}{\gamma} - 2e_i v_i v_e P_Z + (v_e^2 + a_e^2)(v_i^2 + a_i^2) P_Z^2$$

$$B_i^0(Q^2) = \frac{-2e_i a_i a_e P_Z}{\gamma Z^0} + 4a_i a_e v_i v_e P_Z^2$$

$$P_Z = \frac{1}{\sin^2 2\theta} \frac{Q^2}{(M_Z^2 + Q^2)}$$

: quarks

polarized coefficients

$$A_i^P(Q^2) = \frac{2e_i v_i a_e P_Z}{\gamma Z^0} - 2v_e a_e (v_i^2 + a_i^2) P_Z^2$$

$$B_i^P(Q^2) = \frac{2e_i a_i v_e P_Z}{\gamma Z^0} - 2v_i a_i (v_e^2 + a_e^2) P_Z^2$$

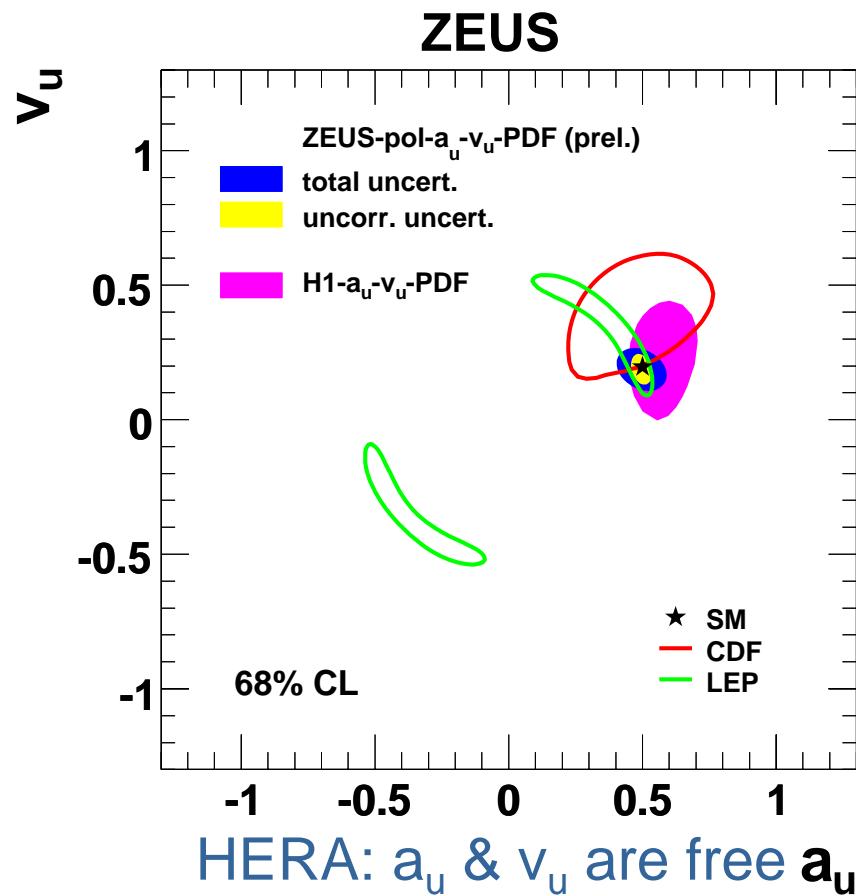
v_e is very small (~ 0.04).
 $P_Z \gg P_Z^2$ (\sim middle Q^2)

unpolarized $xF_3 \rightarrow a_i$,
 polarized $F_2 \rightarrow v_i$

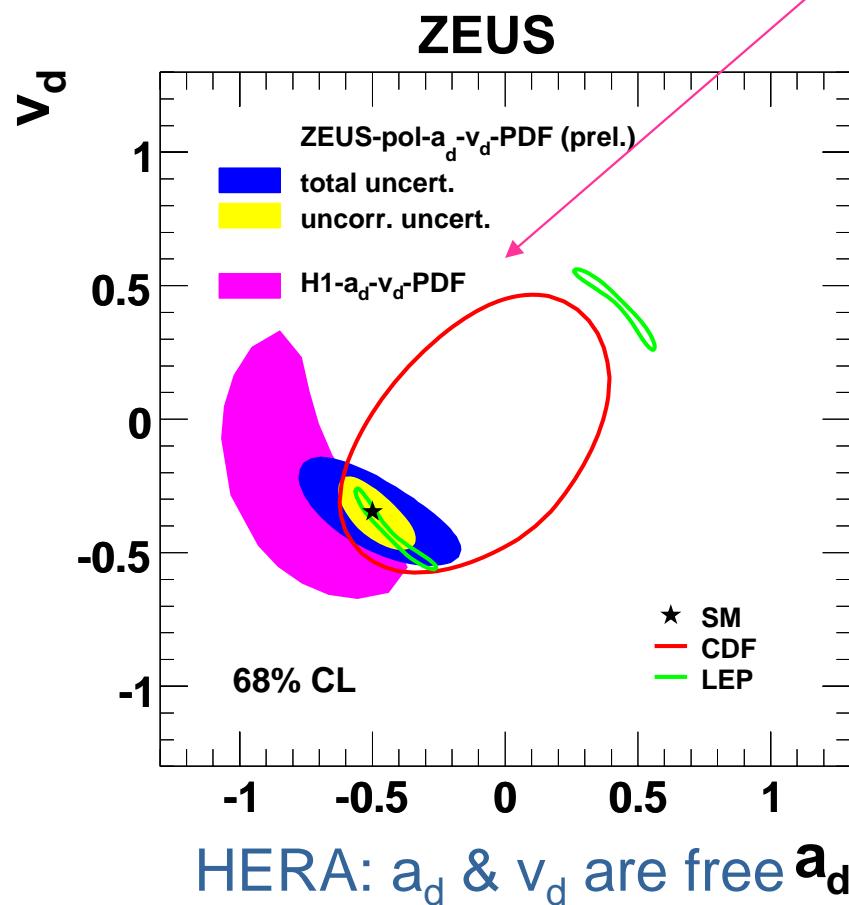
Quark couplings to Z

a_i , v_i and PDFs are determined simultaneously.

H1 performed similar fit but without HERA II data.



- Comparing to H1 fit, HERA II polarized data improves the sensitivity, especially to vector couplings.
- ZEUS-pol- a_i - v_i fit shows excellent constraint on quark couplings.



QCD+EW fit: Using SM relation

- ♦ In SM formalism, $a_q = T_q^3$

$$v_q = T_q^3 - 2e_q \sin^2 \theta_W$$

→ Determine T_u^3 , T_d^3 , $\sin^2 \theta_W$: 3 EW parameters

Note: $\sin^2 \theta_W$ is also in Z exchange term (P_Z)

Results: (preliminary)

$$T_u^3 = 0.47 \pm 0.05 \pm 0.13$$

$$T_d^3 = -0.55 \pm 0.18 \pm 0.35$$

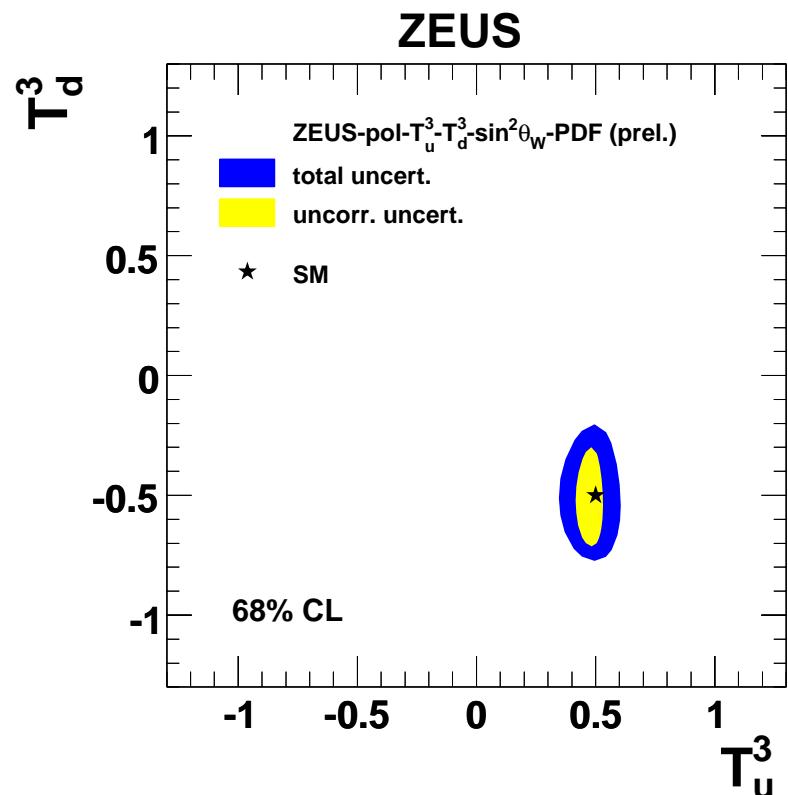
$$\sin^2 \theta_W = 0.231 \pm 0.024 \pm 0.070$$

Good agreement with SM values.

$$T_u^3 = 0.5$$

$$T_d^3 = -0.5$$

$$\sin^2 \theta_W = 0.2315$$



Summary

- ◆ HERA has sensitivities to both EW and QCD physics.
- ◆ We have HERA II data.
 - Large luminosity at high- Q^2 with polarized electrons.
- ◆ New fit including HERA II data: ZEUS-pol fit
 - HERA II data is well described and fitted.
 - Uncertainties of PDFs are reduced.
- ◆ EW parameters are extracted from combined analysis of EW and PDFs (ZEUS-pol-Mw fit, etc).
 - Extracted M_W is consistent with the world average value.
 - Quark couplings are determined with excellent precision. They are well consistent with SM.

Back up slides

Extraction of PDFs at ZEUS

- ◆ PDFs: parameterization @ $Q_0^2 = 7 \text{ GeV}^2$

$x f(x) = A x^b (1 - x)^c (1 + dx)$

 for $xu_v, xd_v, xS, xg, x\Delta (= x\bar{d} - x\bar{u})$

A: Normalization, b: Low x, c: High x, d: smoothing for middle x
- ◆ Constraints
 - Momentum and number sum rule $\rightarrow A_{uv}, A_{dv}, A_g$
 - Equal behaviour of u_v and d_v at low x $\rightarrow b_{uv} = b_{dv}$
 - Δ : consistent with Gottfried sum rule and Drell Yan (CCFR)
- ◆ 11 free parameters
- ◆ DGLAP evolution at NLO (MSbar)
- ◆ Heavy quarks are treated in variable flavour-number scheme of Thorne and Roberts.
- ◆ Corr. syst. uncertainties are evaluated using OFFSET method.

PDF Parameterization

u-valence ($x u_v$)	$A_{uv} x^{b_{uv}} (1-x)^{c_{uv}} (1+d_{uv}x)$
d-valence ($x d_v$)	$A_{dv} x^{b_{dv}} (1-x)^{c_{dv}} (1+d_{dv}x)$
Sea ($x S$)	$A_S x^{b_S} (1-x)^{c_S}$
gluon ($x g$)	$A_g x^{b_g} (1-x)^{c_g} (1+d_g x)$
dbar-ubar ($x \Delta$)	$0.27 x^{0.5} (1-x)^{c_\Delta}$

Constraints

- Momentum and number sum rule
- Equal behaviour of u_v and d_v at low x
- Δ : consistent with Gottfried sum rule and Drell Yan

11 free parameters

OFFSET method

χ^2 is defined as

$$\chi^2 = \sum_i \frac{[F_i^{\text{QCD}}(p) + \sum_{\lambda} s_{\lambda} \Delta_{i\lambda}^{\text{sys}} - F_i^{\text{meas}}]^2}{(\sigma_i^{\text{stat}}{}^2 + \sigma_i^{\text{unc.sys}}{}^2)} + \sum_{\lambda} s_{\lambda}^2$$

F_i^{QCD} : prediction from QCD

F_i^{meas} : measured data point

s_{λ} : fit parameter of systematic uncertainty $\Delta_{i\lambda}^{\text{sys}}$: correlated systematic uncertainty

σ_i^{stat} : statistical uncertainty

$\sigma_i^{\text{unc.sys}}$: uncorrelated systematic uncertainty

1. Central values are extracted without any correlated systematic uncertainties ($s_{\lambda}=0$).
2. For each source of correlated systematic uncertainty (i.e. for each λ):
 - Data points are shifted to the limit of the uncertainty ($s_{\lambda}=\pm 1$).
 - Deviation from the central value is extracted by re-doing the fit.
3. Add all deviations in quadrature

No assumption of gaussian shape for correlated systematic uncertainties.

Conservative method.

HERA I : ZEUS-JETS fit

First fit using HERA jets data.

→ Making use of full potential of ZEUS data (and alone) in HERA I.

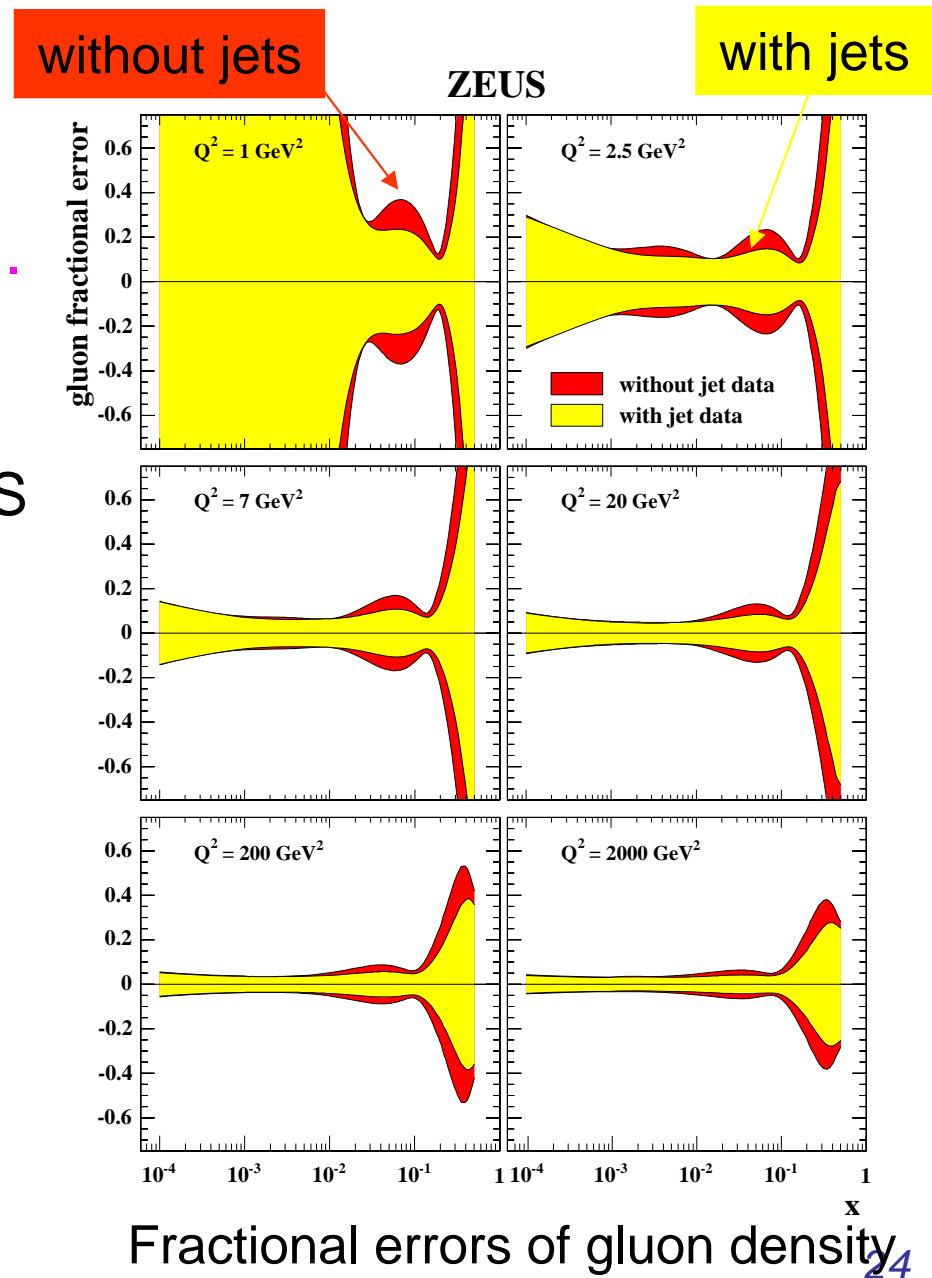
- HERA I inclusive NC/CC cross sections (94-00)
- Inclusive jets cross sections in DIS (96-97)
- Dijets in photoproduction (96-97)

Single experiment

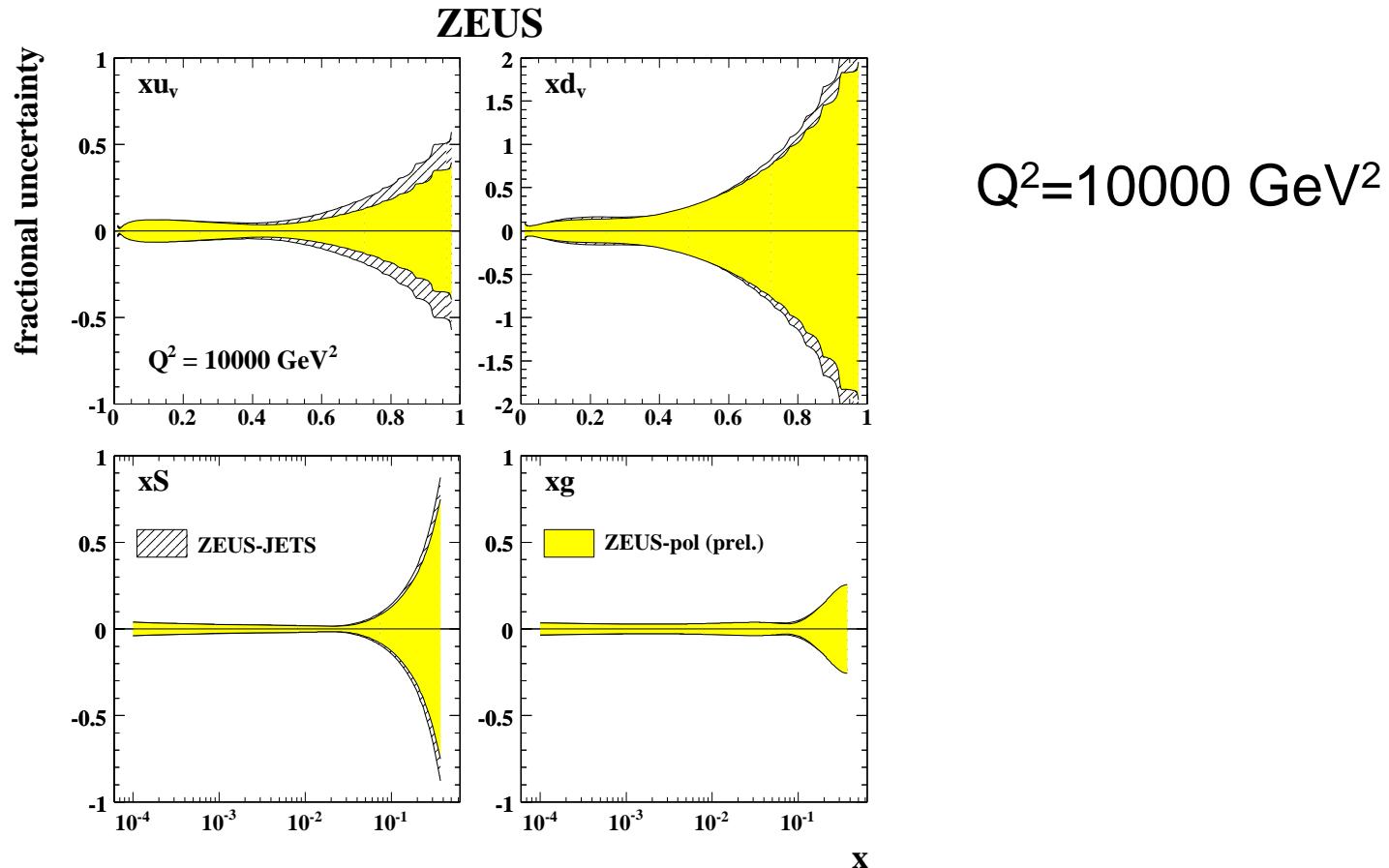
→ systematic uncertainties are well understood.

Jets cross sections

→ sensitive to gluon density.



PDF uncertainties at very High Q^2



- ◆ Improvement of PDF uncertainties is also seen at $Q^2=10^4 \text{ GeV}^2$. Good news for LHC physics.
- ◆ HERA is now running with positron beam.
→ Further improvement can be expected in future.

Extraction of quark couplings to Z

Axial/vector couplings of u/d-type quark: 4 couplings

→ 2 of them are free and fitted together with PDFs: 4 fits in total

Results (preliminary)	a_u	a_d	v_u	v_d
SM	0.5	-0.5	0.196	-0.346
ZEUS-pol- a_u - v_u fit	$0.50 \pm 0.04 \pm 0.09$	fixed	$0.19 \pm 0.06 \pm 0.06$	fixed
ZEUS-pol- a_d - v_d fit	fixed	$-0.49 \pm 0.14 \pm 0.28$	fixed	$-0.37 \pm 0.14 \pm 0.16$
ZEUS-pol- a_u - a_d fit	$0.48 \pm 0.06 \pm 0.10$	$-0.55 \pm 0.10 \pm 0.21$	fixed	fixed
ZEUS-pol- v_u - v_d fit	fixed	fixed	$0.12 \pm 0.10 \pm 0.05$	$-0.47 \pm 0.15 \pm 0.19$

- ◆ Note: These fits parameterize the couplings in most general way.
- ◆ They are in good agreement with SM predictions.
→ Contours will be shown in the next slides.

Extraction of M_W (2)

- ◆ Determination of BOTH G_F and M_W
(ZEUS-pol- G_F - M_W fit)

$$\frac{G_F^2}{4\pi x} \frac{M_W^4}{(Q^2 + M_W^2)^2}$$

$$G_F = 1.127 \pm 0.013 \pm 0.014 \times 10^{-5} [\text{GeV}^{-2}]$$

$$M_W = 82.8 \pm 1.5 \pm 1.3 [\text{GeV}]$$

preliminary

- ◆ Determination of M_W as more general ‘propagator mass’ with general coupling g
(ZEUS-pol- g - M_W fit)

$$\frac{1}{4\pi x} \frac{g^2}{(Q^2 + M_W^2)^2}$$

$$g = 0.0772 \pm 0.0021 \pm 0.0019$$

$$M_W = 82.8 \pm 1.5 \pm 1.3 [\text{GeV}]$$

preliminary

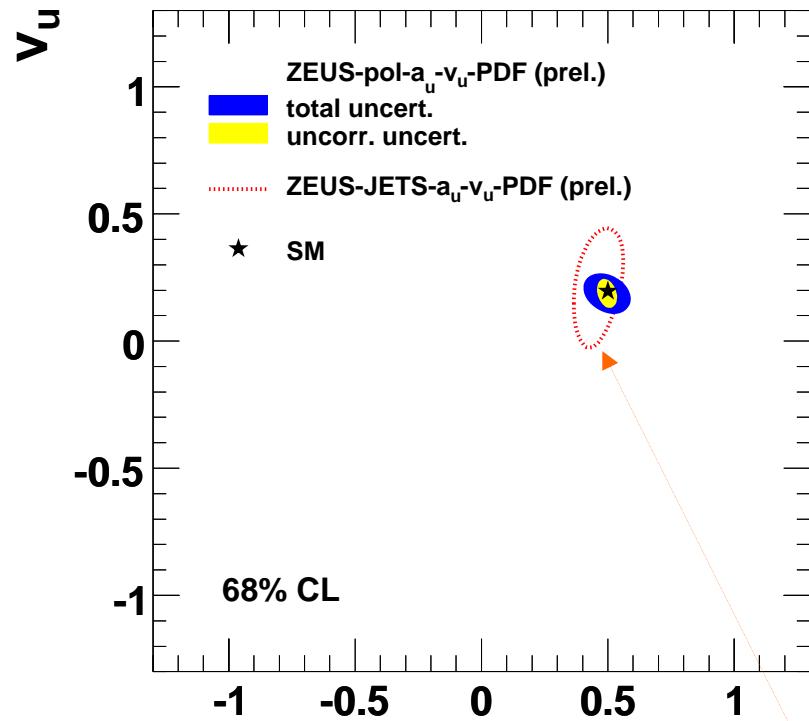
- ◆ They are in good agreement with the world average values.

$$G_F = 1.16639 \times 10^{-5} \text{ GeV}^{-2}$$

$$M_W = 80.4 \text{ GeV}$$

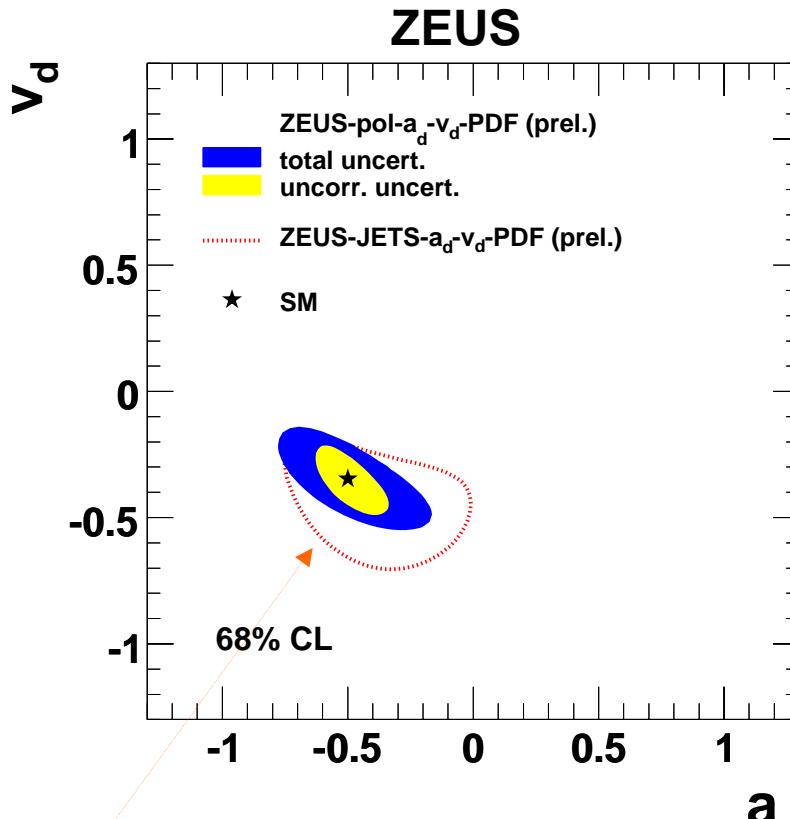
$$g = G_F M_W^2 = 0.07542$$

a_i vs. v_i ZEUS



	a_u	a_d	v_u	v_d
SM	0.5	-0.5	0.196	0.346
ZEUS-pol- a_u - v_u fit	0.50 $\pm 0.04 \pm 0.09$	fixed	0.19 $\pm 0.05 \pm 0.06$	fixed
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ZEUS-pol- a_u - a_d fit	0.48 $\pm 0.06 \pm 0.10$	-0.55 $\pm 0.10 \pm 0.21$	fixed	fixed
ZEUS-pol- v_u - v_d fit	fixed	0.12 $\pm 0.10 \pm 0.05$	0.47 $\pm 0.15 \pm 0.19$	fixed

ZEUS-pol- a_u - v_u
 a_d, v_d : fixed



	a_u	a_d	v_u	v_d
SM	0.5	-0.5	0.196	0.346
ZEUS-pol- a_u - v_u fit	0.50 $\pm 0.04 \pm 0.09$	fixed	0.19 $\pm 0.05 \pm 0.06$	fixed
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ZEUS-pol- a_u - a_d fit	0.48 $\pm 0.06 \pm 0.10$	-0.55 $\pm 0.10 \pm 0.21$	fixed	fixed
ZEUS-pol- v_u - v_d fit	fixed	0.12 $\pm 0.10 \pm 0.05$	0.47 $\pm 0.15 \pm 0.19$	fixed

ZEUS-pol- a_d - v_d
 a_u, v_u : fixed

We also extract couplings without HERA II data with same parameter settings (--- ZEUS-JETS- a_i - v_i fit)

HERA II data constrains the quark couplings well.
 They agree well with SM prediction.

v_u VS. v_d

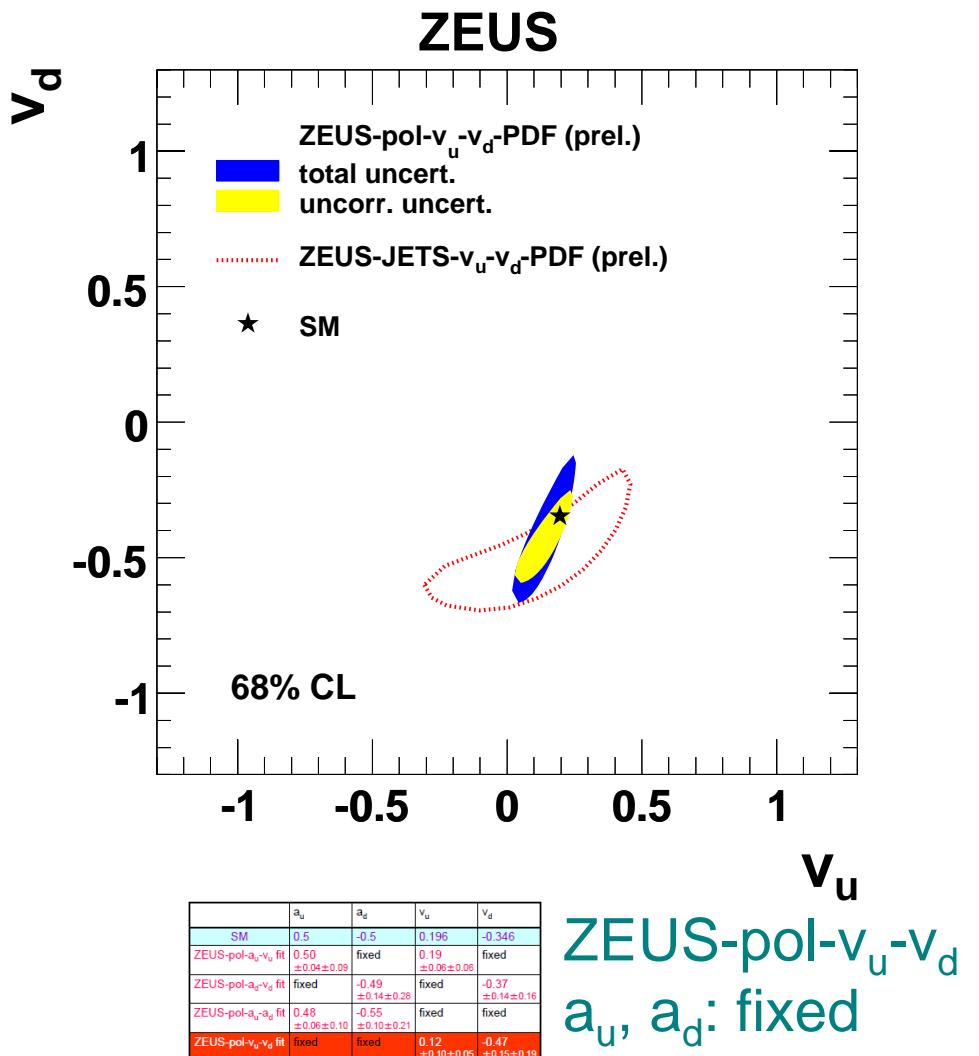
Now we have polarized data!

→ Vector couplings benefit from extra sensitivity from polarized F_2 .

Reminder:

unpolarized $xF_3 \rightarrow a_i$,
polarized $F_2 \rightarrow v_i$

- ◆ v_u and v_d are determined well by the fit with HERA II.
– especially on v_u .



Right handed Isospin

- Introduce right handed isospin, $T^3_{q,R}$, which should be 0 in SM,

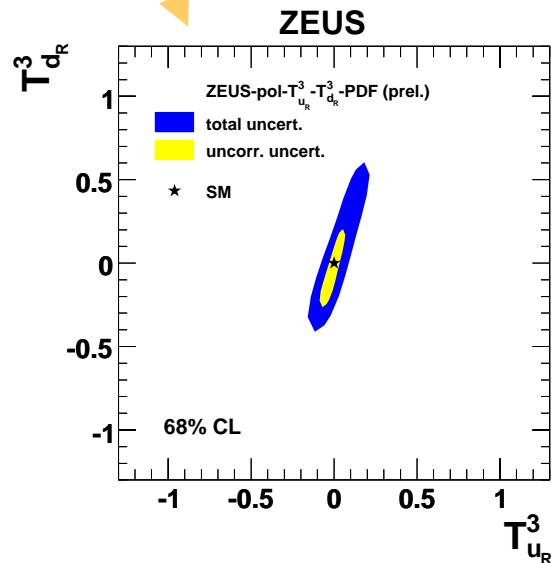
$$a_q = T^3_{q,L} + T^3_{q,R}$$

$T^3_{q,L}$ are fixed:

$$\nu_q = T^3_{q,L} - T^3_{q,R} - 2e_q \sin^2 \theta_W$$

$$T^3_{u,L} = 1/2, T^3_{d,L} = -1/2$$

Results (preliminary)	$T^3 u_R$	$T^3 d_R$	$\sin^2 \theta_W$
ZEUS-pol- $T^3_{u,R}$ - $T^3_{d,R}$ fit	-0.04 $\pm 0.06 \pm 0.13$	-0.14 $\pm 0.18 \pm 0.33$	0.2315 fixed
ZEUS-pol- $T^3_{u,R}$ - $T^3_{d,R}$ - $\sin^2 \theta_W$ fit	-0.07 $\pm 0.07 \pm 0.07$	-0.26 $\pm 0.19 \pm 0.19$	0.238 $\pm 0.011 \pm 0.023$



No deviation from SM is seen.
They are well constrained by the fits.