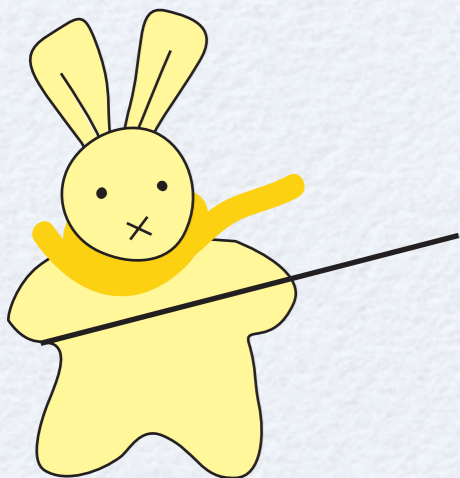


Presence and Absence of Degeneracies in the LBL NuOscExp

Masafumi Koike (Saitama Univ.)
Masako Saito (Saitama Univ.)



Neutrino Oscillation

- Evolution equation of neutrinos

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & & \\ & \delta m_{21}^2 & \\ & & \delta m_{31}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} a & & \\ & 0 & \\ & & 0 \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \equiv H \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

- Quadratic mass differences

$$\delta m_{ij}^2 \equiv m_i^2 - m_j^2$$

$$\delta m_{21}^2 \simeq 8 \times 10^{-5} \text{ eV}^2,$$

$$|\delta m_{31}^2| \simeq (2 - 2.5) \times 10^{-3} \text{ eV}^2$$

(Fogli *et al.* 2005)

$$\begin{cases} \delta m_{31}^2 > 0 & \text{"Normal" hierarchy} \\ \delta m_{31}^2 < 0 & \text{"Inverted" hierarchy} \end{cases}$$

- Mixing matrix, mixing angles and CP phase(s)

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{\text{CP}}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{\text{CP}}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & & \\ & e^{i\phi_1} & \\ & & e^{i\phi_2} \end{pmatrix}$$

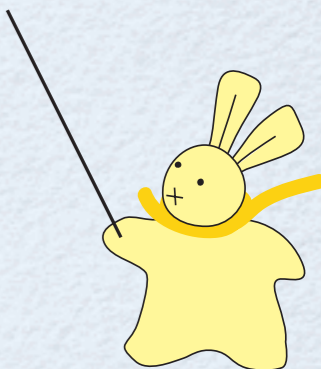
$$c_{ij} \equiv \cos \theta_{ij} \quad s_{ij} \equiv \sin \theta_{ij}$$

$$\sin^2 \theta_{12} \sim 0.3, \quad \sin^2 \theta_{23} \sim 0.45, \quad \sin^2 \theta_{13} \lesssim 0.04 \quad (\text{Fogli } et \text{ al. } 2005)$$

- Matter effect

$$a = 2\sqrt{2}G_{\text{F}}n_e E = 7.63 \times 10^{-5} \text{ eV}^2 \frac{\rho}{\text{g cm}^3} \frac{E}{\text{GeV}}$$

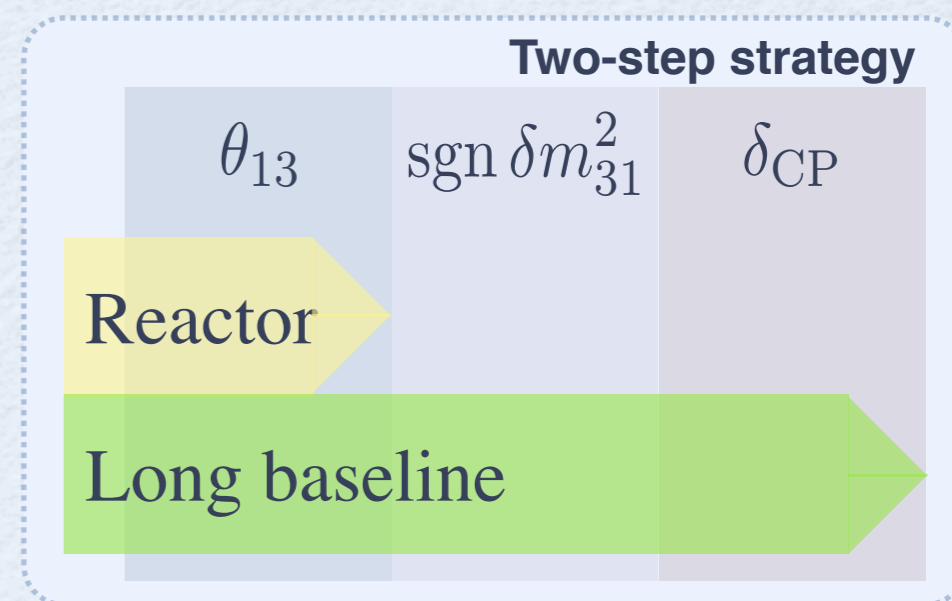
- Unknowns to date: θ_{13} , $\text{sgn } \delta m_{31}^2$, δ_{CP} , ϕ_1, ϕ_2



Scope of this Talk

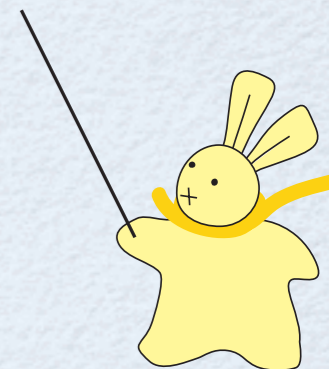
● Situations considered

- θ_{13} is constrained in advance of the future generation of long-baseline experiment.
- The value of θ_{13} is not too small ($\sin^2 2\theta_{13} \gtrsim 0.01$) so that the CP -violation effect is accessible.

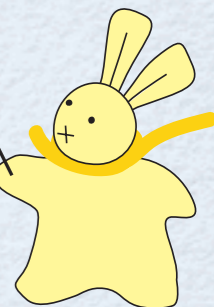
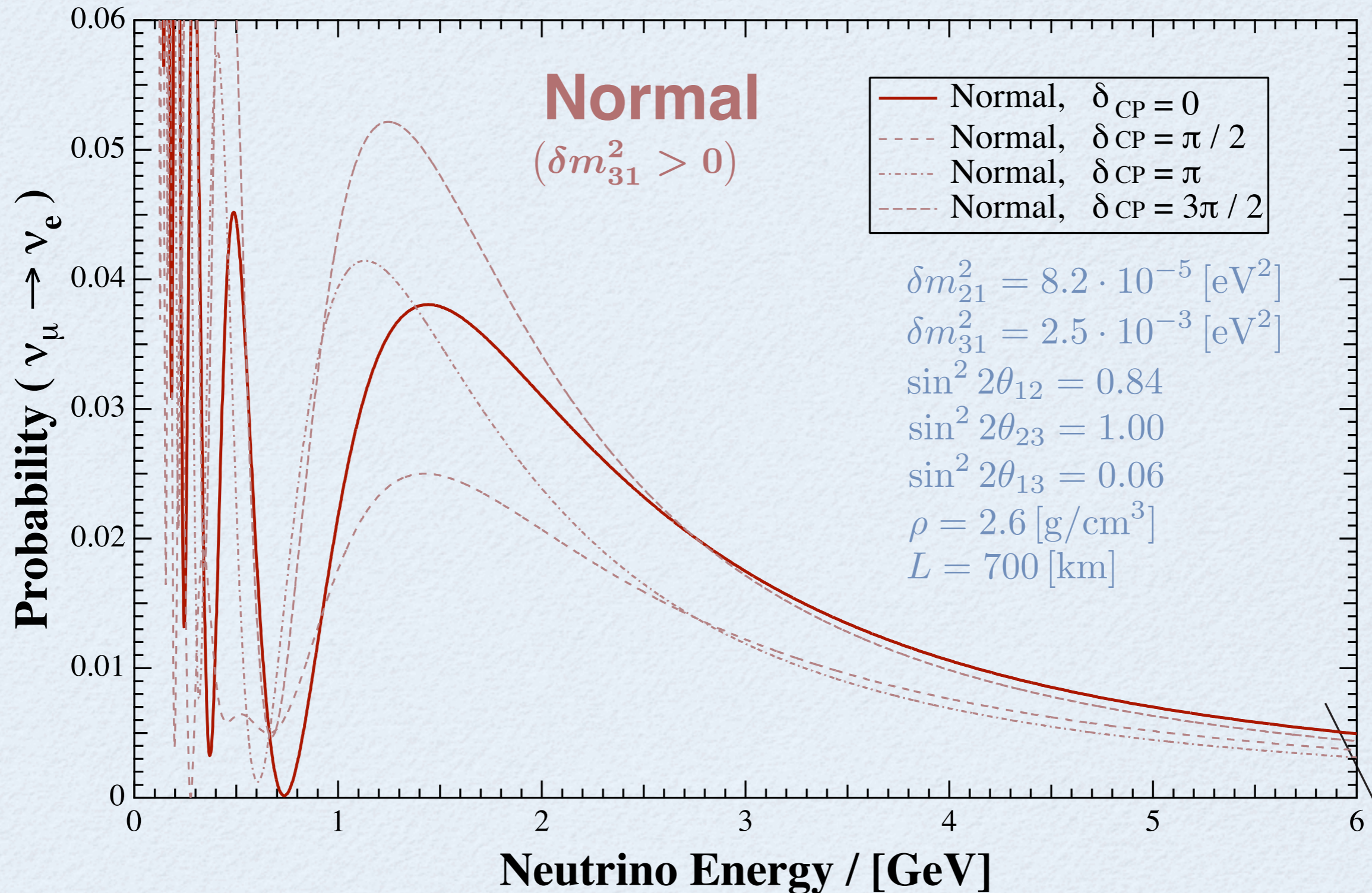


● Purpose

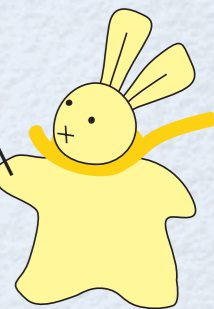
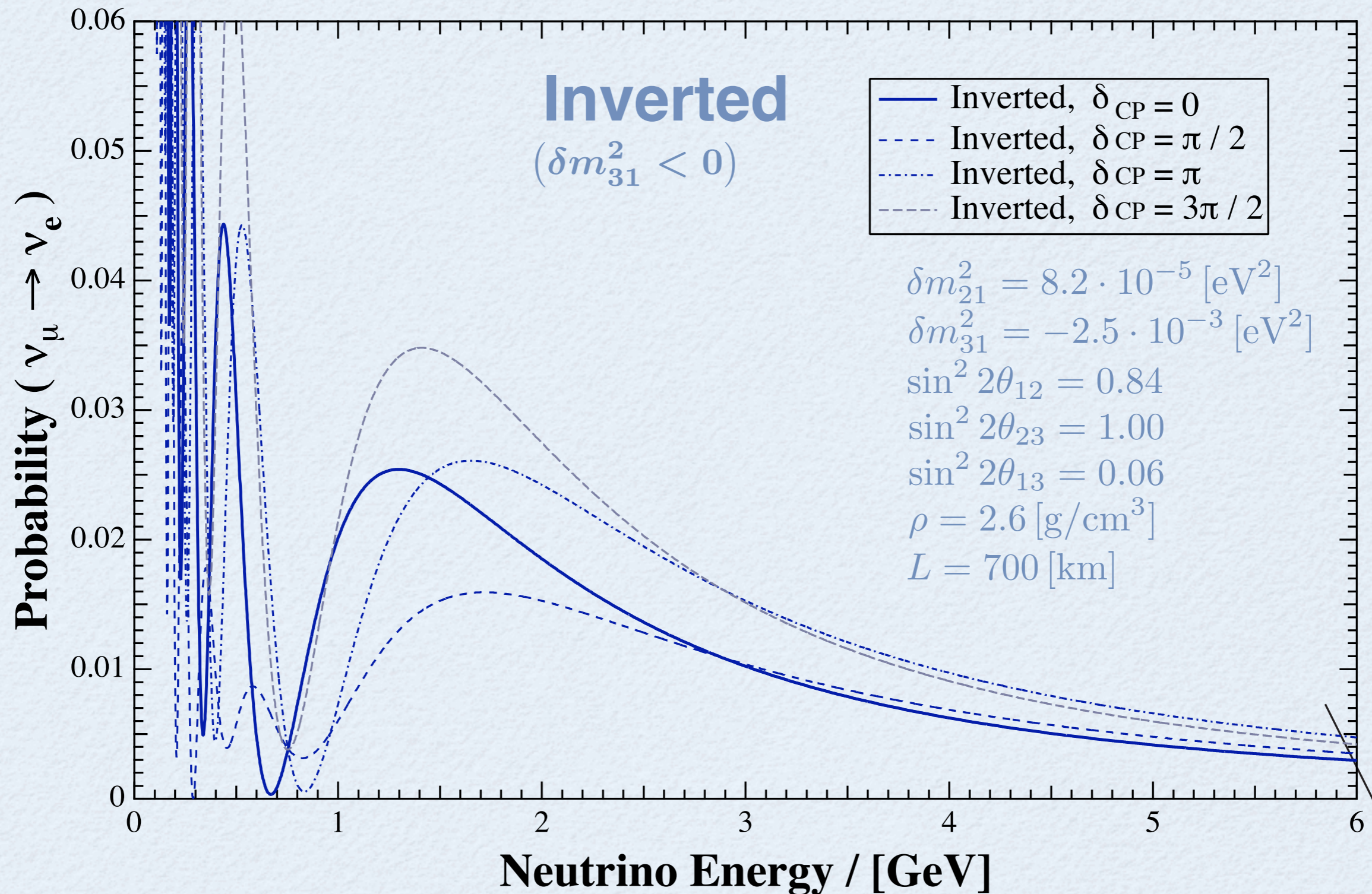
- Provide a perspective of the presence and absence of the degeneracies regarding hierarchy ($\text{sgn } \delta m_{31}^2$) and δ_{CP} .
- Discuss the way to get out of this degeneracy.



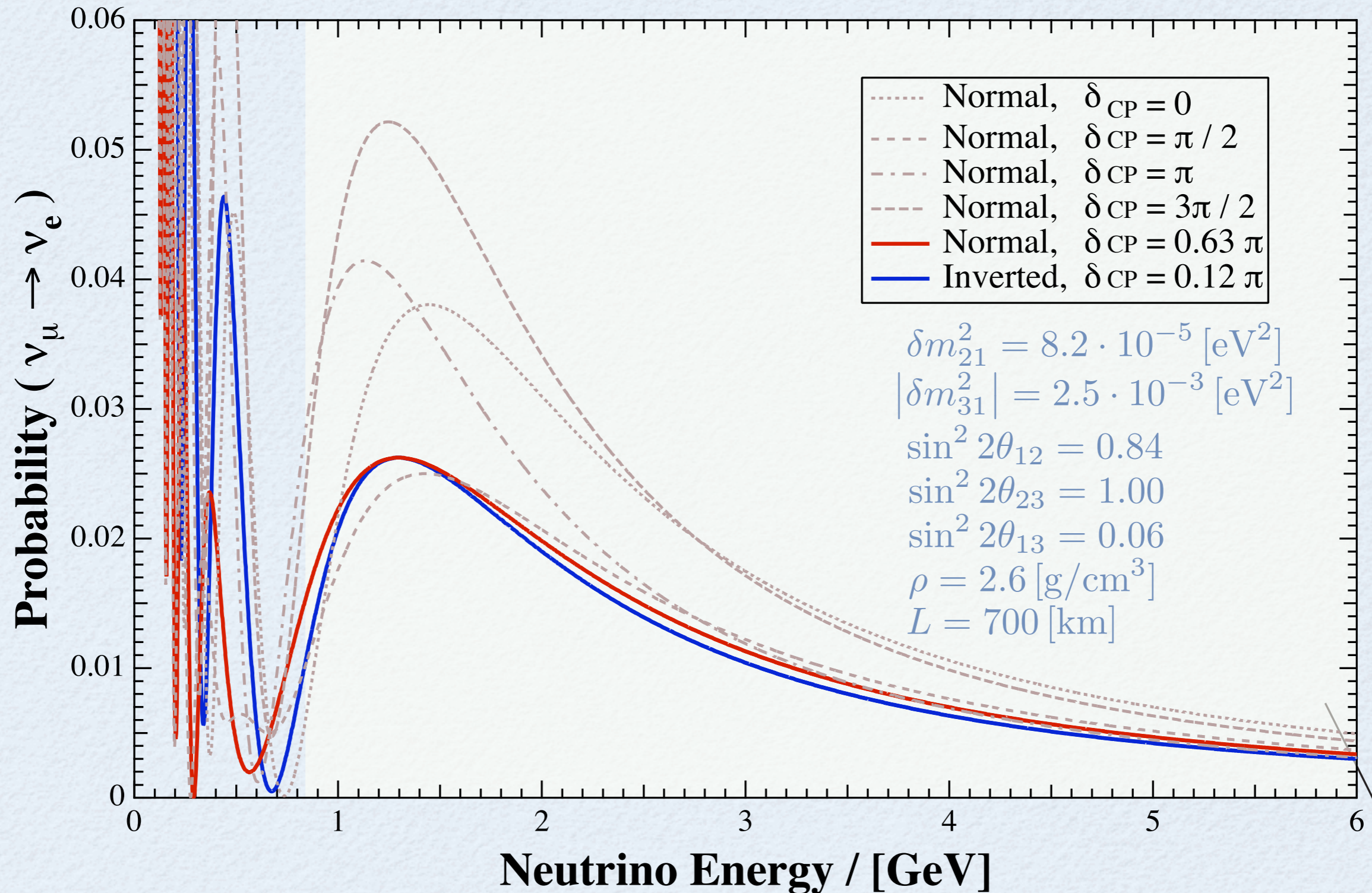
Oscillation Probabilities



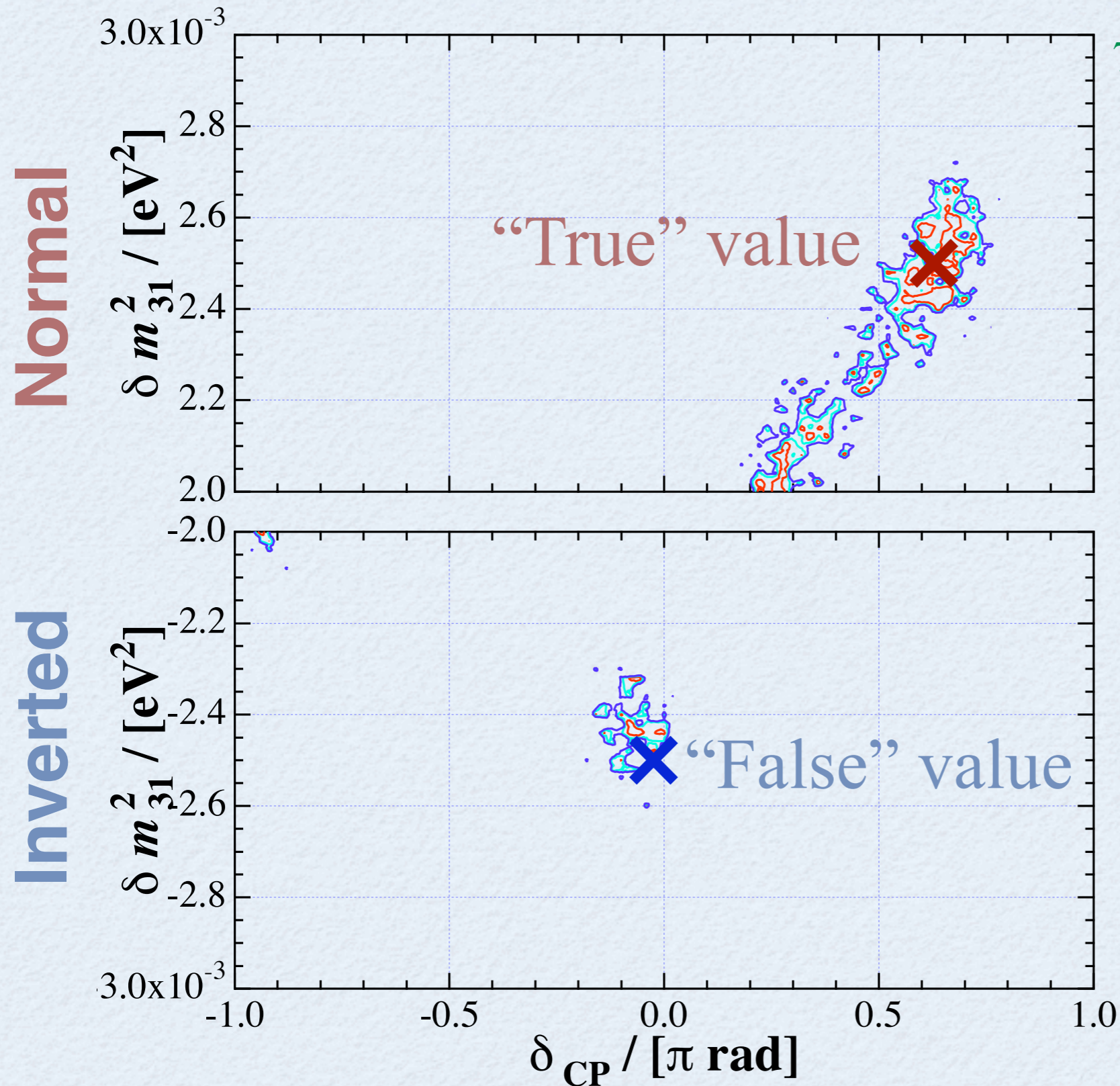
Oscillation Probabilities



Oscillation Probabilities



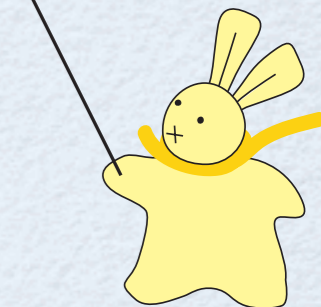
Distinguishability of Hierarchies



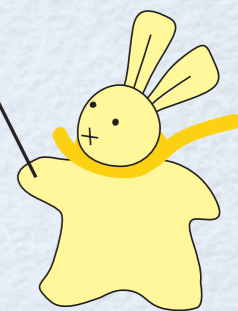
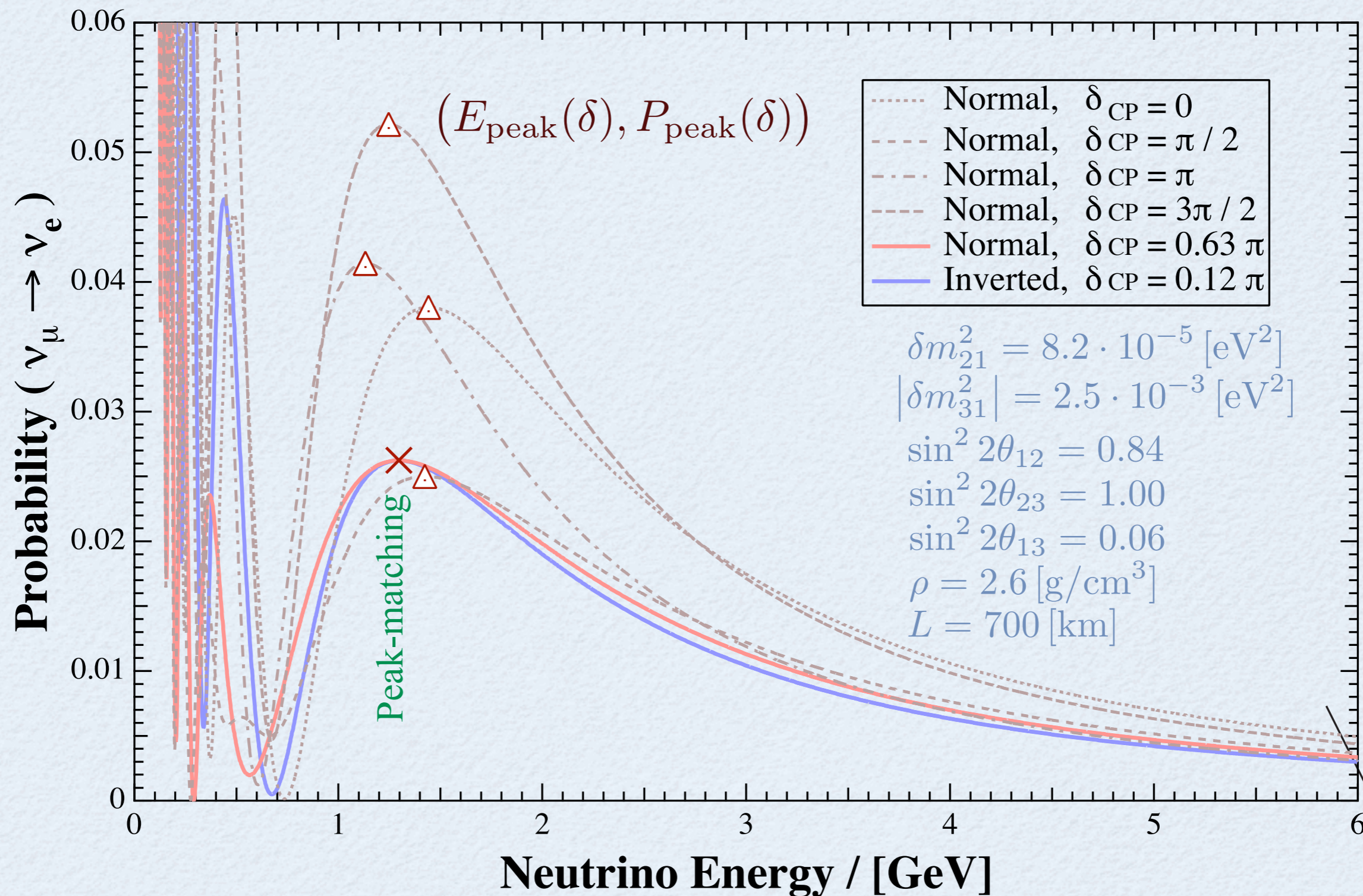
700 km

Beam flux taken from the AGS upgrade plan at BNL.

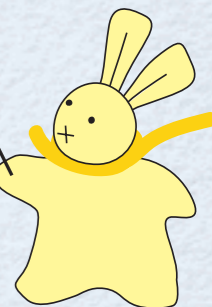
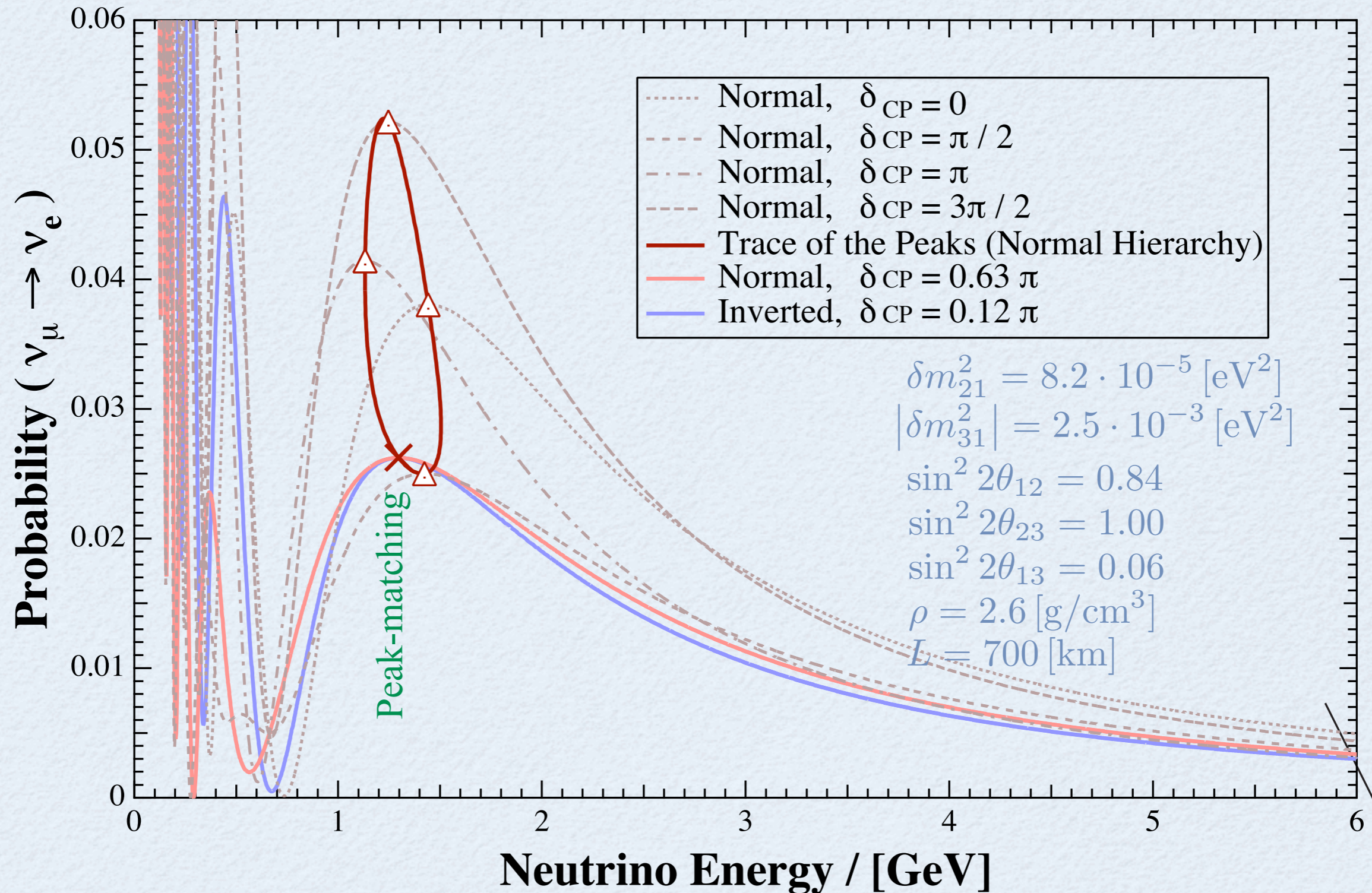
Reference: M. V. Diwan *et al.*, Phys. Rev. D **68**, 012002 (2003).



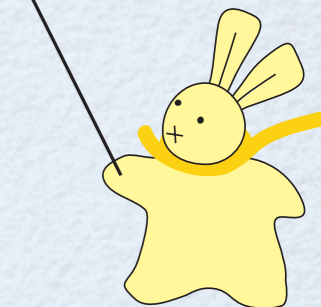
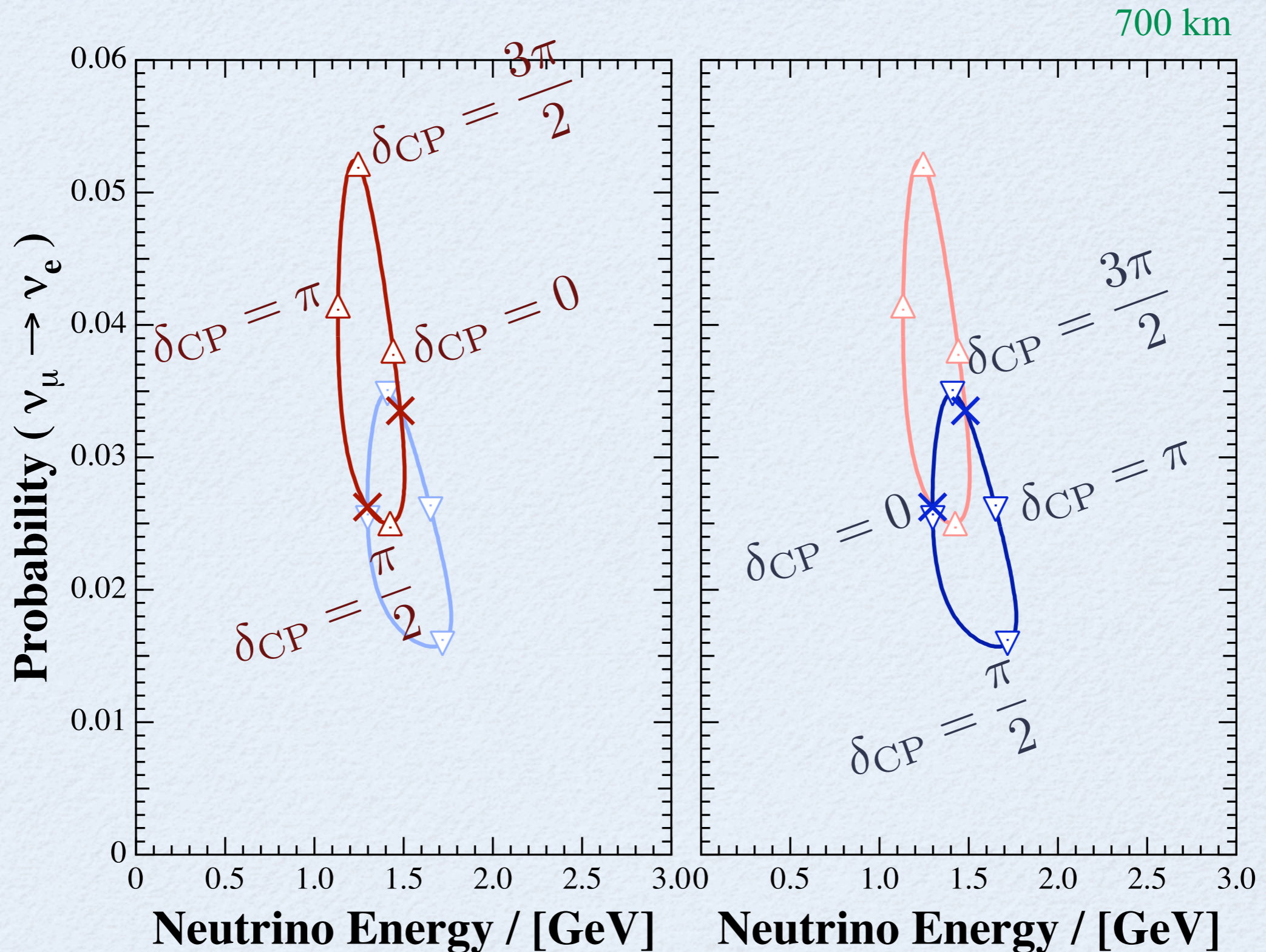
Oscillation Peaks



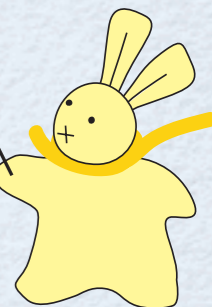
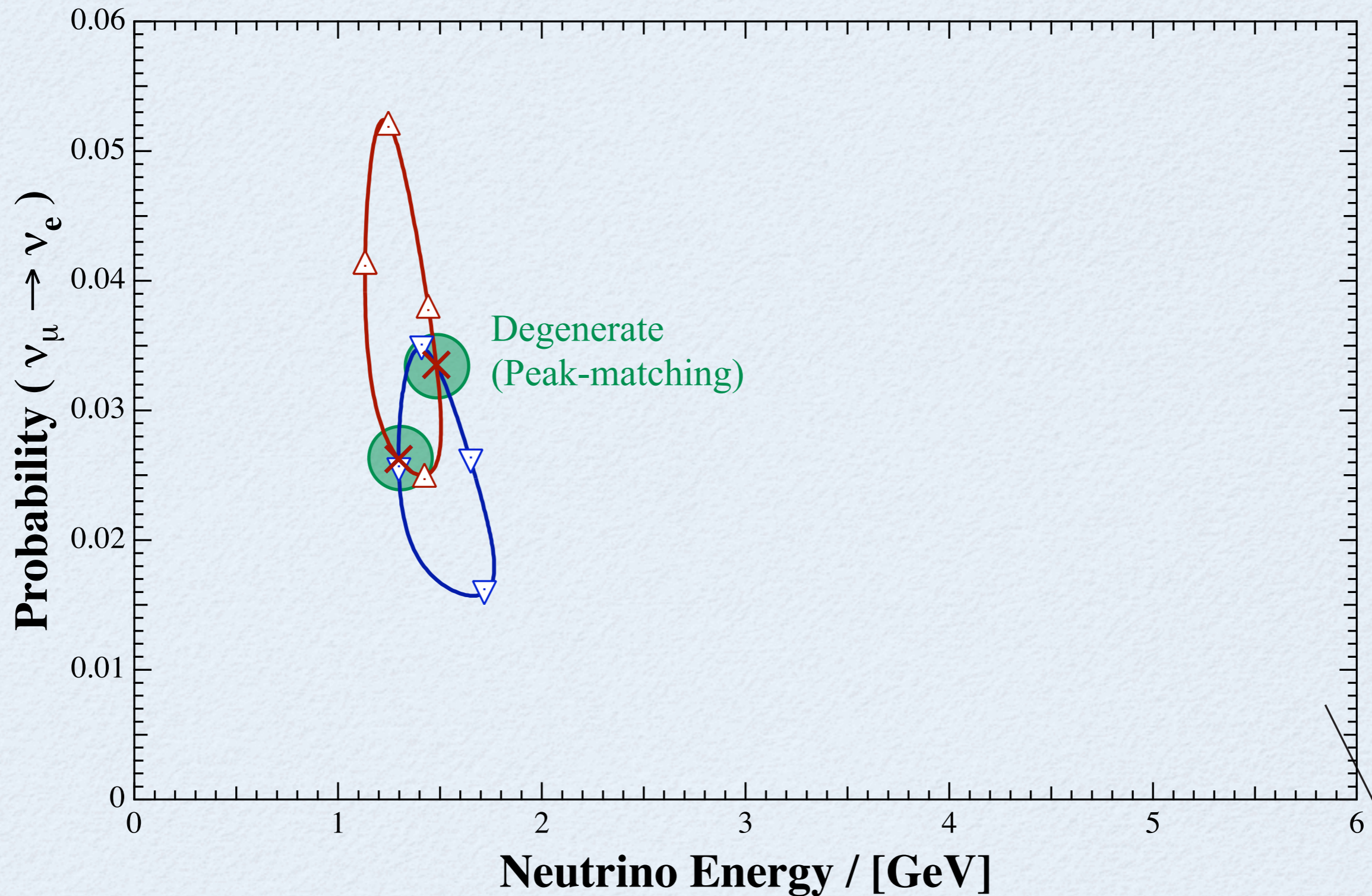
Ring of the Peaks



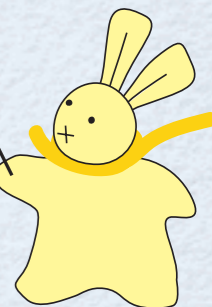
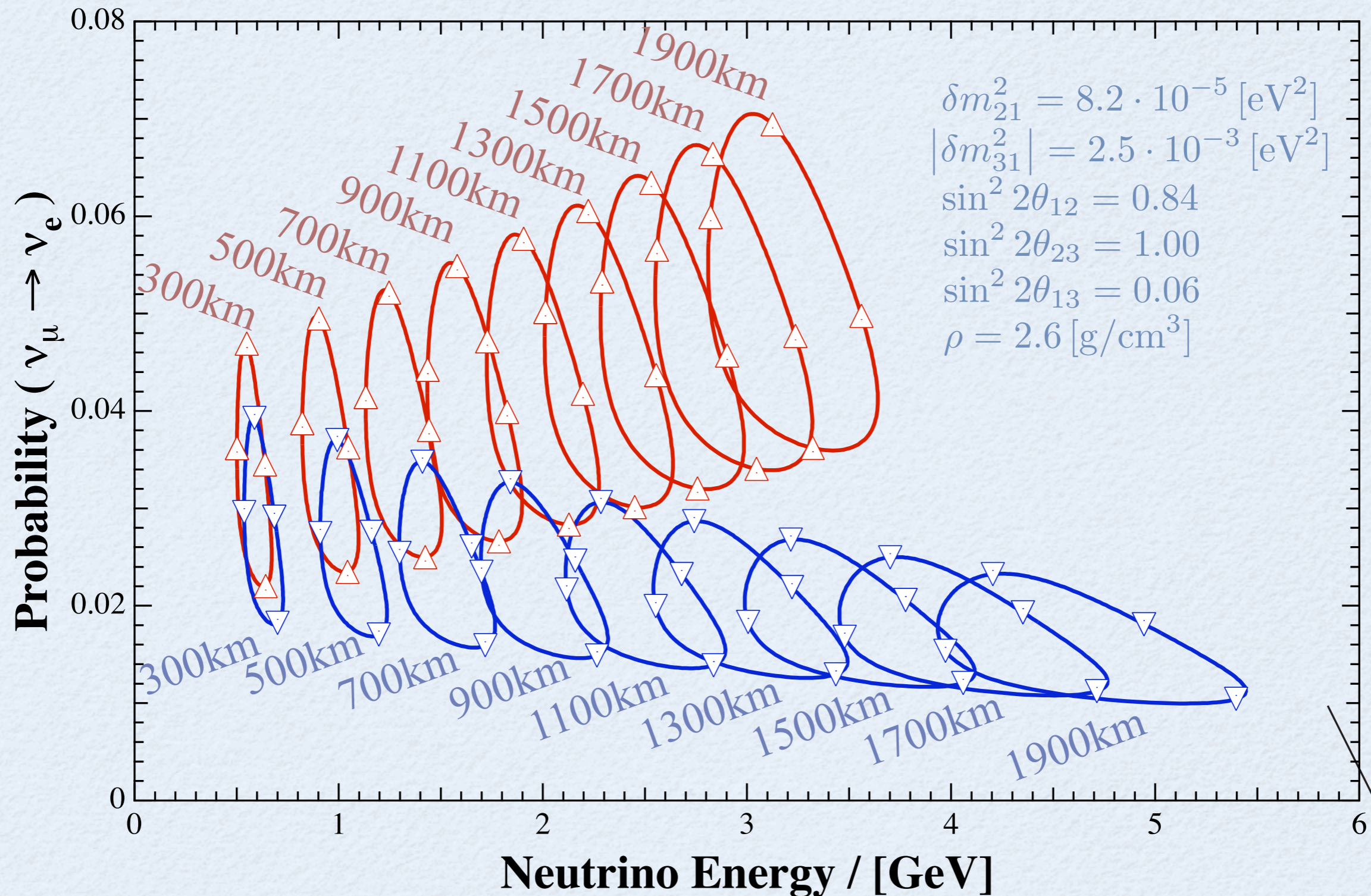
Movement on the Rings



Crosspoints of Rings



The Road of the Rings

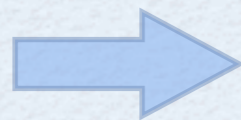


Analytic Expressions

- AKS (Arafune-Koike-Sato) approximation

J. Arafune, MK, J. Sato (1997)

$$i \frac{d\nu}{dt} = H\nu = (H_0 + H_1)\nu$$



$$\nu(x) = S(x)\nu(0)$$

$$S(x) = e^{-iH_0x} \text{Texp} \left[-i \int_0^x ds H_1^{(I)}(s) \right] \quad \left(H_1^{(I)}(x) \equiv e^{iH_0x} H_1 e^{-iH_0x} \right)$$

$$= e^{-iH_0x} - e^{-iH_0x} i \int_0^x ds H_1^{(I)}(s) + \dots$$

$$P(\nu_\alpha \rightarrow \nu_\beta) = |S_{\beta\alpha}|^2 \quad (\{\alpha, \beta\} \in \{e, \mu, \tau\})$$

- Conditions of applicability

$$(i) \delta m_{21}^2 \ll \delta m_{31}^2 ,$$

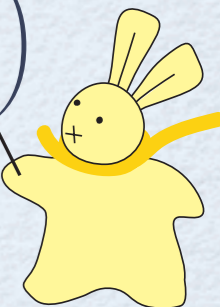
$$(ii) a \ll \delta m_{31}^2 ,$$

$$(iii) \frac{aL}{2E} \ll 1$$

$$\left(\frac{\rho}{[\text{g cm}^3]} \frac{E}{[\text{GeV}]} \ll \frac{\delta m_{31}^2}{7.56 \cdot 10^{-5} [\text{eV}^2]} \right) \quad \left(\frac{\rho}{[\text{g cm}^3]} \frac{L}{[\text{km}]} \ll 5200 \right)$$



Short baseline approximation



Oscillation Probability

- AKS (Arafune-Koike-Sato) second-order approximation (omitted in part)

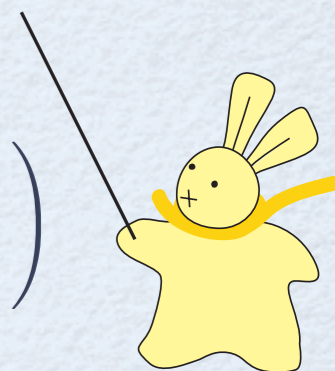
$$P(\nu_\mu \rightarrow \nu_e; E) = 4l \left[C(E) \sin^2 \Theta(E) + D(E) \right],$$

$$C(E) = 1 + 2 \frac{\Delta_m}{\Delta_{31}} (1 - 2s_{13}^2) - \Delta_{21} \frac{j}{l} \sin \delta - \Delta_{21} \frac{\Delta_m}{\Delta_{31}} \frac{j}{l} \left(\sin \delta + \frac{\Delta_{31}}{2} \cos \delta \right) + \frac{\Delta_{21}^2}{2} \left[\frac{j}{l} \cos \delta + (1 - 2s_{12}^2) \right] \frac{j}{l} \cos \delta + 3 \frac{\Delta_m^2}{\Delta_{31}^2}$$

$$\Theta(E) = \frac{\Delta_{31}}{2} - \frac{\Delta_m}{2} (1 - 2s_{13}^2) + \frac{\Delta_{21}}{2} \left(\frac{j}{l} \cos \delta - s_{12}^2 \right) - \frac{\Delta_{21}}{2} \frac{\Delta_m}{\Delta_{31}} \frac{j}{l} \left(\cos \delta + \frac{\Delta_{31}}{2} \sin \delta \right) + \frac{\Delta_{21}^2}{2} \left[\frac{j}{l} \cos \delta + \frac{1}{2} (1 - 2s_{12}^2) \right] \frac{j}{l} \sin \delta$$

$$D(E) = \frac{\Delta_{21}^2}{4} \frac{j^2}{l^2} \sin^2 \delta$$

$$\left(\Delta_{ij} = \frac{\delta m_{ij}^2 L}{2E}, \quad \Delta_m = \frac{aL}{2E} \quad l = c_{13}^2 s_{13}^2 s_{23}^2, \quad j = c_{13}^2 s_{13} c_{23} s_{23} c_{12} s_{12} \right)$$



Peak of the Oscillation Probability

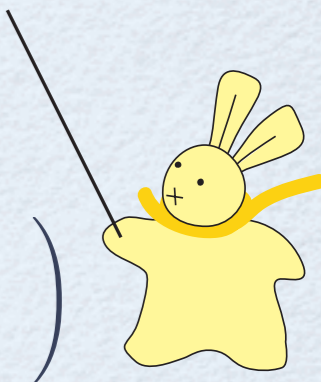
- AKS (Arafune-Koike-Sato) second-order approximation (omitted in part)

$$E_{\text{peak},n} = \frac{|\delta m_{31}^2|L}{2\Pi} \left\{ 1 \mp \frac{\Delta_m}{\Pi} (1 - 2s_{13}^2) \left(1 - \frac{4}{\Pi^2}\right) \mp R s_{12}^2 \pm R \frac{j}{l} \left(\cos \delta \pm \frac{2}{\Pi} \sin \delta\right) \right. \\ \mp \frac{\Delta_m}{2} R \frac{j}{l} \left[\left(1 + \frac{8}{\Pi^2} - \frac{64}{\Pi^4}\right) \sin \delta \pm \frac{2}{\Pi} \left(1 - \frac{8}{\Pi^2}\right) \cos \delta \right] \\ \pm R^2 \frac{\Pi}{2} \frac{j}{l} (1 - 2s_{12}^2) \left(\sin \delta \mp \frac{4}{\Pi} \cos \delta\right) + \frac{\Delta_m^2}{\Pi^2} \left(1 - \frac{12}{\Pi^2} + \frac{48}{\Pi^4}\right) \\ \left. + R^2 \frac{j^2}{l^2} \left(\pm \Pi \cos \delta \sin \delta + 1 - 3 \cos^2 \delta + \frac{4}{\Pi^2} \sin^2 \delta\right) \right\}$$

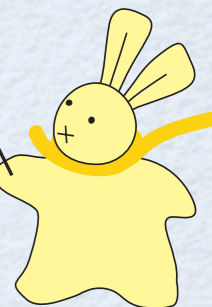
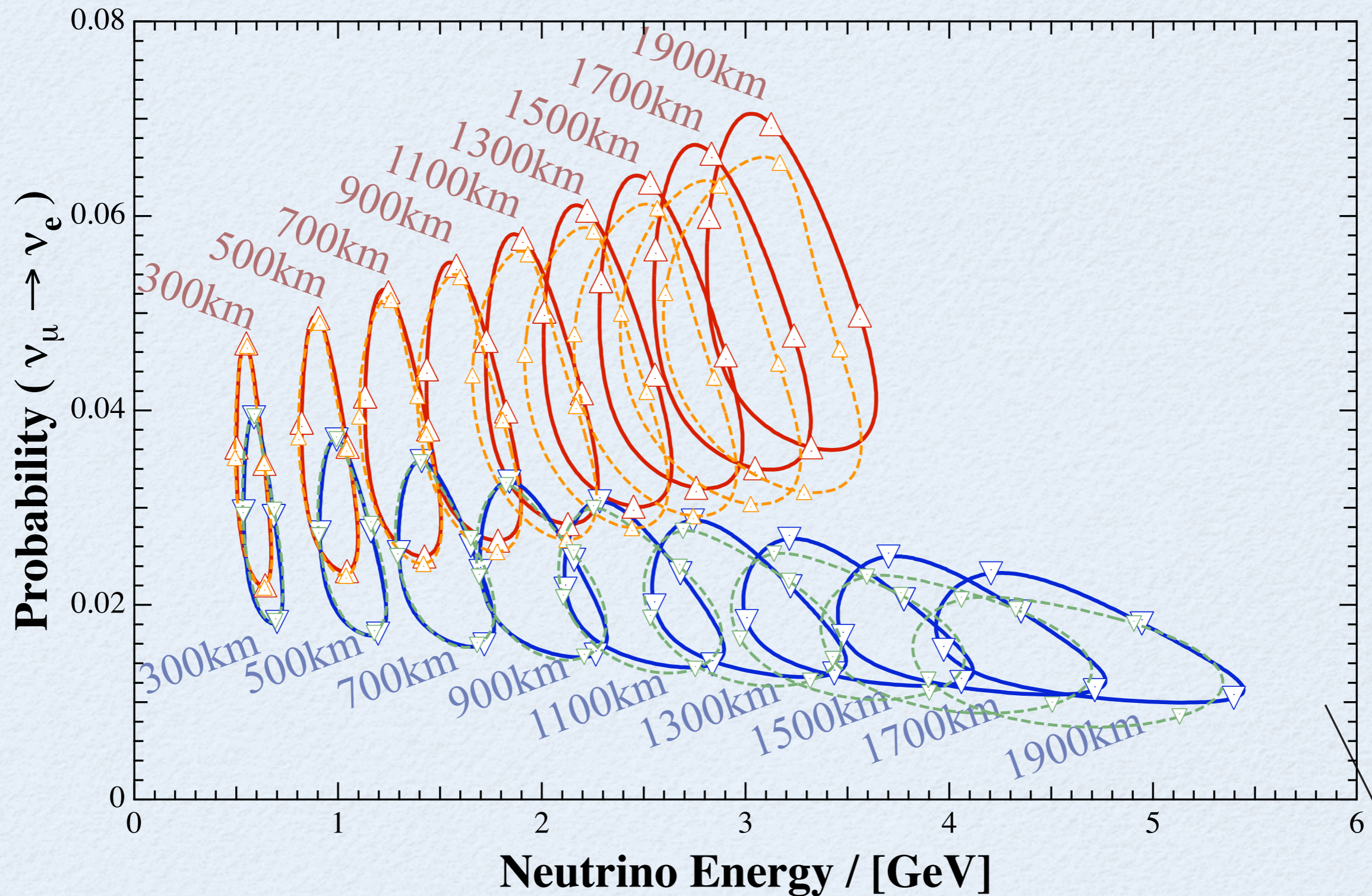
$$P_{\text{peak},n} \equiv P(\nu_\mu \rightarrow \nu_e, E_{\text{peak},n})$$

$$= 4l \left\{ 1 \pm 2 \frac{\Delta_m}{\Pi} (1 - 2s_{13}^2) - R \Pi \frac{j}{l} \sin \delta - \frac{\Pi}{2} \Delta_m R \frac{j}{l} \left[\left(1 - \frac{4}{\Pi^2}\right) \cos \delta \mp \frac{4}{\Pi} \left(1 - \frac{2}{\Pi^2}\right) \sin \delta \right] \right. \\ \left. + R^2 \frac{\Pi^2}{2} \frac{j}{l} \left[(1 - 2s_{12}^2) \cos \delta \mp \frac{2}{\Pi} s_{12}^2 \sin \delta \right] + \frac{\Delta_m^2}{\Pi^2} \left(1 + \frac{4}{\Pi^2}\right) \right. \\ \left. + \frac{1}{4} R^2 \Pi^2 \frac{j^2}{l^2} \left(1 + \cos^2 \delta \pm \frac{4}{\Pi} \cos \delta \sin \delta + \frac{4}{\Pi^2} \sin^2 \delta\right) \right\}$$

$$\left(R \equiv \frac{\delta m_{21}^2}{|\delta m_{31}^2|}, \quad \Pi \equiv (2n + 1)\pi \quad (n = 0, 1, 2, \dots) \right)$$



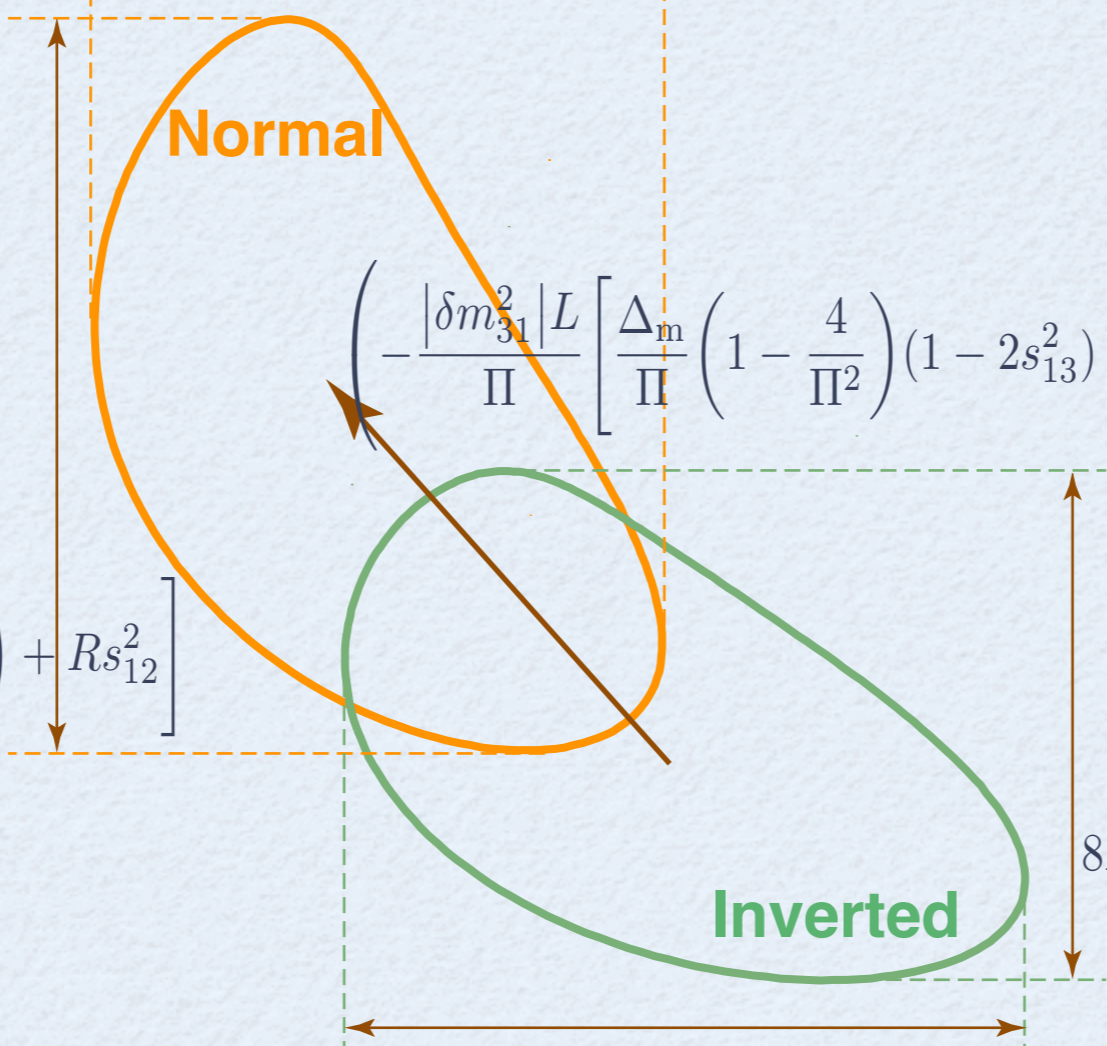
Numerical vs. Analytic



Size of Rings

$$\delta m_{21}^2 L \frac{j}{l} \left(1 - 2 \frac{\Delta_m}{\Pi} \frac{1 - \frac{32}{\Pi^4}}{1 + \frac{4}{\Pi^2}} - R \frac{1 - 2s_{12}^2}{1 + \frac{4}{\Pi^2}} \right)$$

$$\left(\begin{array}{l} R \equiv \frac{\delta m_{21}^2}{|\delta m_{31}^2|} \\ \Pi \equiv (2n + 1)\pi \quad (n = 0, 1, 2, \dots) \end{array} \right)$$



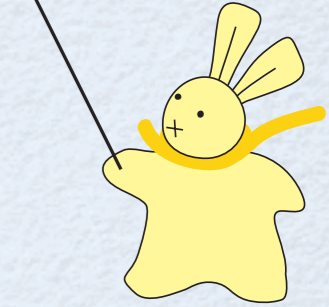
$$\left(-\frac{|\delta m_{31}^2| L}{\Pi} \left[\frac{\Delta_m}{\Pi} \left(1 - \frac{4}{\Pi^2} \right) (1 - 2s_{13}^2) + R s_{12}^2 \right], \frac{16l \Delta_m}{\Pi} (1 - 2s_{13}^2) \right)$$

$$8Rj \sqrt{1 + \frac{4}{\Pi^2}} \left[1 + 2 \frac{\Delta_m}{\Pi} \left(1 - \frac{2}{\Pi^2} \right) + R s_{12}^2 \right]$$

$$8Rj \sqrt{1 + \frac{4}{\Pi^2}} \left[1 - 2 \frac{\Delta_m}{\Pi} \left(1 - \frac{2}{\Pi^2} \right) - R s_{12}^2 \right]$$

Recall:
 $\Delta_m \propto L$

$$\delta m_{21}^2 L \frac{j}{l} \left(1 + 2 \frac{\Delta_m}{\Pi} \frac{1 - \frac{32}{\Pi^4}}{1 + \frac{4}{\Pi^2}} + R \frac{1 - 2s_{12}^2}{1 + \frac{4}{\Pi^2}} \right)$$



Baseline Length of Ring Separation

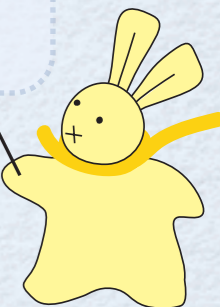
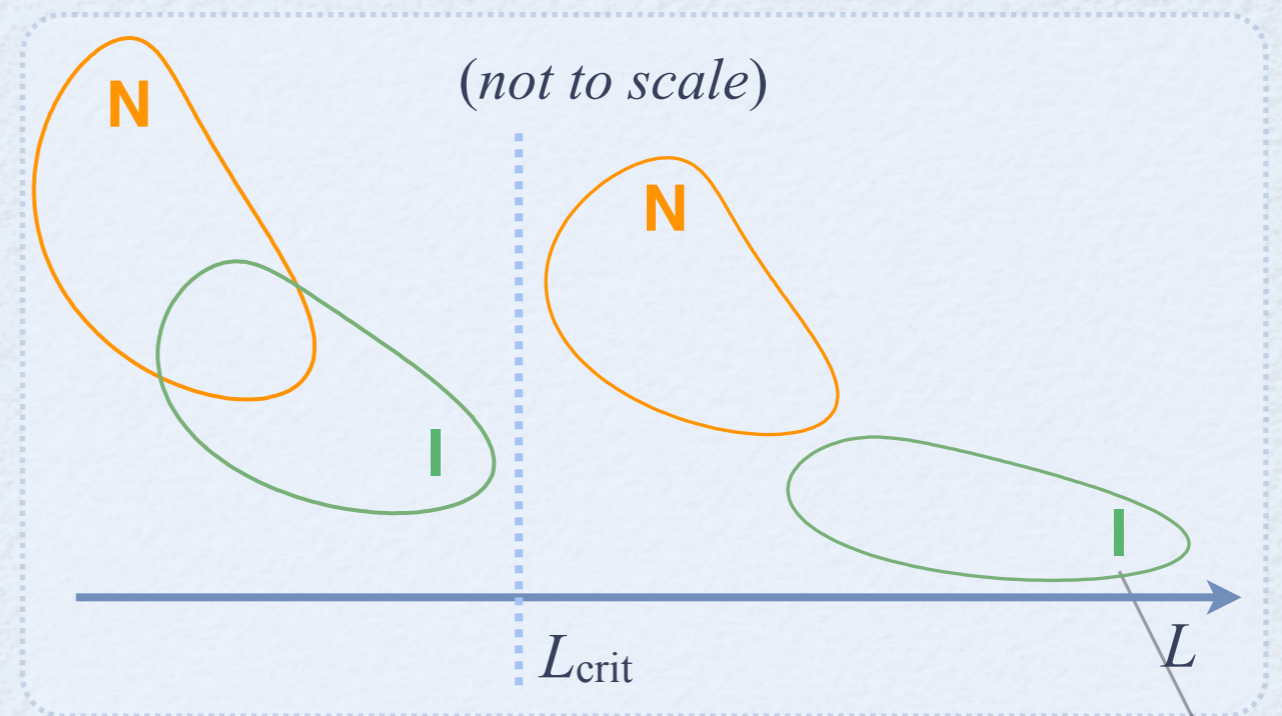
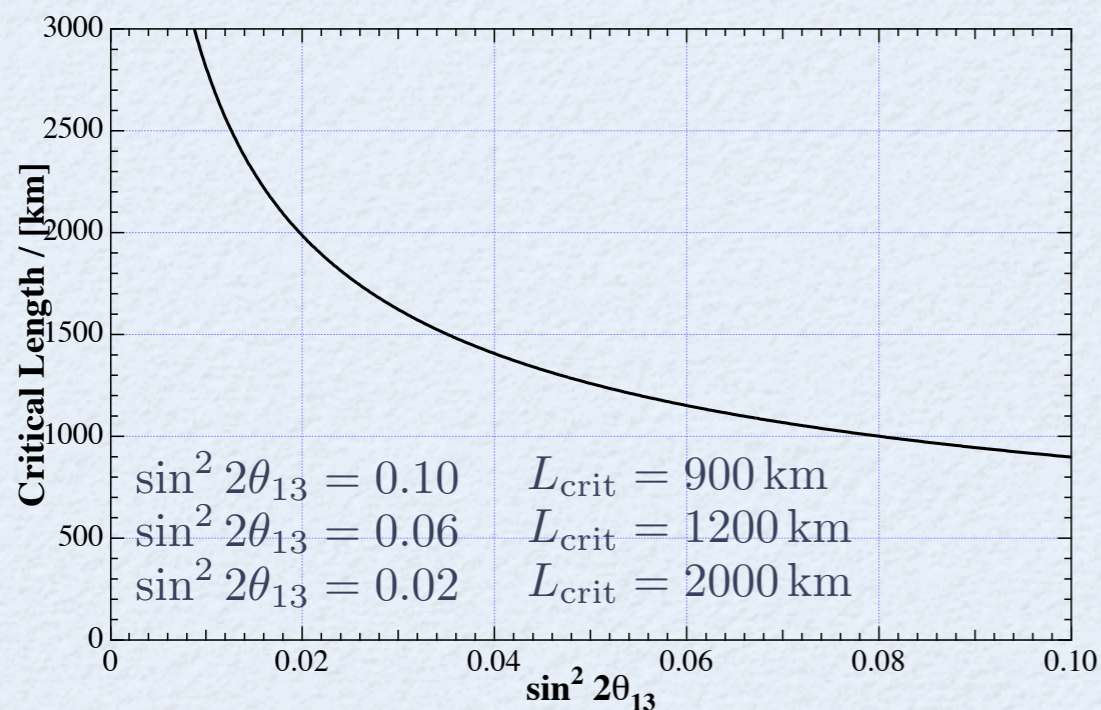
- Rings, and thus hierarchies, separate at a long baseline $L > L_{\text{crit}}$

$$L_{\text{crit}} = \frac{1}{a'} \frac{\delta m_{21}^2}{|\delta m_{31}^2|} \frac{1}{1 - 2s_{13}^2} \frac{\Pi}{1 - \frac{12}{\Pi^2} + \frac{64}{\Pi^4}} \left[-\left(1 - \frac{8}{\Pi^2}\right) s_{12}^2 + \sqrt{\left(1 - \frac{12}{\Pi^2} + \frac{64}{\Pi^4}\right) \frac{c_{23}^2 c_{12}^2 s_{12}^2}{s_{23}^2 s_{13}^2} - \frac{4}{\Pi^2} s_{12}^4} \right]$$

$$\approx \frac{1}{a'} \frac{\delta m_{21}^2}{|\delta m_{31}^2|} \frac{\Pi}{\sqrt{1 - \frac{12}{\Pi^2} + \frac{64}{\Pi^4}}} \frac{c_{23} c_{12} s_{12}}{s_{23} s_{13} (1 - 2s_{13}^2)} \sim \frac{1}{s_{13}} \quad (\text{up to first order})$$

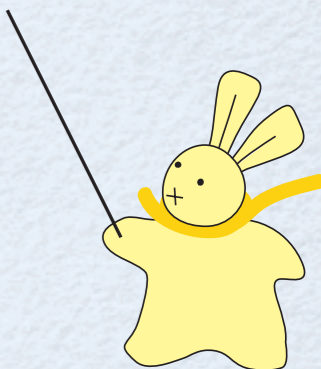
$$\left(\frac{1}{a'} \equiv \frac{1}{\sqrt{2} G_F n_e} = \frac{5.17 \cdot 10^3 \text{ [km]}}{\frac{\rho}{[\text{g cm}^{-3}]}} \right)$$

For our example parameter set,

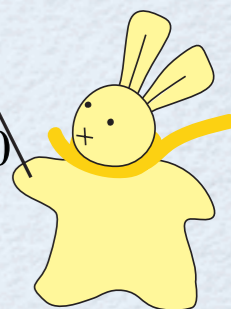
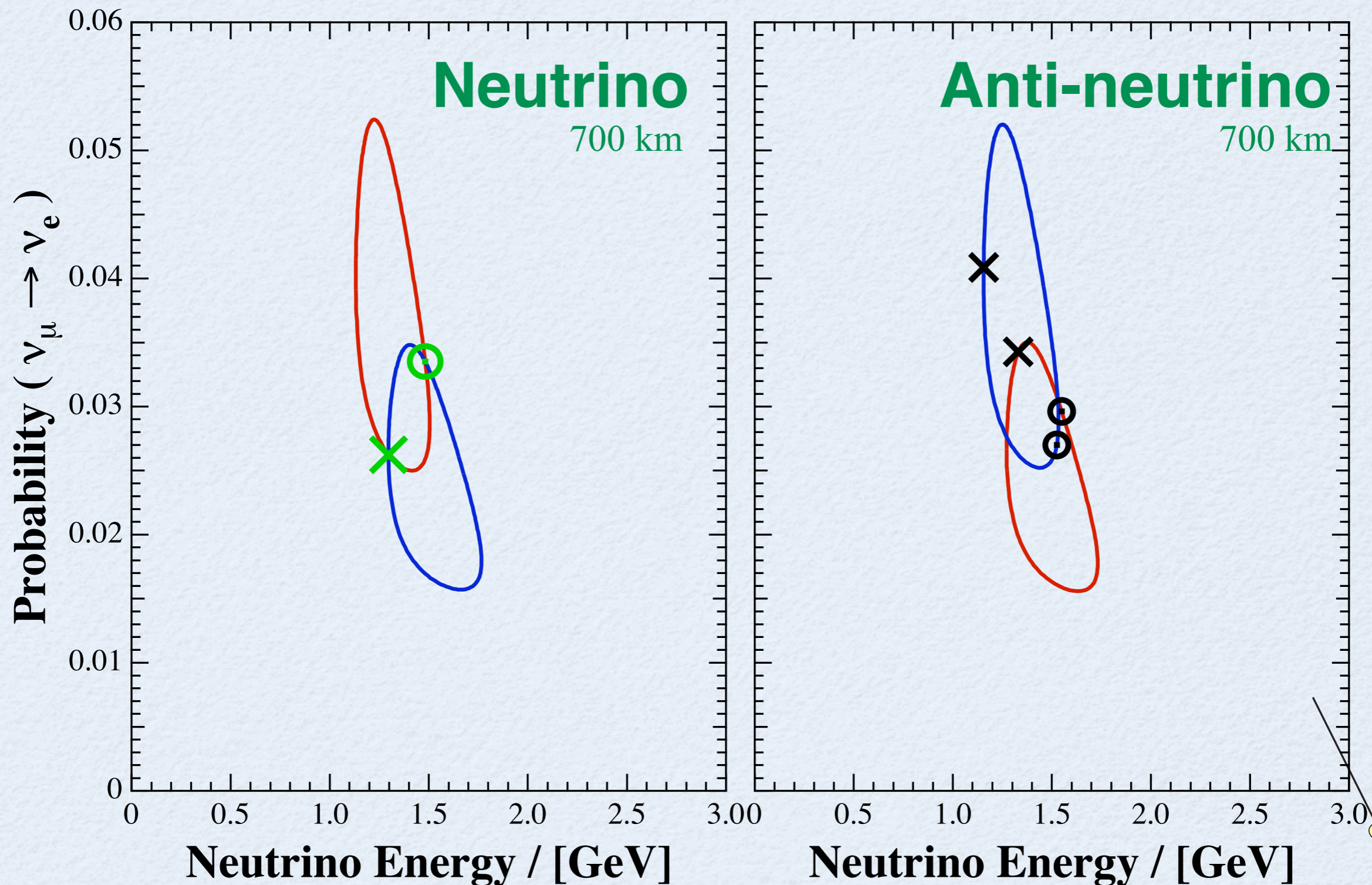


Resolving the Degeneracy

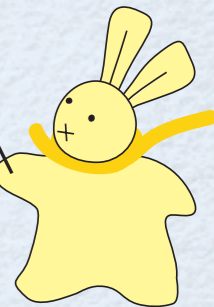
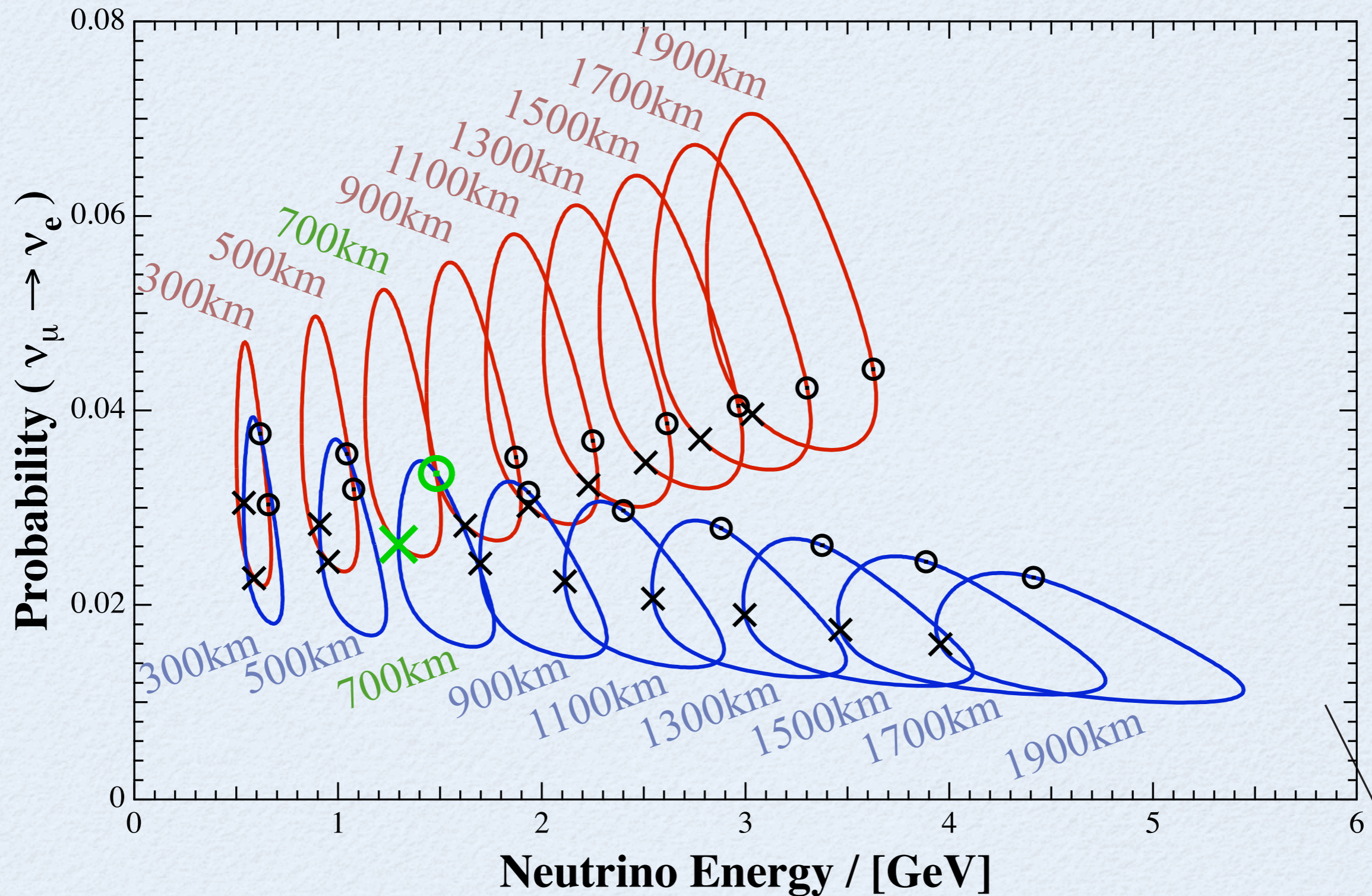
- The way out from the degeneracy (we are discussing)?
 - Go for a long-length baseline, $L > L_{\text{crit}}$.
 - Employ anti-neutrino beams together with neutrino beams.
 - Combine two different baseline lengths. (“The Two Towers”)
 - Push the detection to the lower-energy neutrinos.



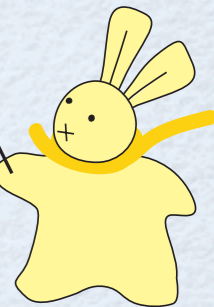
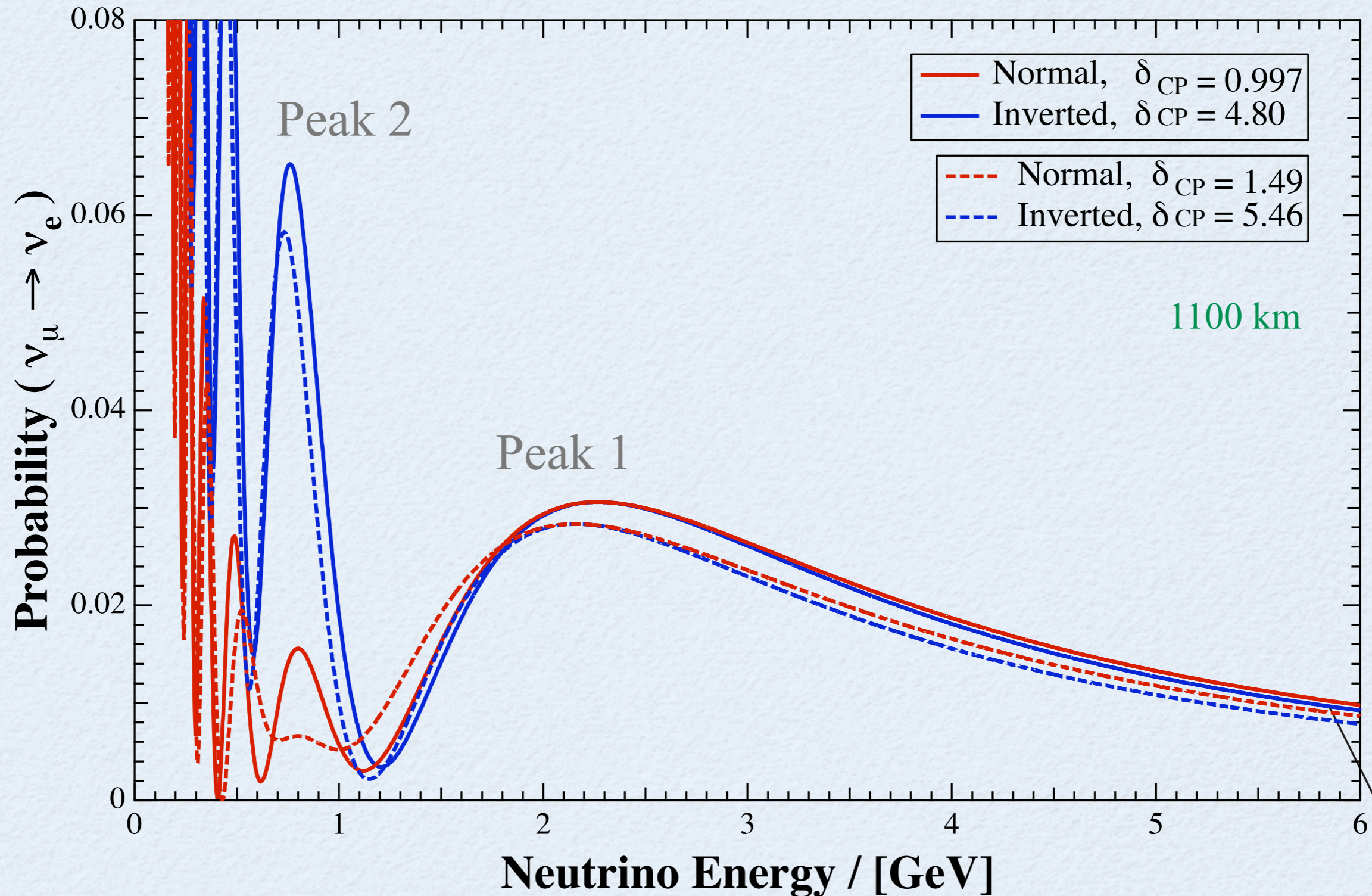
Use of Anti-neutrinos



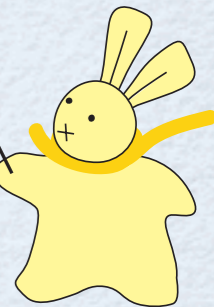
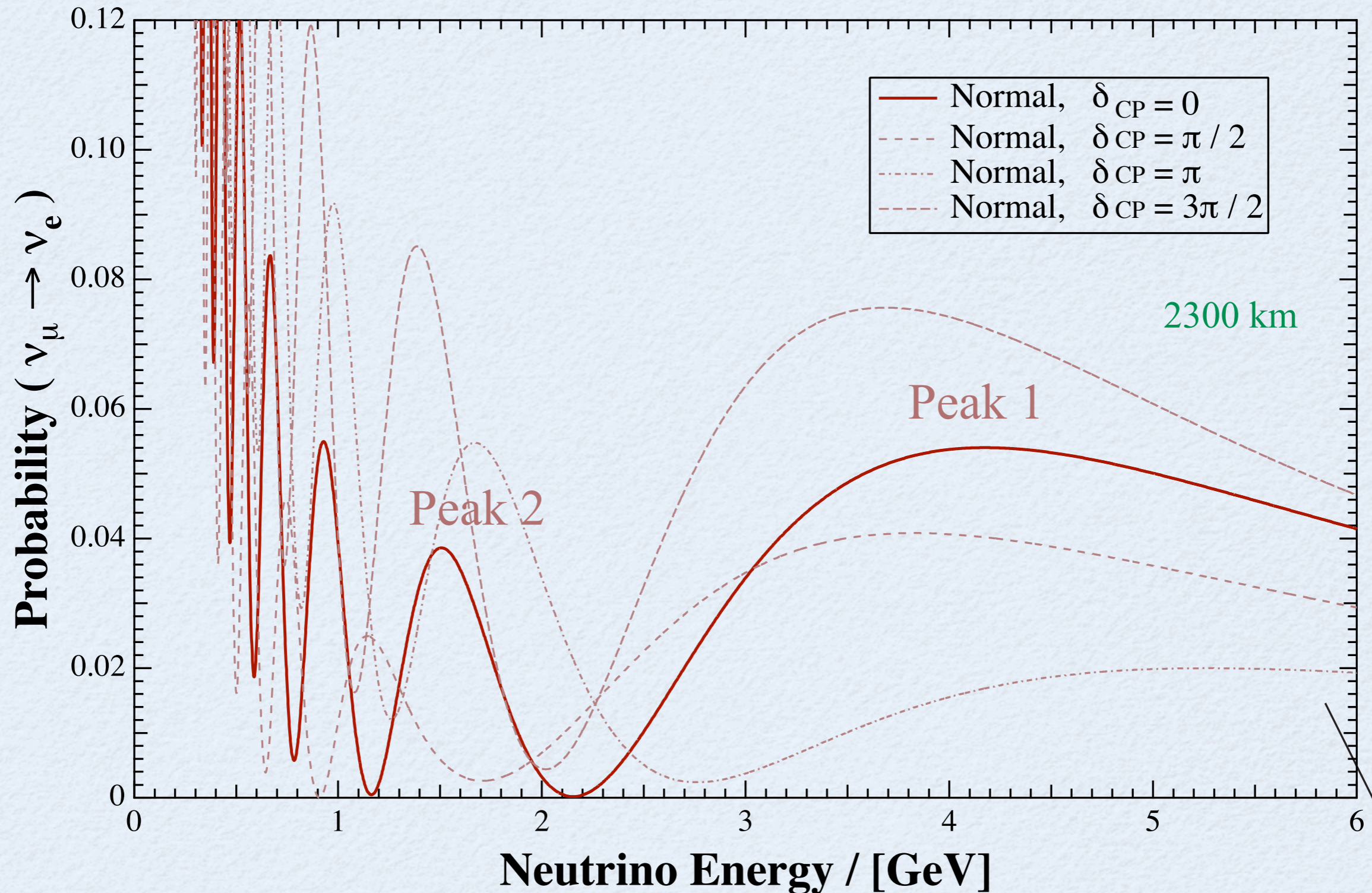
Use of Another Distance



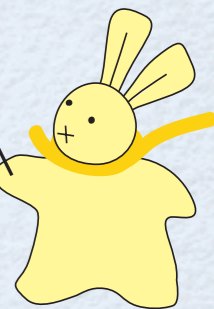
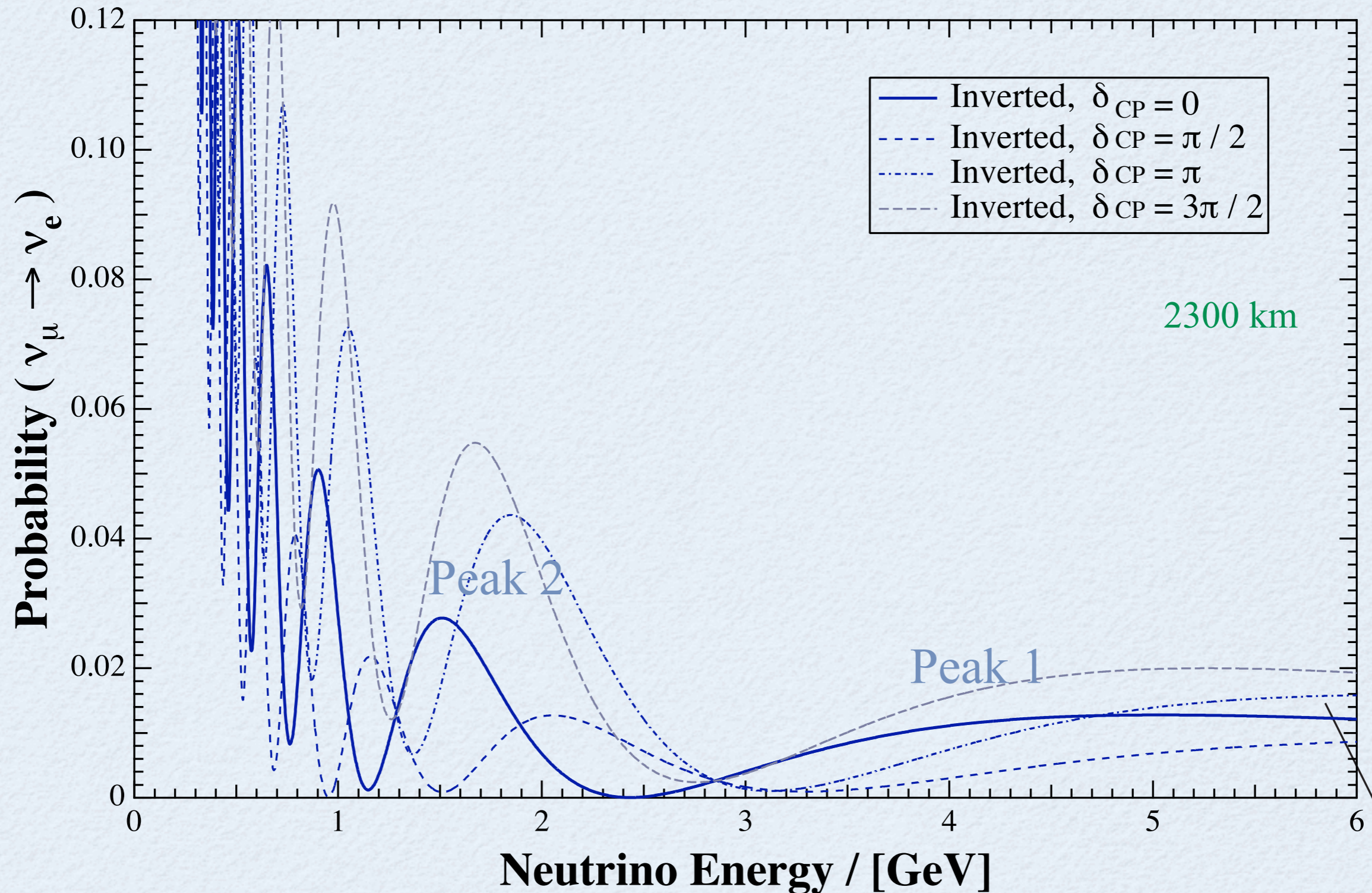
Use of Another Peak



Use of Another Peak



Use of Another Peak



Conclusions, Outlooks

- Determination of neutrino parameters may be complicated due to the degeneracy.
- The peak of the oscillation is a good representative of the whole spectra.
 - Peak-matching leads to the mutually “similar” oscillation spectra.
- The analysis of the peak position provides a perspective of the presence and absence of parameter degeneracies in the long baseline experiments.
- The parameter-searching power can be systematically analyzed.
 - The road of the rings: Various baseline length
 - Blurred rings (*to be done*): Ambiguities of the oscillation parameters
 - Another road of the rings (*to be done*): Combination of neutrinos and anti-neutrinos

