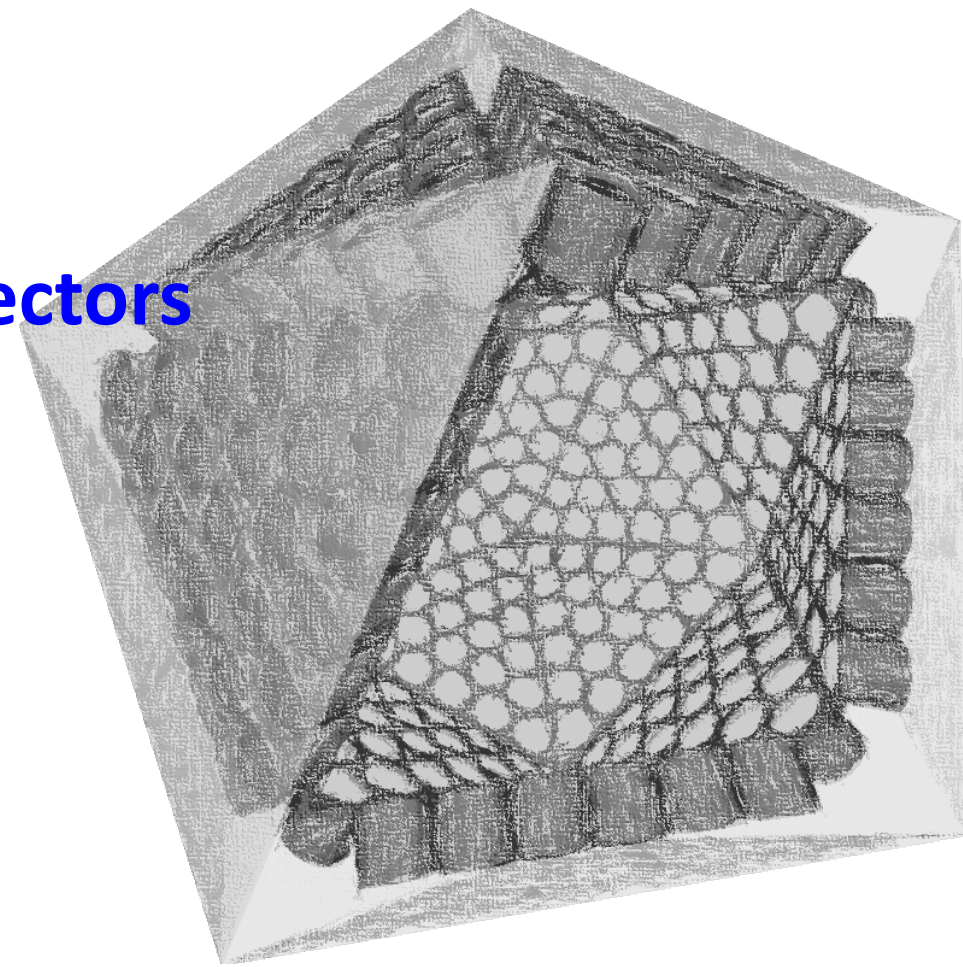


Distributed Imaging for Scintillation Detectors

(of neutrinos and all sorts of other things)

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Most of the work done by J.Dalmasson, A.Jamil, S.Kravitz,
M.Malek, K.Wells, J.Bentley, S.Steven, J.Su



The evil



Joseph Liouville

Cassandra



"our John"

It is a frustration of mine that we are still using ~the same organic scintillators developed in the 1920's, readout by vacuum tubes also developed in the 1920's. Somehow our field has not been reached by modern advances in material science.

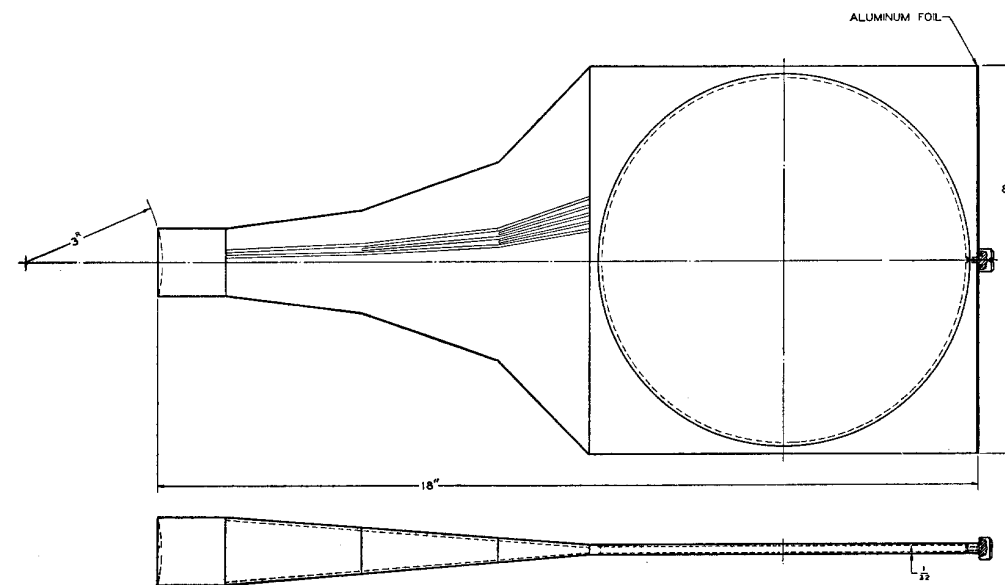
So, it is fitting that I start with a couple of old papers, generically pointed out to me by John, probably while on some boring KamLAND shift...



R. L. Garwin "The Design of Liquid Scintillation Cells"
Rev Sci Inst 23, 755 (1952)

R. L. Garwin "The Collection of Light from Scintillation
Counters" Rev Sci Inst 31, 1010 (1960)

These papers, basically, mark the beginning of "adiabatic lightguides" and can be summarized by saying that evil M. Liouville is the origin of lots of trouble.



Applied to large scintillation detectors, this principle puts serious constraints on the possibility of imaging tracks. This is because there are very few photons produced by a scintillator* and these photons are produced with an emittance

$$\eta = V \times 4\pi$$

where V is the entire volume of the detector and 4π is the full solid angle.

Quibbles and notes:

- 1) Strictly speaking this is not true, for a point-like source $\eta = \varepsilon 4\pi$, because V is small, but, unfortunately, we do not know where the event is in the detector, so, in practice the emittance contains the entire volume of the detector.
- 2) The emittance also contains Δt and ΔE , or equivalent. For photons one has really only access to Δt . Indeed, Δt can be traded for other variables. An infinitely reflective sphere can extract all photons out of a tiny hole by waiting an arbitrarily long time. In practice infinitely reflective things do not exist and I am not sure it is wise to give up timing, so I will ignore this possibility.

* This is not to say scintillators are “stupid”. Typical energy conversion efficiencies are in the neighborhood of 10%, which is only one order of magnitude different from 100% and this explains why it is so hard to make better scintillators!

Hence imaging in a scintillator with a (few) camera(s) and a finite size lens is a non-starter:

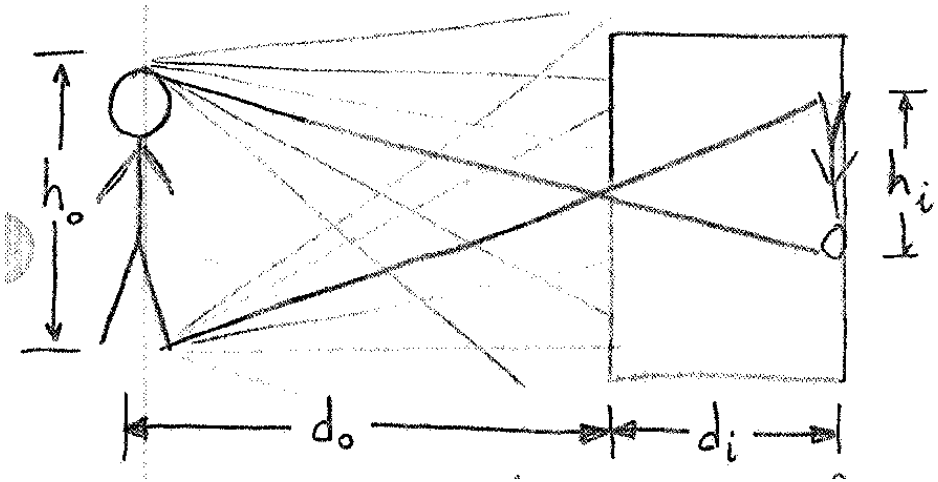
- A conventional lens arrangement selects few photons to produce images
- The smaller the aperture the more depth of field and, because of the huge emittance, we need a pinhole!

Not enough photons!

This is irrespective of the sensitivity/gain/whatever of the camera.

*Yet, lots of things could be done with imaging,
witness the many existing segmented detectors.*

Two extreme cases of “imaging”



Pinhole camera:

Pro: image infinitely sharp

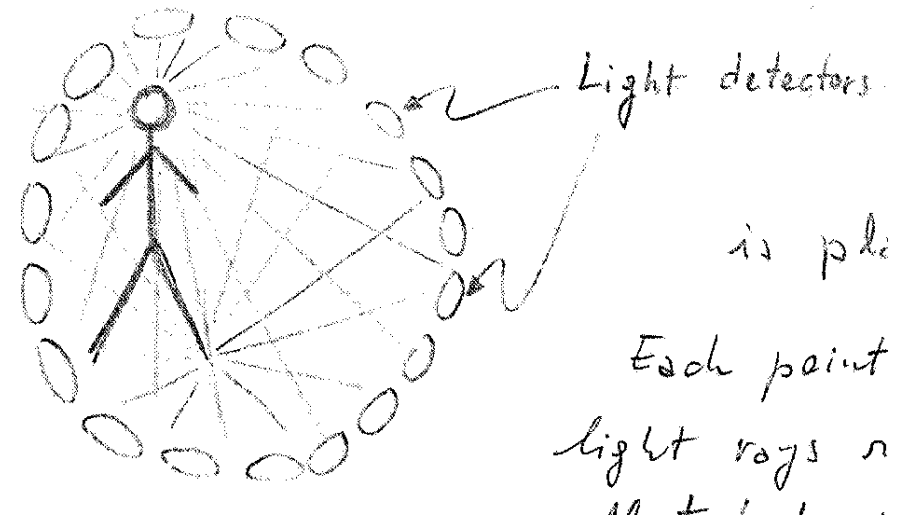
Con: infinitely inefficient

Large liquid scintillation detector:

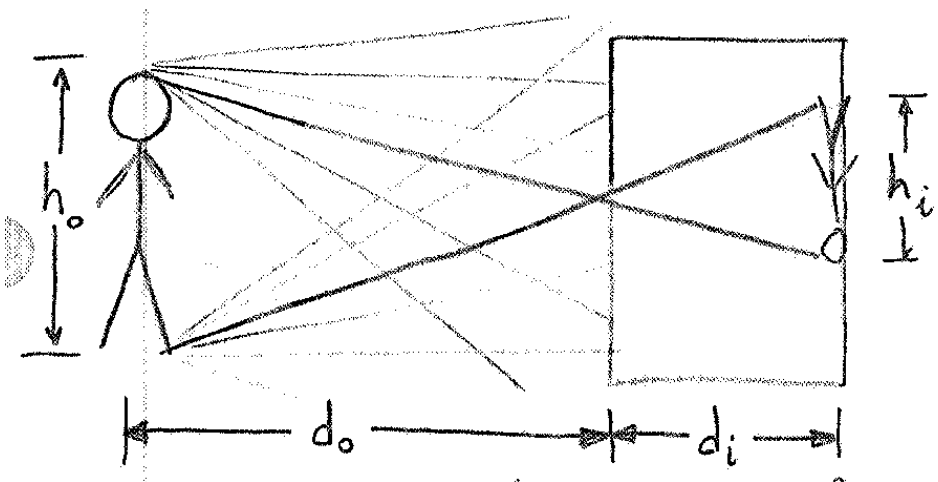
Pro: collect every single photon

Con: no imaging

Example 1: “Proximity focusing”



Two extreme cases of “imaging”



Pinhole camera:

Pro: image infinitely sharp

Con: infinitely inefficient

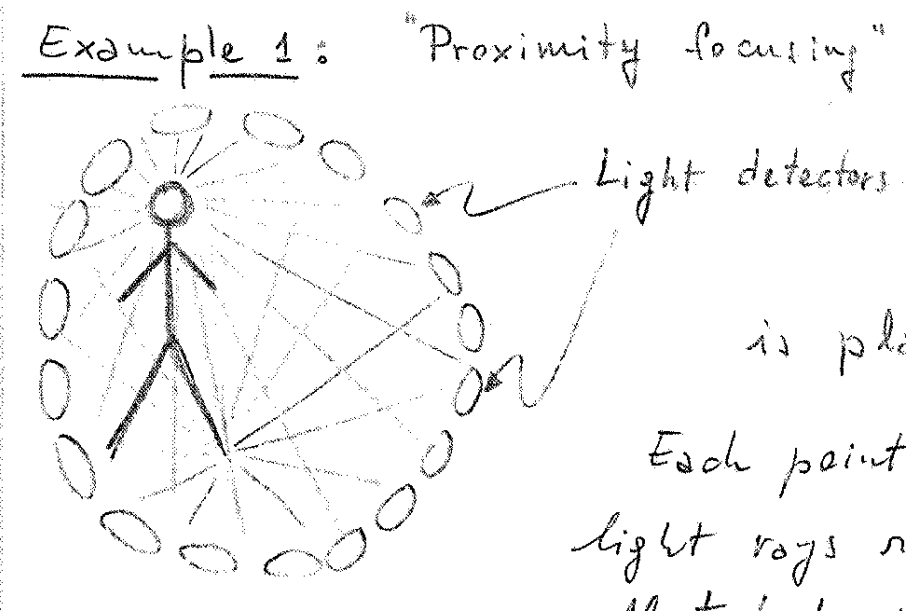
Large liquid scintillation detector:

Pro: collect every single photon

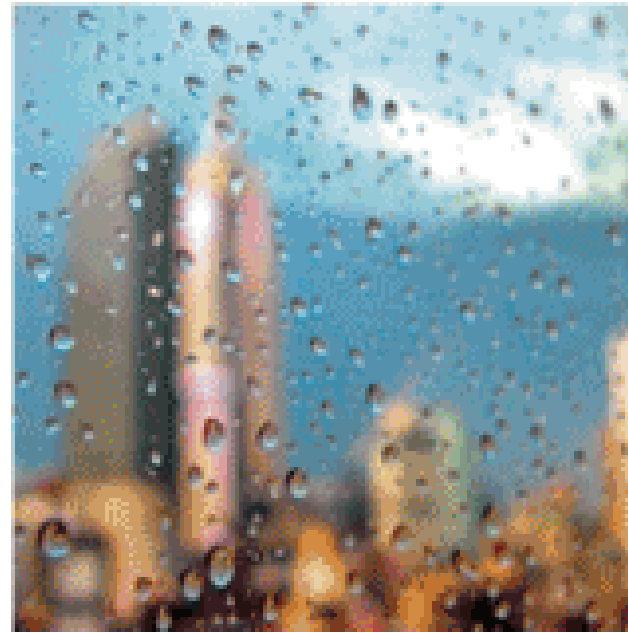
Con: no imaging

However, note that the rays hitting a particular PMT here do contain angular information.

...if we only knew how to record it!



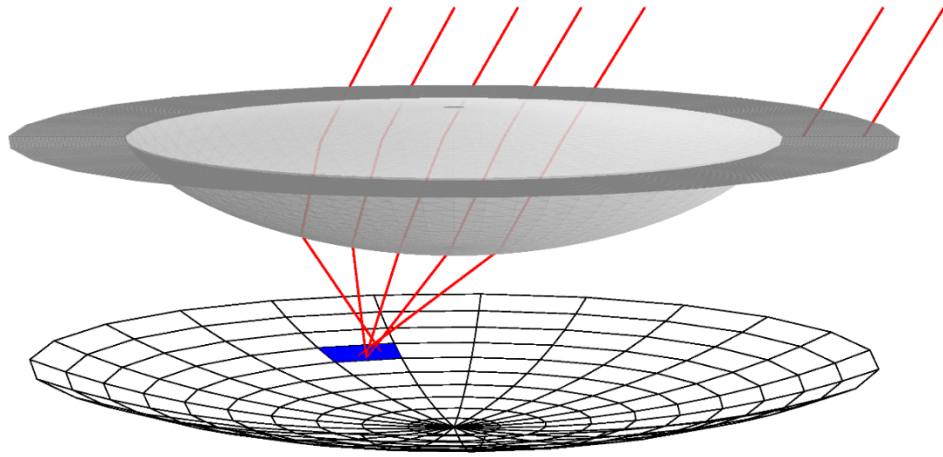
It turns out that measuring angles of incoming photons (or wavefronts) is done in some exotic (“plenoptic”) cameras [e.g. Lytro (US), Raytrix (Germany), now bought/sold/merged by the usual mindless MBAs].
This allows 3D imaging out of a single lens or “offline refocusing”.



Idea:

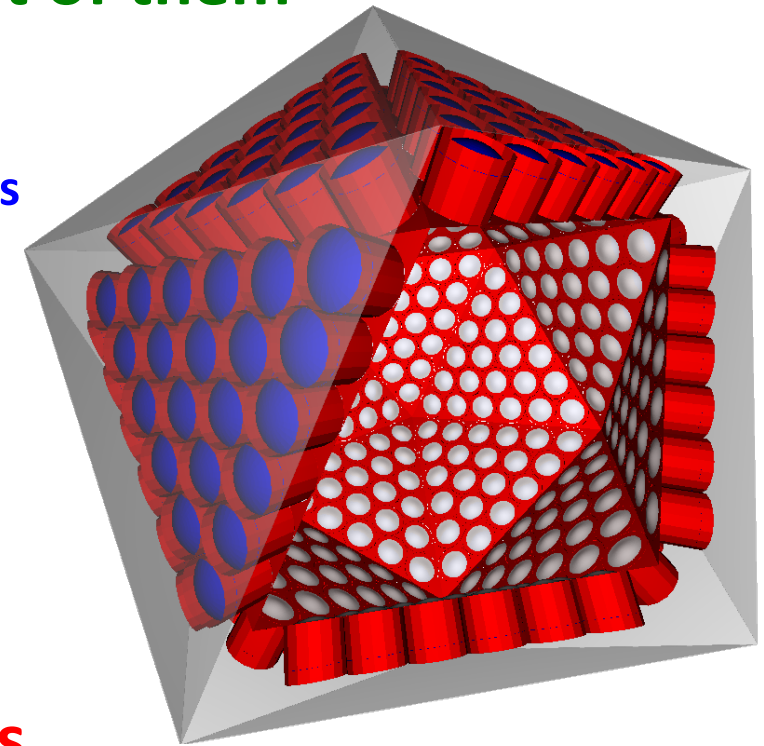
- Interpose a lenslet array in front of the CCD, with the CCD in the focus
- Each lenslet transforms the angle of the incident wavefront into a position on a subarray of the CCD
- Obtain an array of images also containing information on the depth
- Need a very large pixel count, as in some way this encodes the depth
- Needs plenty of software processing

R. Ng et al. “Light Field Photography with a Hand-Held Plenoptic Camera.” Stanford University Computer Science Tech Report CSTR 2005-02, April 2005.



Use a similar principle:

- further pixelate large photodetectors (think of the KamLAND PMTs pixelated)
- use lenses in front of them



Topological reconstruction then proceeds as tracking in a TPC: Each photon has

- a position of impact -- on the “optical module”, with uncertainty related to its diameter
- a direction of origin -- related to the quality of the lens and the size of the pixels

*Warning: don't get confused, in this “mode” depth of focus means nothing.
All lenses are focused to infinity!*

Note that no lens “owns” an image, not enough photons.

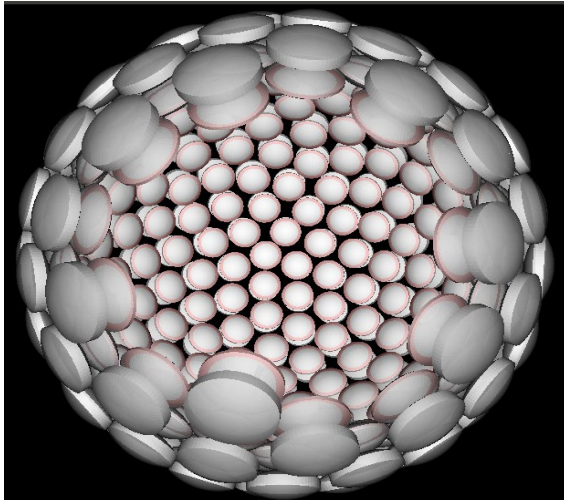
This is an important difference with respect to plenoptic cameras.

Here only putting together many lenses results in imaging.

**Detectors parameters are quite realistic, in terms of q.e. (33%)
and scintillator yield (8000 photons/MeV)**

But to get started a number of approximations/simplifications are used:

- Lens is stopped down for performance and the light outside of the active lens is suppressed (one could collect it for energy)
- Absorbing baffles between lens modules (simpler reconstruction)
- The detector is a icosahedron (simpler geometry)
 - ➔ Inscribed sphere is 14.8m for this case, basically like KamLAND



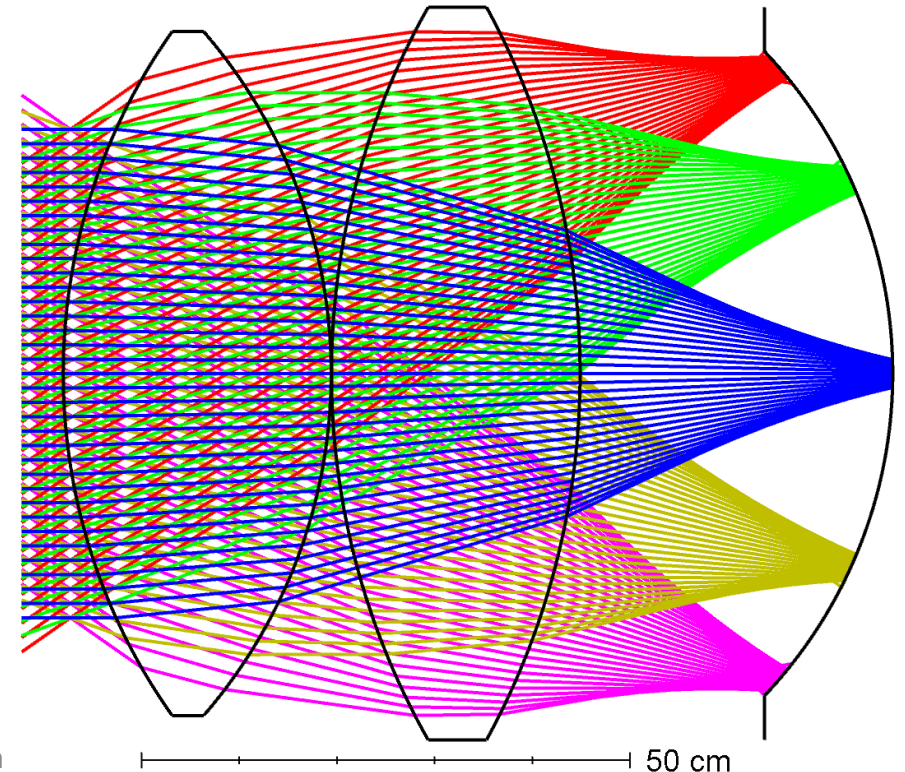
By now we also have a spherical detector model.

**Note the 2 lenses, photodetectors not shown
Performance is similar, but this is still work in progress**

Lens design is non-trivial (in fact, this held us up for a long time)

- Want a realistic system, in particular has to work in scintillator with $n_{\text{scint}} = 1.5$
- Require a very large NA (this is a very non-paraxial system, so forget most of what they taught you in school!)
- The focal surface is a spherical segment (simplifies the lens system)
- At least chromatic aberrations are not an issue (since scintillation light is ~monochromatic)
- Turns out that a single asphere is not necessarily the best solution, because “good practice” is to limit large changes of refractive at single surfaces
- From a first look air lenses appear less good, but I do not think we have investigated this enough

Parameter	Value
Number of elements	2 (spherical)
$\text{NA}_{\text{MAX}} (=n_{\text{scint}} \sin[\text{atan}(R_{\text{pupil}}/f_{\text{eff}})])$	0.64
Field of view	$\pm 40^\circ$
$R_{\text{lens}}/R_{\text{focal_surface}}$	1
Angular resolution	4°
n_{lens} (glass: S-NPH2, Ohara)	1.98



How to optimize a detector

Set the total number of pixels

→ our choice is 10^5

This sets the ~cost of the detector

(for 2000 KamLAND-style PMTs would mean
~8x8 pixels in the place of one 20-inch PMT)

Change the size of the lenses, hence changing the number of pixels/lens and find the best performance, given some desired property.

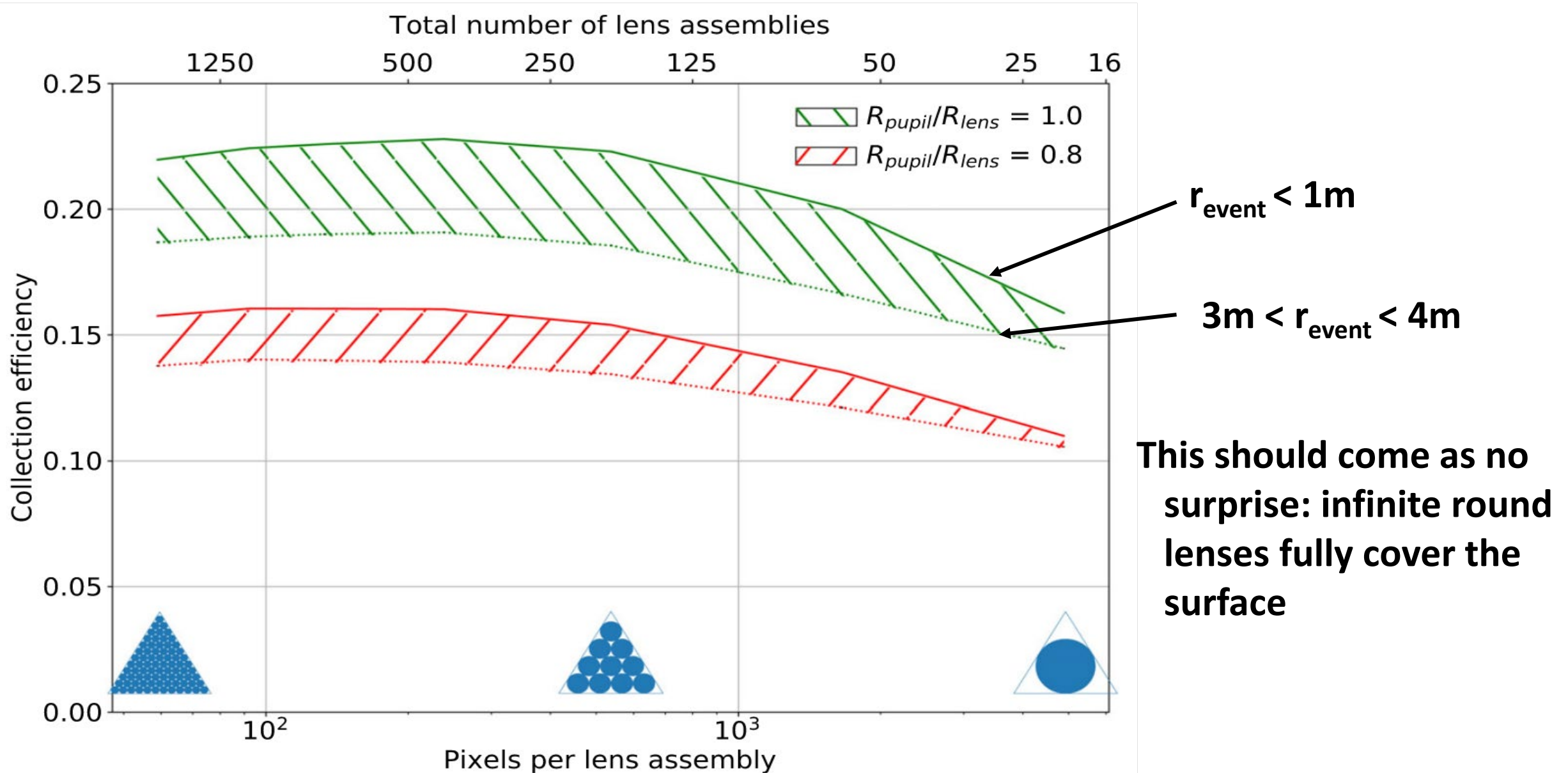
Larger lenses → more pixels/lens → better angular resolution for the direction of the photons → worse position resolution for the origin of the photon track

Smaller lenses → fewer pixels/lens → worse angular resolution → better position resolution

Note that lenses scale by similar transformations, so no need to re-design the lens in the optimization process.

*I do not have a “proof”, but I believe there
is a theorem that says that this is the optimal way to optimize imaging.*

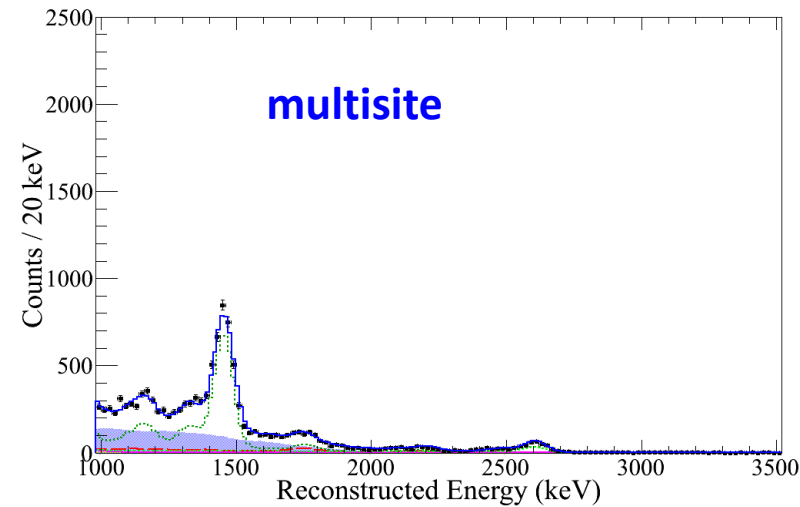
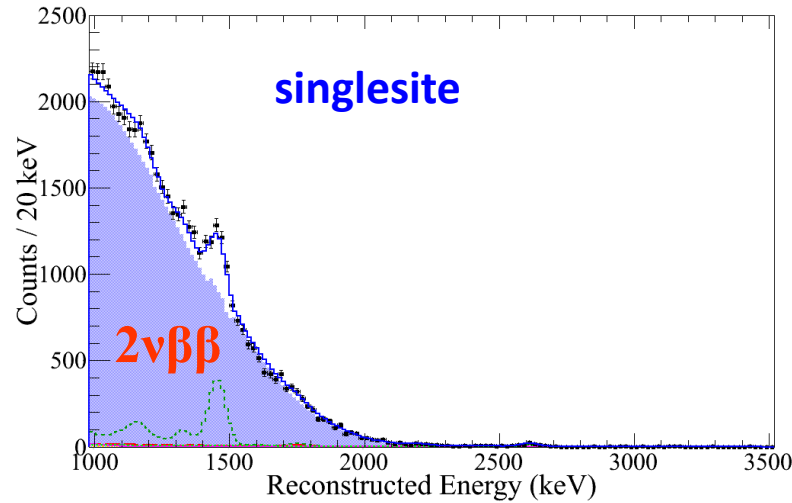
Results (2MeV deposited in the scintillator)



We choose as desired property the ability of discerning γ s from single e^- at a fixed energy of 2MeV

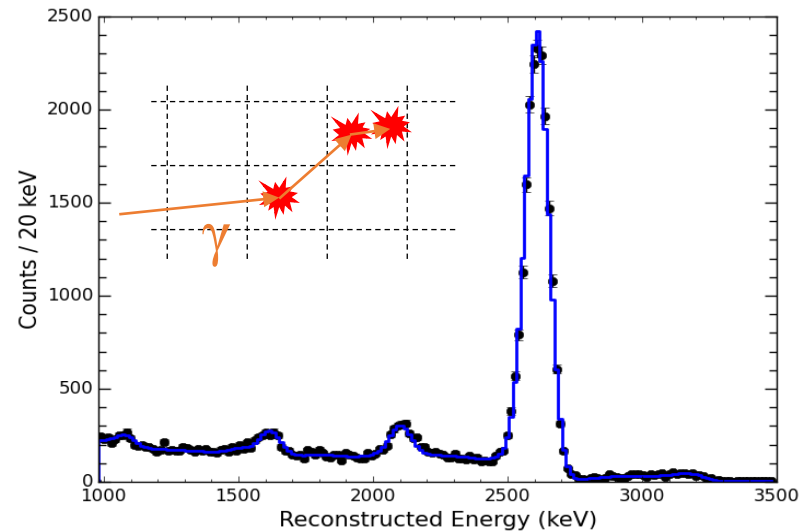
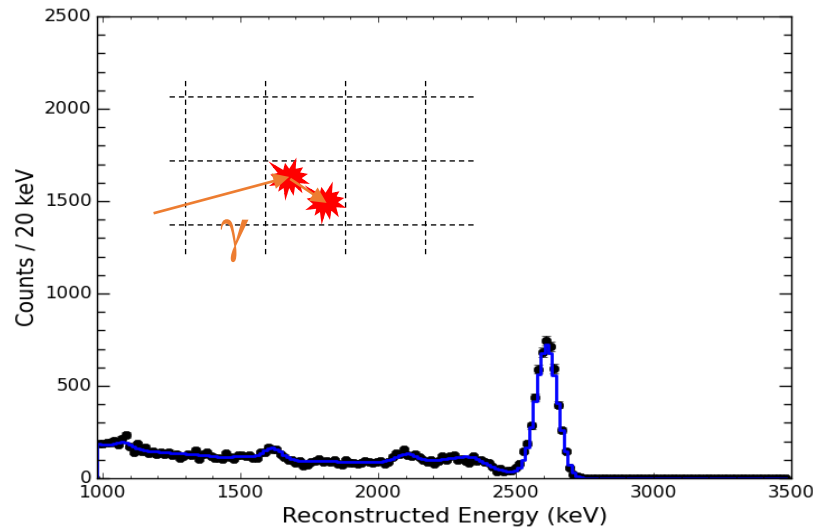
This is something that timing alone cannot do, and yet is it very important to reject some backgrounds

Low background
data

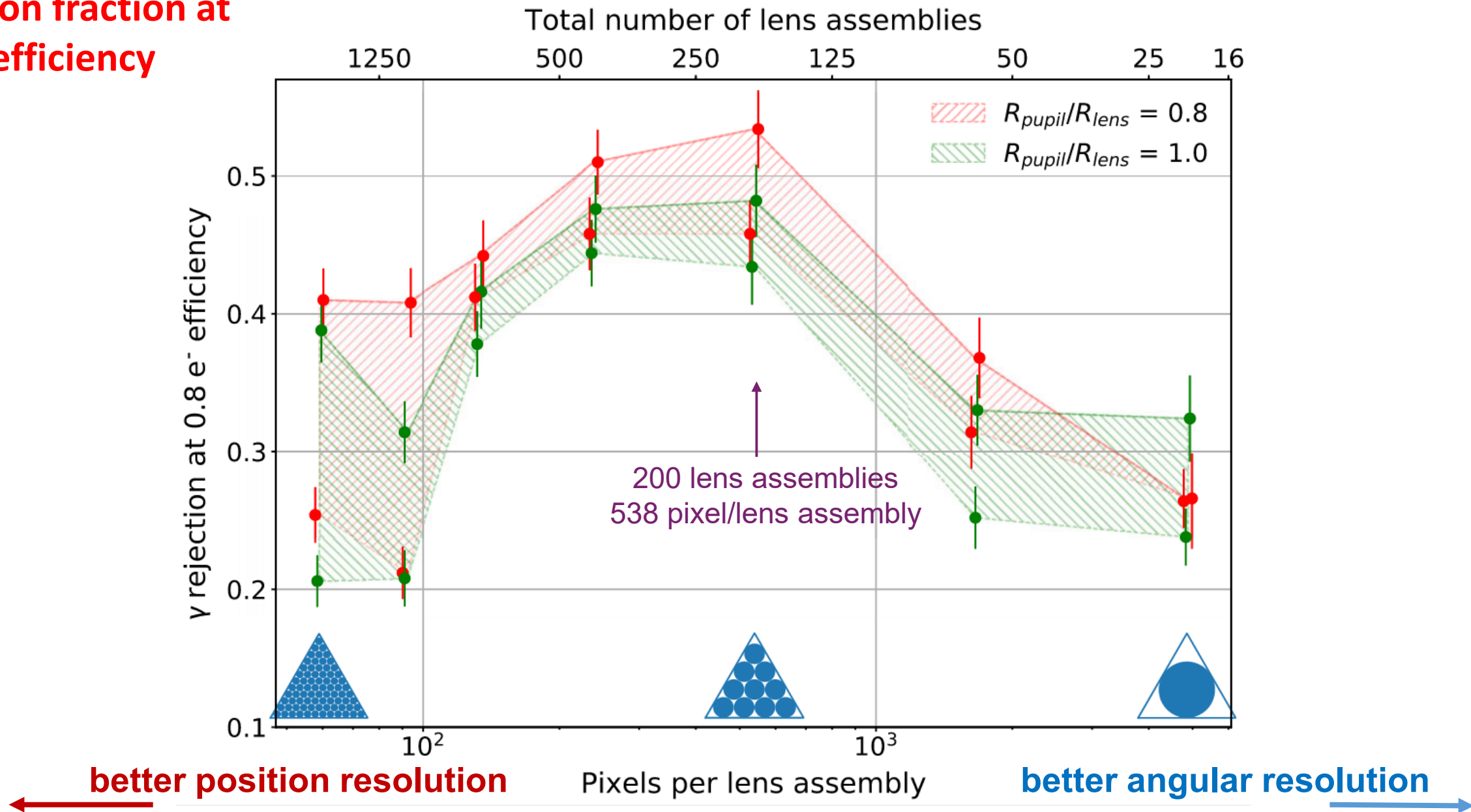


EXO-200 data

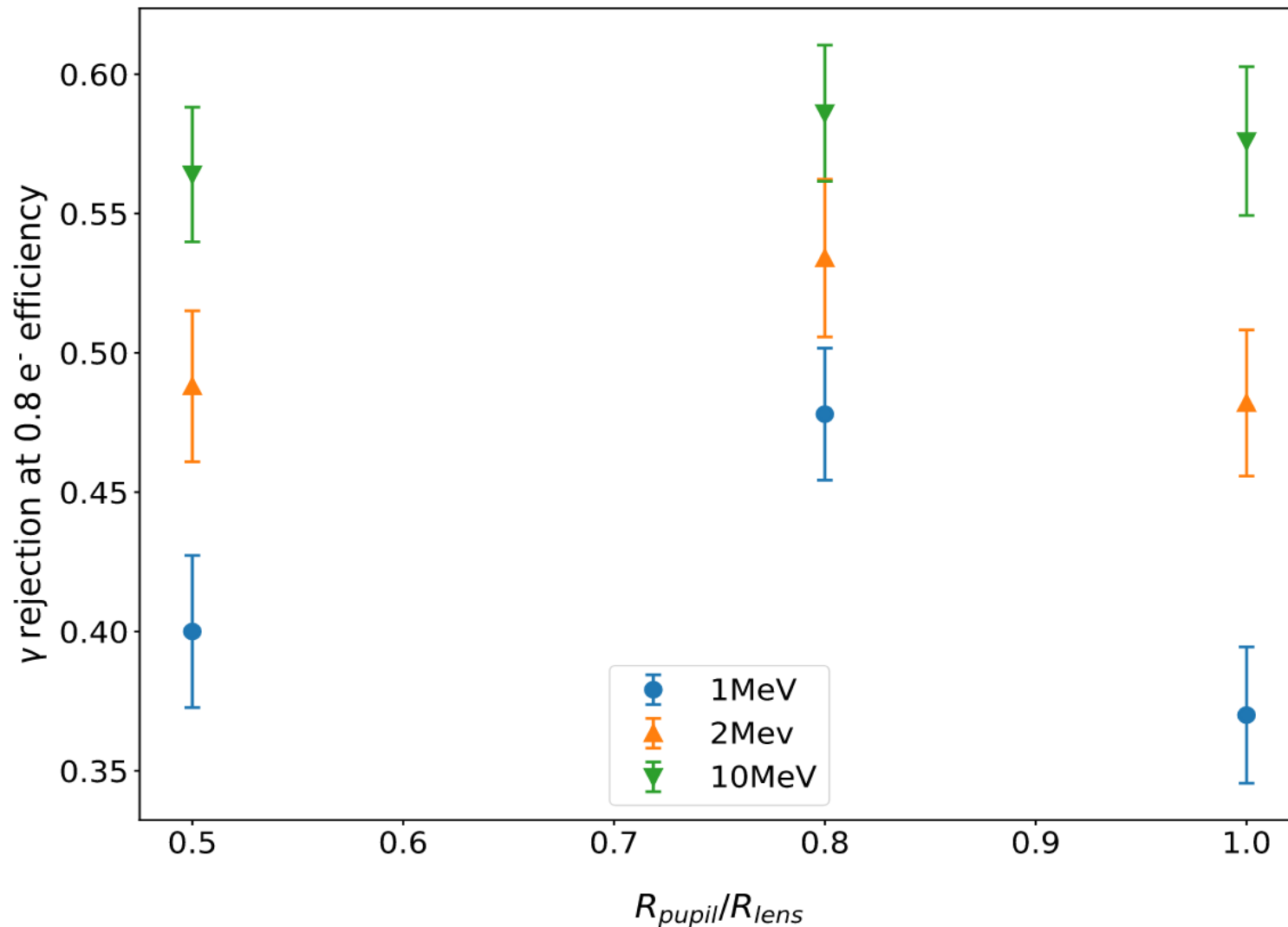
^{228}Th calibration
source



γ rejection fraction at 80% e^- efficiency

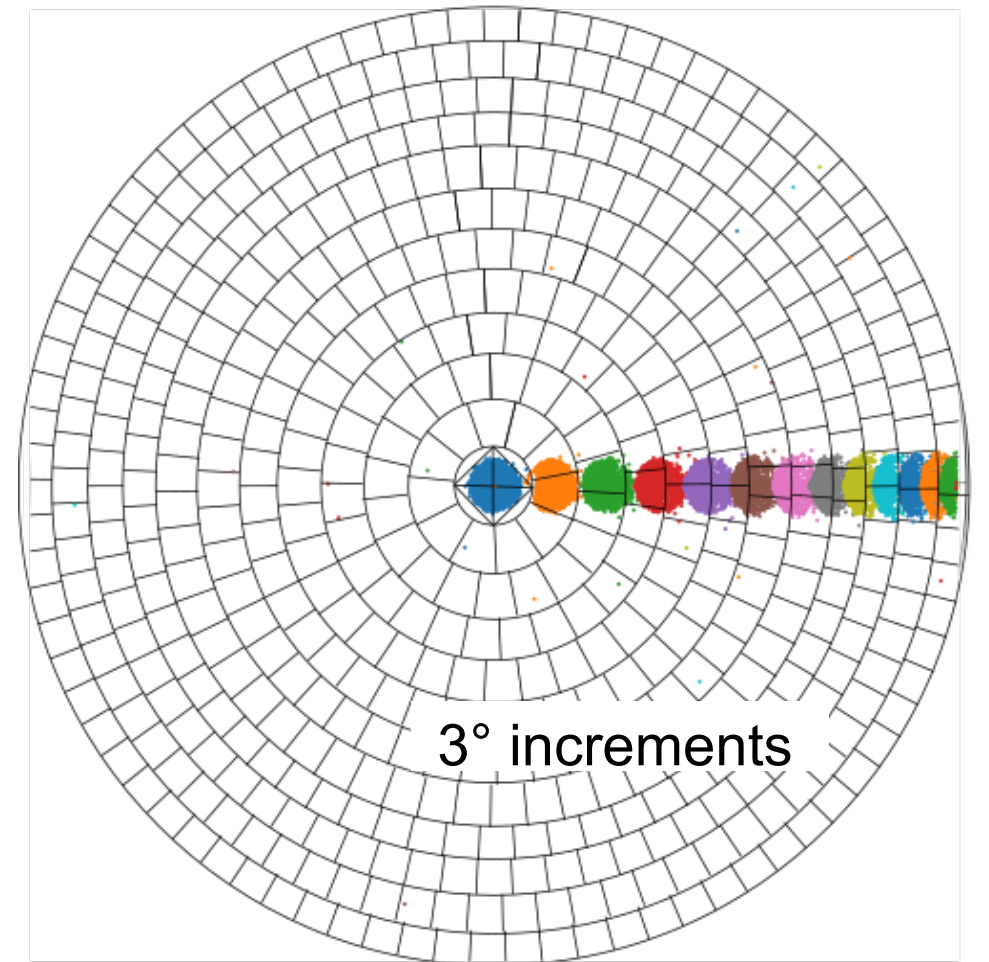
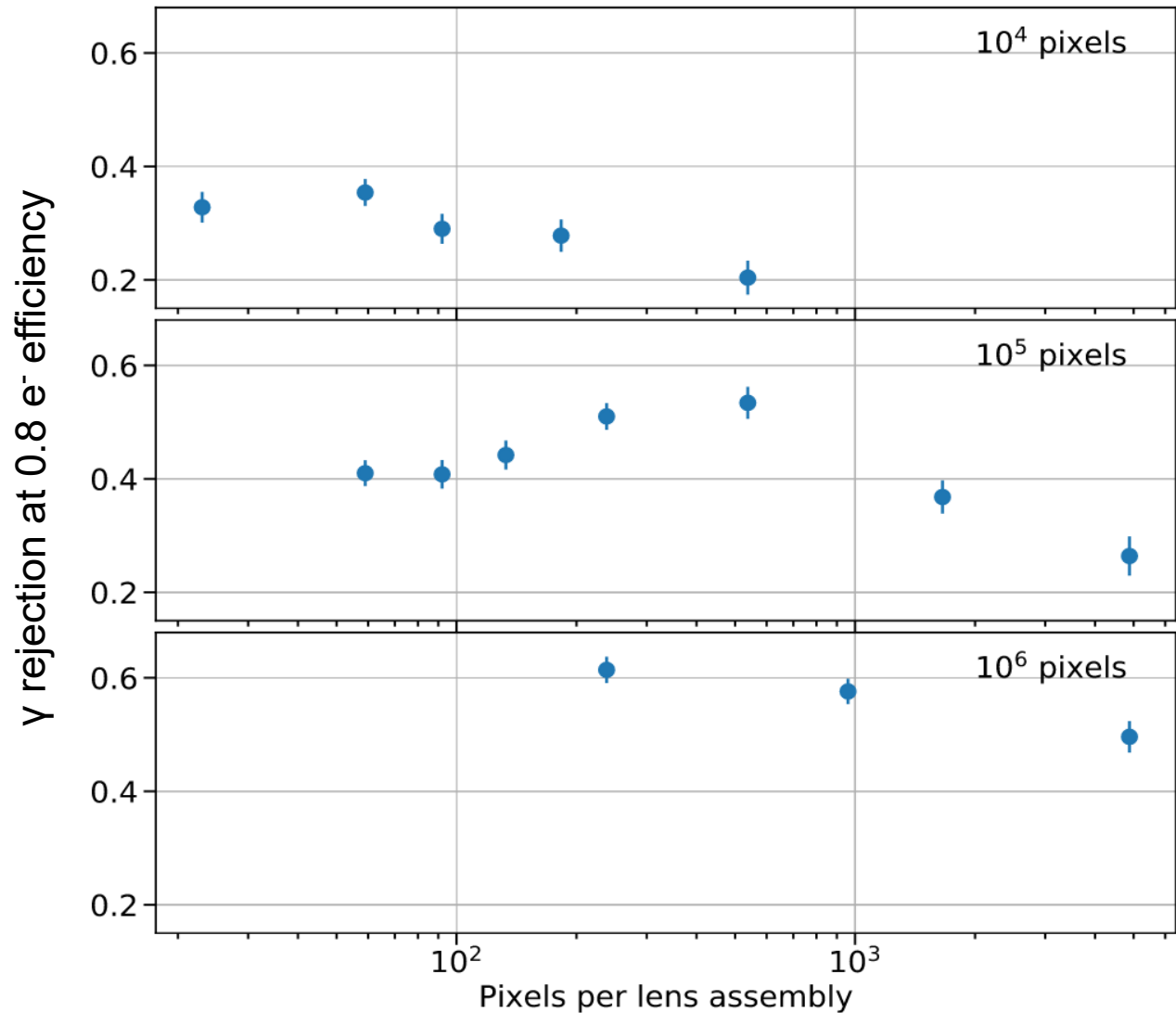


Higher energy achieves better performance and is less sensitive to the system segmentation



200 lens assemblies,
538 pixels/lens, $r_{event} < 1m$

Tradeoff between lens quality and number of pixels (these two have to be matched to each other)



Conclusions

We have found the optimal way to do imaging in scintillation detectors.

Note that, unlike reconstruction by time of arrival (resolution is constant, no matter the detector size) in this new technique $\delta d/D = \text{const}$ (D is the detector size), so that in small detectors, for the same number of pixels the resolution improves.

There may be uses in applied domains, where smaller detectors and better resolution are required.

For large detectors the timing information should be also used and will no doubt improve things (but we have not tried).

Need to investigate machine learning techniques → should be ideal for this.

Need help! Others pls help us and improve things further.

Details in J.Dalmasson et al., Phys. Rev. D 97 (2018) 052006.



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Now at FBK, Trento

And, most important, we need to develop better photodetectors!!

How come CRTs have all but disappeared from our desks while PMTs still reign in our labs?