

# Galactic Cosmic-Rays after AMS02

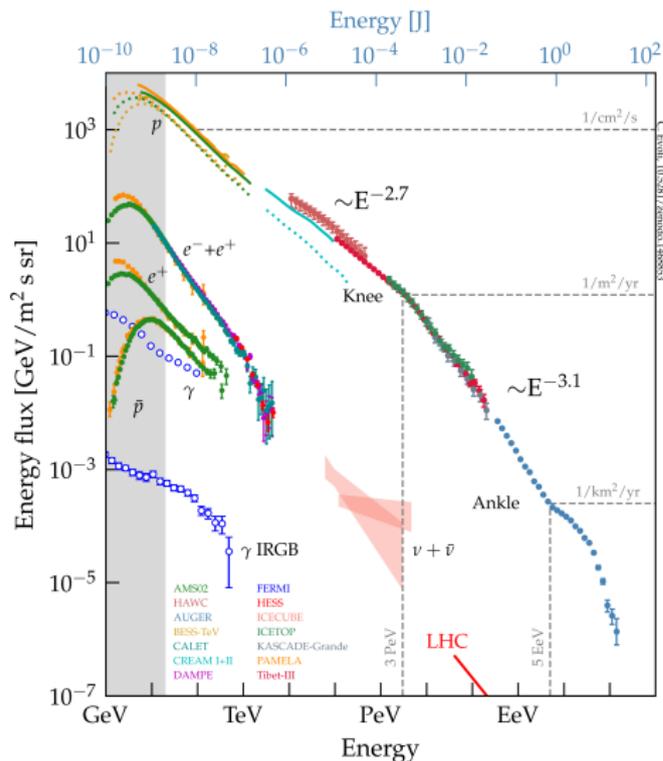
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2nd Cosmic-ray Antideuteron Workshop @ UCLA  
March 27, 2019



# The cosmic-ray spectrum



- ▶ **Non-thermal:** Almost a perfect power-law over more than 11 energy decades.
- ▶ Evidence of departures from a perfect power-law: the **knee** and the **ankle** features.
- ▶ Spectrum cut-off at  $\gtrsim 10^{20}$  eV.
- ▶ Particles observed at energy higher than any terrestrial laboratory.
- ▶ Composition at  $R \sim 10$  GV:
  - $\sim 99.2\%$  are nuclei
  - $\sim 84\%$  protons
  - $\sim 15\%$  He
  - $\sim 1\%$  heavier nuclei
  - $\sim 0.7\%$  are electrons

# The classical questions in CR physics

Gabici+, arXiv tomorrow(?)

- ▶ Which classes of sources contribute to the CR flux in different energy ranges?
- ▶ Which are the relevant processes responsible for CR confinement in the Galaxy?
- ▶ Are CR nuclei and electrons accelerated by the same sources?
- ▶ What is the origin of CR anti-matter?
- ▶ What is the role of CRs in the ISM? (e.g., for star formation)

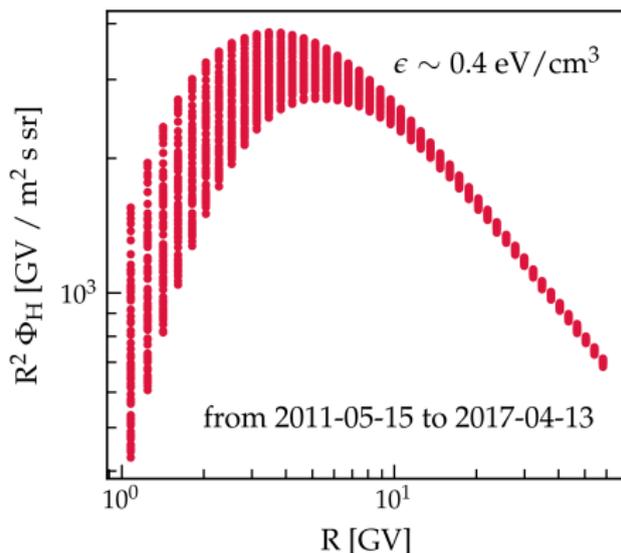
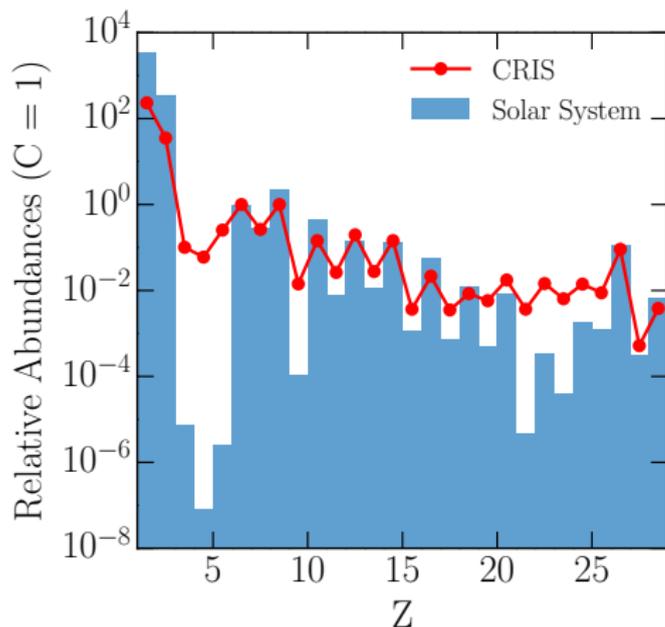


Figure: The TOA proton flux as measured by AMS02 at different times.

## LiBeB as cosmic-ray clocks



- ▶ If we assume that acceleration takes place in the average interstellar medium then this component must be produced during propagation (from that the term **secondary**).

# The grammage pillar

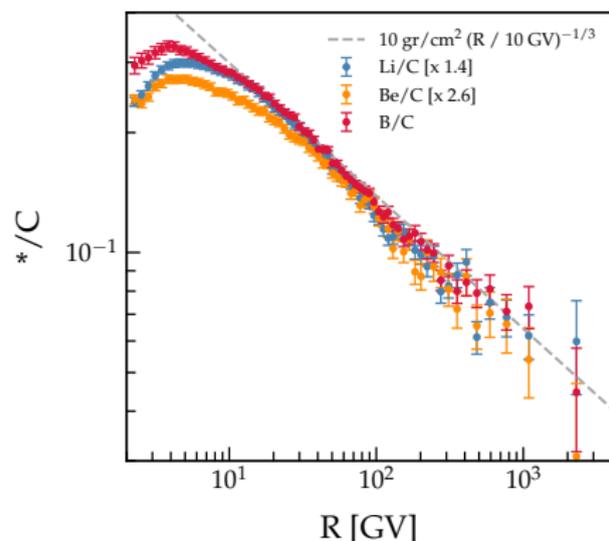


Figure: Secondary-over-primary ratios from AMS02.

- ▶ From this plot it follows the more robust evidence of **diffusion** so far:

$$B/C \sim \frac{X}{\bar{m}_{\text{ISM}}/\sigma_{C \rightarrow B}}$$

- ▶ following:

$$\tau_{\text{esc}}(10 \text{ GV}) \sim \frac{X(R)}{\bar{n}_{\text{ISM}}\mu v} \sim 90 \text{ Myr}$$

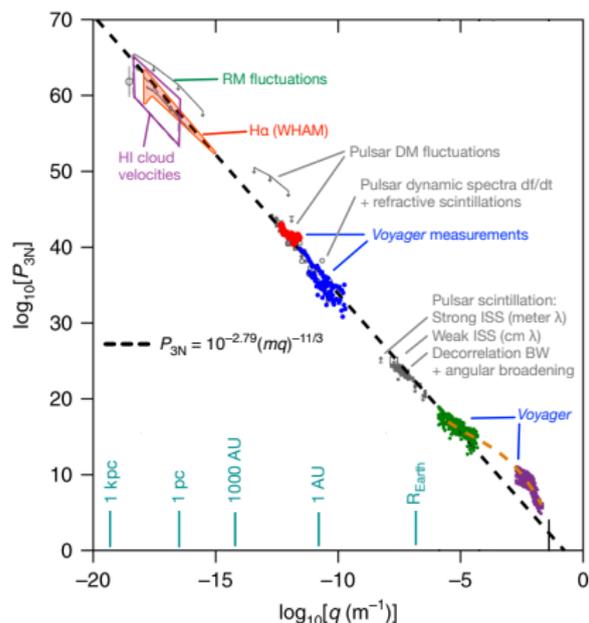
- ▶ while

$$\tau_{\text{ball}} \sim R_G/v \sim 3 \times 10^4 \text{ yr}$$

- ▶ The escape time is **energy dependent** and (roughly) scales like  $R^{-1/3}$



# The interstellar turbulence



Electron-density fluctuations in the ISM  
 [Armstrong+, ApJ 1995 - Chepurnov & Lazarian,  
 ApJ 2010 - Lee & Lee, Nature Astr. 2019]

- ▶ Turbulence is stirred by Supernovae at a typical scale  $L \sim 10 - 100$  pc
- ▶ Fluctuations of velocity and magnetic field are Alfvénic
- ▶ They have a Kolmogorov  $\alpha \sim -5/3$  spectrum (density is a passive tracer so it has the same spectrum:  $\delta n_e \sim \delta B^2$ ):

$$W(k)dk \equiv \frac{\langle \delta B \rangle^2(k)}{B_0^2} = \frac{2}{3} \frac{\eta_B}{k_0} \left( \frac{k}{k_0} \right)^{-\alpha}$$

- ▶ where  $k_0 = L^{-1}$  and the level of turbulence is

$$\eta_B = \int_{k_0}^{\infty} dk W(k) \sim 0.1 \div 0.01$$

# Charged particle in a turbulent field

Jokipii, ApJ 1966

- ▶ The turbulent field produces a **small fluctuation** with respect to the regular component

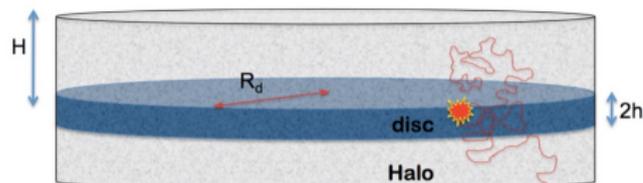
$$\langle \delta B^2 \rangle(k) \ll B_0^2 \text{ for } k \gg k_0$$

- ▶ The particle interacts resonantly with the waves, when the condition  $k_{\text{res}}^{-1} \sim r_L(p)$  is met
- ▶ The diffusion coefficient becomes:

$$D_{\text{QLT}}(p) = \frac{vr_L}{3} \frac{1}{k_{\text{res}} W(k_{\text{res}})} \sim \frac{3 \times 10^{27}}{\eta_B} \left( \frac{p}{\text{GeV}/c} \right)^{2-\alpha}$$

- ▶  $\lambda \sim \text{kpc}$  for  $k_{\text{res}} W(k_{\text{res}}) \sim 10^{-6}$  at scales  $\sim \text{A.U.}$
- ▶ that is just another example of the problem: **little things affect big things**

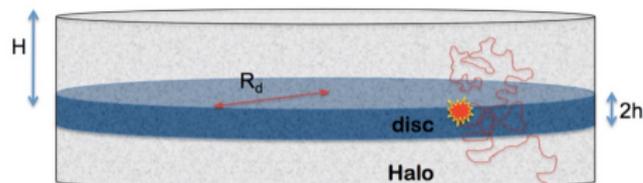
# The CR transport equation in the halo model



$$-\frac{\partial}{\partial z} \left( D_z \frac{\partial f_\alpha}{\partial z} \right) + u \frac{\partial f_\alpha}{\partial z} - \frac{du}{dz} \frac{p}{3} \frac{\partial f_\alpha}{\partial p} = q_{\text{SN}} - \frac{1}{p^2} \frac{\partial}{\partial p} [p^2 \dot{p} f_\alpha] - \frac{f_\alpha}{\tau_\alpha^{\text{in}}} + \sum_{\alpha' > \alpha} b_{\alpha' \alpha} \frac{f_{\alpha'}}{\tau_{\alpha'}^{\text{in}}}$$

► Spatial diffusion:  $\vec{\nabla} \cdot \vec{J}$

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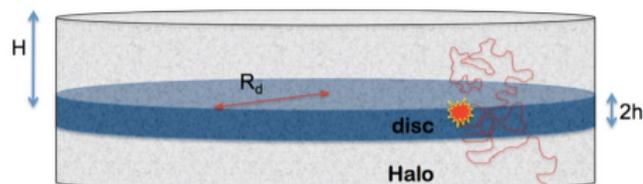


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► Spatial diffusion:  $\vec{\nabla} \cdot \vec{J}$

► Advection by Galactic winds/outflows:  $u = u_w + v_A \sim v_A$

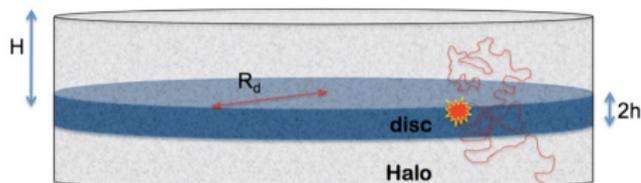
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- ▶ Spatial diffusion:  $\vec{\nabla} \cdot \vec{J}$
- ▶ Advection by Galactic winds/outflows:  $u = u_w + v_A \sim v_A$
- ▶ Source term proportional to Galactic SN profile

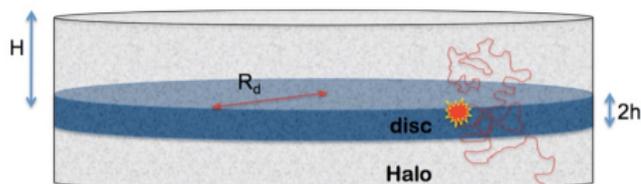
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- ▶ Source term proportional to Galactic SN profile
- ▶ Energy losses: ionization, Bremsstrahlung, IC, Synchrotron, ...
- ▶ Production/destruction of nuclei due to inelastic scattering (or decay)

## Predictions of the standard picture

For a primary CR species (e.g., H, C, O) at **high energy** we can ignore energy gain/losses, and the transport equation can be simplified as:

$$\frac{\partial f}{\partial t} = Q_0(p)\delta(z) + \frac{\partial}{\partial z} \left[ D \frac{\partial f}{\partial z} \right]$$

For  $z \neq 0$  one has:

$$D \frac{\partial f}{\partial z} = \text{constant} \rightarrow f(z) = f_0 \left( 1 - \frac{z}{H} \right)$$

where we used the definition of a halo:  $f(z = \pm H) = 0$ .

The typical solution gives (assuming injection  $Q \propto p^{-\gamma}$ ):

$$f_0(p) = \frac{Q_0(p)}{2A_d} \frac{H}{D(p)} \sim p^{-\gamma-\delta}$$

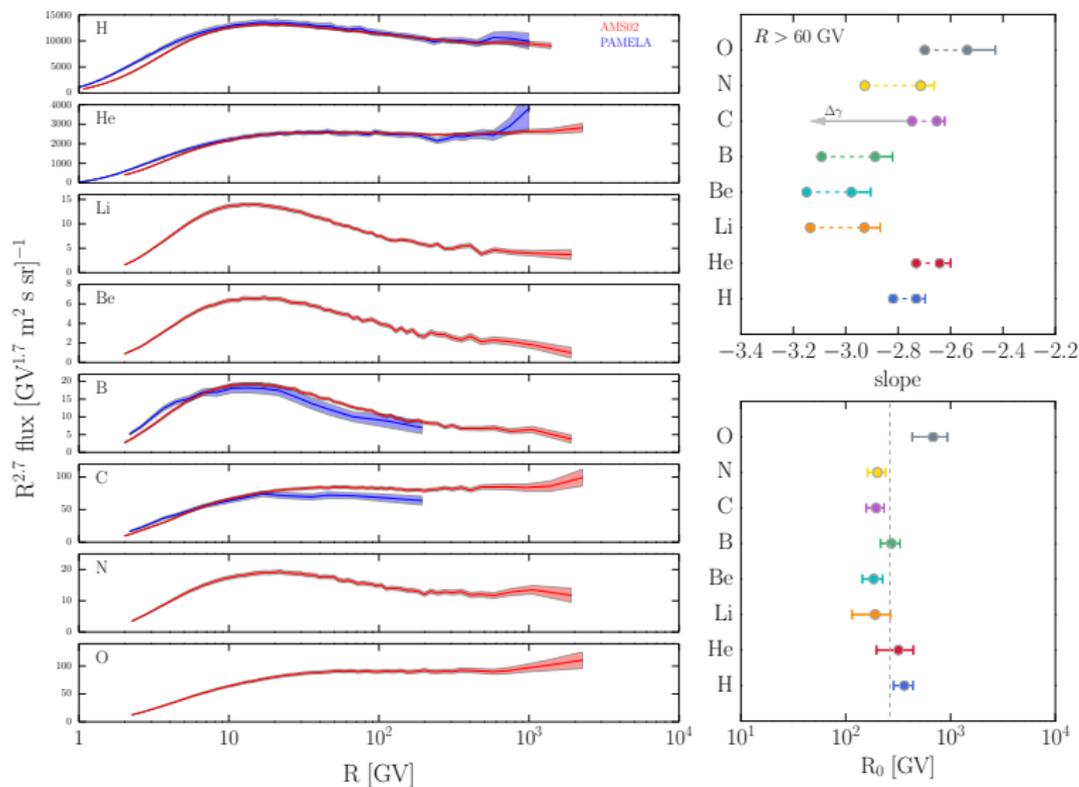
For a secondary (e.g., Li, Be, B) the source term is proportional to the primary density:

$$Q_B \sim \bar{n}_{\text{ISM}} c \sigma_{C \rightarrow B} N_C \rightarrow \frac{B}{C} \sim \frac{H}{D_0} p^{-\delta}$$

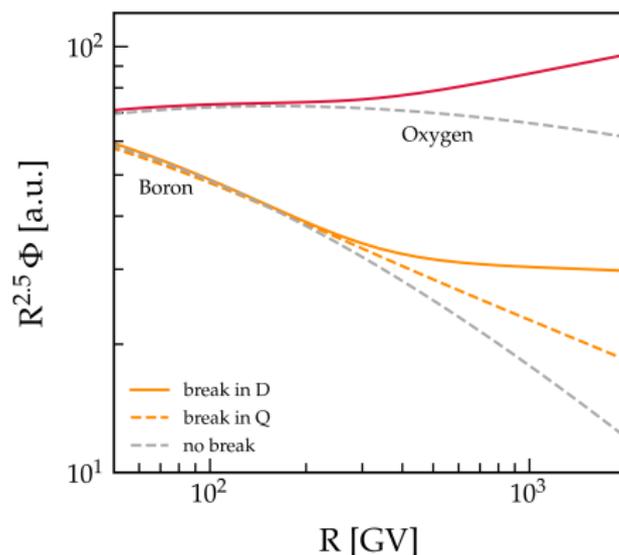
where we use  $\bar{n}_{\text{ISM}} = n_{\text{disk}} h/H$ .

# Unprecedented data precision: The rigidity break

Adriani+, Science 2011 - Aguilar+, PRLs 2013 and so on

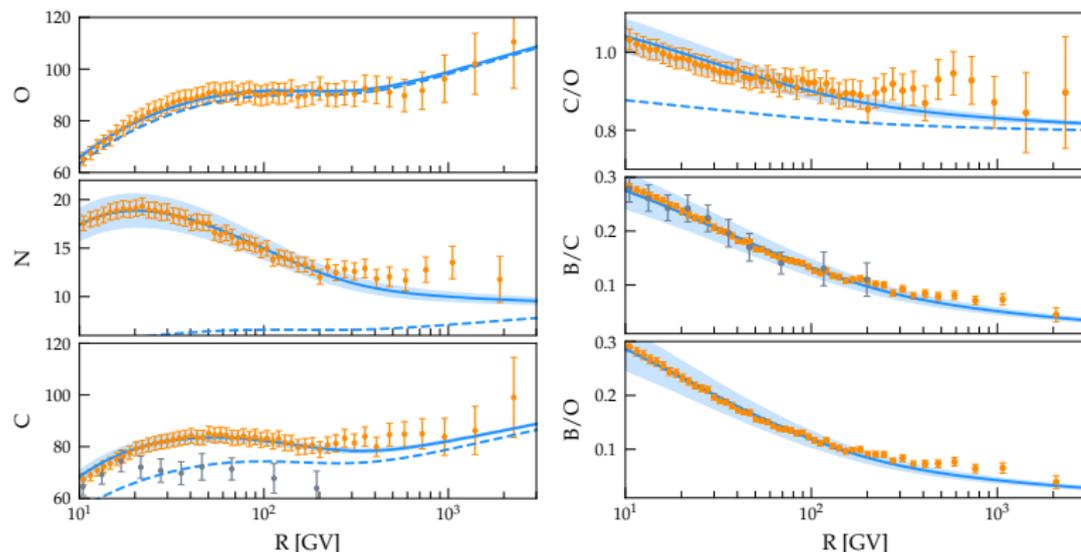


# The break is a propagation matter



- ▶ We conclude from the data that the observed spectral hardening at  $\sim 300$  GV is due to a change of regime in particle **diffusion**
- ▶ Similar conclusion from a Bayesian analysis in [Genolini+, PRL 2018]
- ▶ Physical mechanisms able to explain the break are presented in [Blasi, Amato & Serpico, PRL 2012 - Tomassetti, ApJL 752 (2012) 13]

# Fitting the nuclei heavier than He



- ▶ Modelling  $D$  with a smooth break:

$$D(R) = \beta D_0 \frac{(R/\text{GV})^\delta}{[1 + (R/R_b)^{\Delta\delta/s}]^s},$$

- ▶ we find  $\delta = 0.64$ ,  $D_0/H = 0.25 \times 10^{28}$  cm/s<sup>2</sup>,  $\Delta\delta = 0.2$ ,  $u = 7$  km/s and  $\gamma = 4.26$
- ▶ **B/C and C/O as grammage indicators are severely limited by our knowledge of cross-sections.**

# The problem with cross-sections: need for new measurements

Genolini+, PRC 2018 - Reinert & Winkler, JCAP 2018 - Evoli+, JCAP 2018

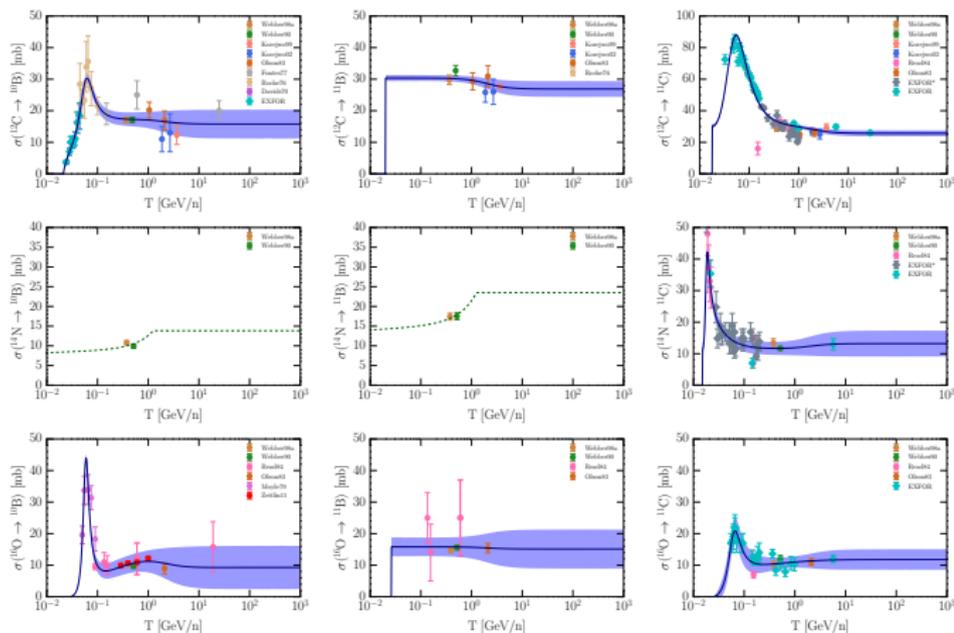
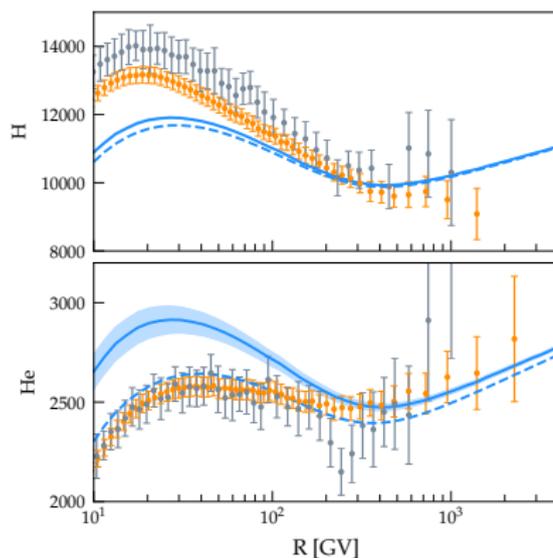


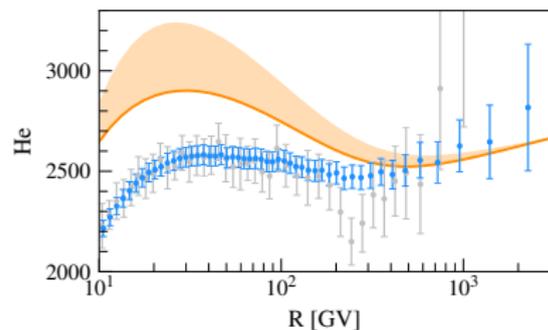
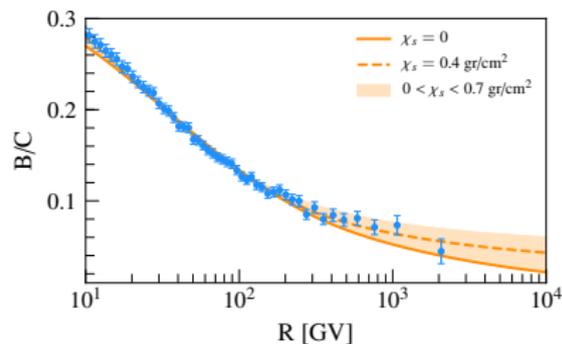
Figure: Cross-sections for Boron production by CNO spallation on Hydrogen target as a function of kinetic energy per nucleon. Data are taken from GALPROP and from EXFOR database.

# The injection drama



- ▶ H is **softer** than nuclei, while He is **harder**
- ▶ At odds with what one would expect in the case of pure rigidity dependent acceleration [Serpico, ICRC 2015].
- ▶ Problematic even for models of the difference between H and He injection based on the different  $A/Z$  at shocks [Hanusch+, Apj 2019].

# Grammage at the source



- ▶ To provide a better fit of high-energy B/C we account for an additional contribution to the grammage traversed by CRs
- ▶ The grammage due to confinement inside a SNR can be easily estimated as [Aloisio, Blasi & Serpico, A&A 2015]

$$X_{\text{SNR}} \sim 0.2 \text{ g/cm}^{-2}$$

- ▶ It is important at high-energy since the harder spectrum
- ▶ B/C can constrain  $\chi_s \lesssim 0.7 \text{ gr/cm}^2$
- ▶ However the injection problem for He gets worse!

## A new scenario for cosmic-ray propagation in the halo

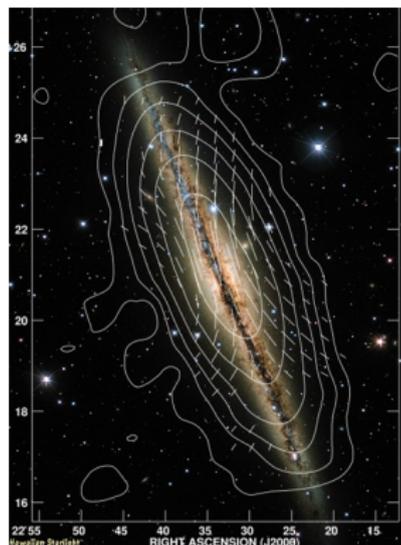
- ▶ By solving the transport equation we obtain a featureless (at least up to the knee) propagated spectrum for each primary species, differently than what is observed.
- ▶ This result remains true even in more sophisticated approaches as GALPROP or DRAGON
- ▶ **What is missing in our physical picture?**

# The halo size $H$

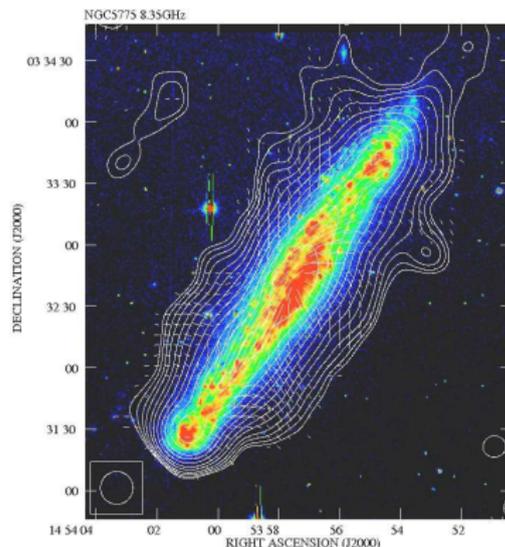
- ▶ Assuming  $f(z = H) = 0$  reflects the requirement of lack of diffusion (infinite diffusion coefficient)
- ▶ May be because  $B \rightarrow 0$ , or because turbulence vanishes (in both cases  $D$  cannot be spatially constant!)
- ▶ Vanishing turbulence may reflect the lack of sources
- ▶ Can be  $H$  dependent on  $p$ ? (remember  $\mathbf{B/C} \sim D/H!$ )
- ▶ **What is the physical meaning of  $H$ ?**

# The radio halo in external galaxies

Credit: MPIfR Bonn



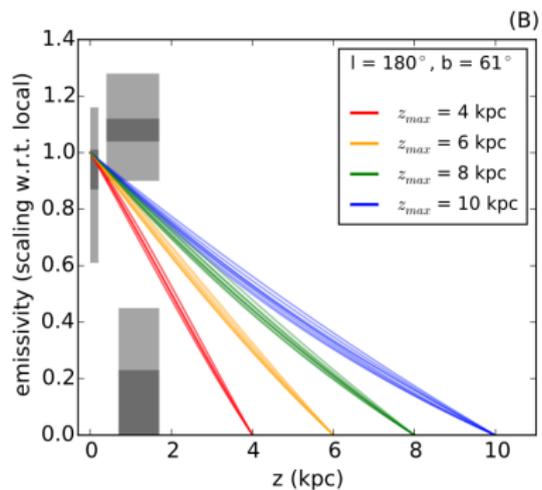
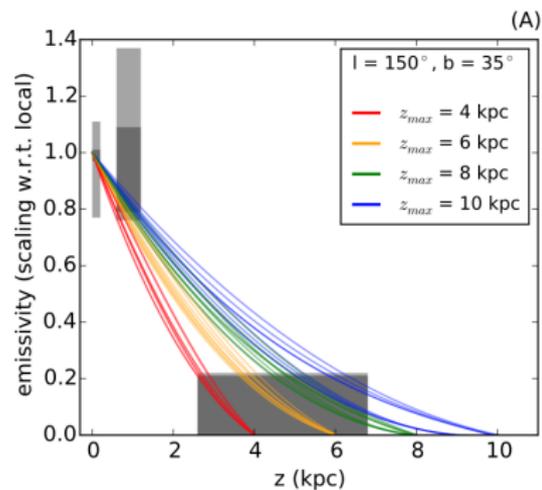
Total radio emission and B-vectors of edge-on galaxy NGC891, observed at 3.6 cm wavelength with the Effelsberg telescope



Total radio intensity and B-vectors of edge-on galaxy NGC 5775, combined from observations at 3.6 cm wavelength with the VLA and Effelsberg telescopes

# The $\gamma$ -halo in our Galaxy

Tibaldo et al., 2015, ApJ



- ▶ Using high-velocity clouds one can measure the emissivity per atom as a function of  $z$  (proportional to  $f$ )
- ▶ Indication of a halo with  $H \sim$  few kpc

# Non-linear cosmic ray transport

Skilling71, Wentzel74

- ▶ The net effect of spatial diffusion is to reduce the momentum of the particles forcing them, eventually, to move at the same speed as the waves  $\sim v_A$
- ▶ If CR stream faster than the waves, the net effect of diffusion is to make waves grow and make CR diffusive motion slow down: this process is known as **self-generation of waves** (notice that self-generated waves are  $k \sim r_L$ )
- ▶ Waves are amplified by CRs through streaming instability:

$$\Gamma_{\text{CR}} = \frac{16\pi^2}{3} \frac{v_A}{kW(k)B_0^2} \left[ v(p)p^4 \frac{\partial f}{\partial z} \right]$$

and are damped by wave-wave interactions that lead the development of a turbulent cascade (NLLD):

$$\Gamma_{\text{NLLD}} = (2c_k)^{-3/2} kv_A(kW)^{1/2}$$

- ▶ What is the typical scale/energy up to which self-generated turbulence is dominant?

# Non-linear cosmic ray transport

Blasi, Amato & Serpico, PRL, 2012

Transition occurs at scale where external turbulence (e.g., from SNe) equals in energy density the self-generated turbulence

$$W_{\text{ext}}(k_{\text{tr}}) = W_{\text{CR}}(k_{\text{tr}})$$

where  $W_{\text{CR}}$  corresponds to  $\Gamma_{\text{CR}} = \Gamma_{\text{NLLD}}$

Assumptions:

- ▶ Quasi-linear theory applies
- ▶ The external turbulence has a Kolmogorov spectrum
- ▶ Main source of damping is non-linear damping
- ▶ Diffusion in external turbulence explains high-energy flux with SNR efficiency of  $\epsilon \sim 10\%$

$$E_{\text{tr}} = 228 \text{ GeV} \left( \frac{R_{d,10}^2 H_3^{-1/3}}{\epsilon_{0.1} E_{51} \mathcal{R}_{30}} \right)^{3/2(\gamma_p-4)} B_{0,\mu}^{(2\gamma_p-5)/2(\gamma_p-4)}$$

# The turbulence evolution equation

Eilek, ApJ 1979

$$\frac{\partial W}{\partial t} = \frac{\partial}{\partial k} \left[ D_{kk} \frac{\partial W}{\partial k} \right] + \frac{\partial}{\partial z} (v_A W) + \Gamma_{\text{CR}} W + Q(k)$$

► Diffusion in  $k$ -space damping:  $D_{kk} = c_k |v_A| k^{7/2} W^{1/2}$

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- ▶ Waves growth due to cosmic-ray streaming:  $\Gamma_{\text{CR}} \propto \partial f / \partial z$

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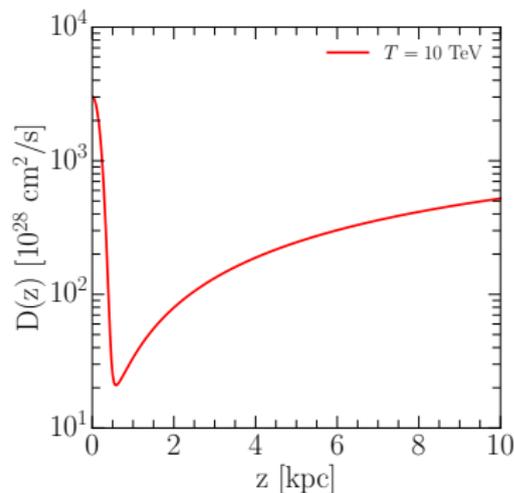
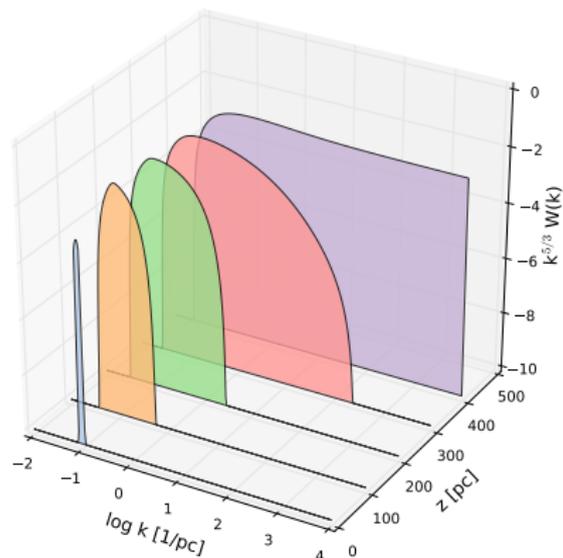
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- ▶ Waves growth due to cosmic-ray streaming:  $\Gamma_{\text{CR}} \propto \partial f / \partial z$
- ▶ External (e.g., SNe) source term  $Q \sim \delta(z) \delta(k - k_0)$
- ▶ In the absence of the instability, it returns a kolmogorov spectrum:  
 $W(k) \sim k^{-5/3}$

# Wave advection $\rightarrow$ the turbulent halo

Evoli, Blasi, Morlino & Aloisio, 2018, PRL



$$\tau_{\text{cascade}} = \tau_{\text{adv}} \rightarrow \frac{k_0^2}{D_{kk}} = \frac{z_{\text{peak}}}{v_A} \rightarrow z_{\text{peak}} \sim \mathcal{O}(\text{kpc})$$

# Non-linear cosmic ray transport: diffusion coefficient

Evoli, Blasi, Morlino & Aloisio, 2018, PRL

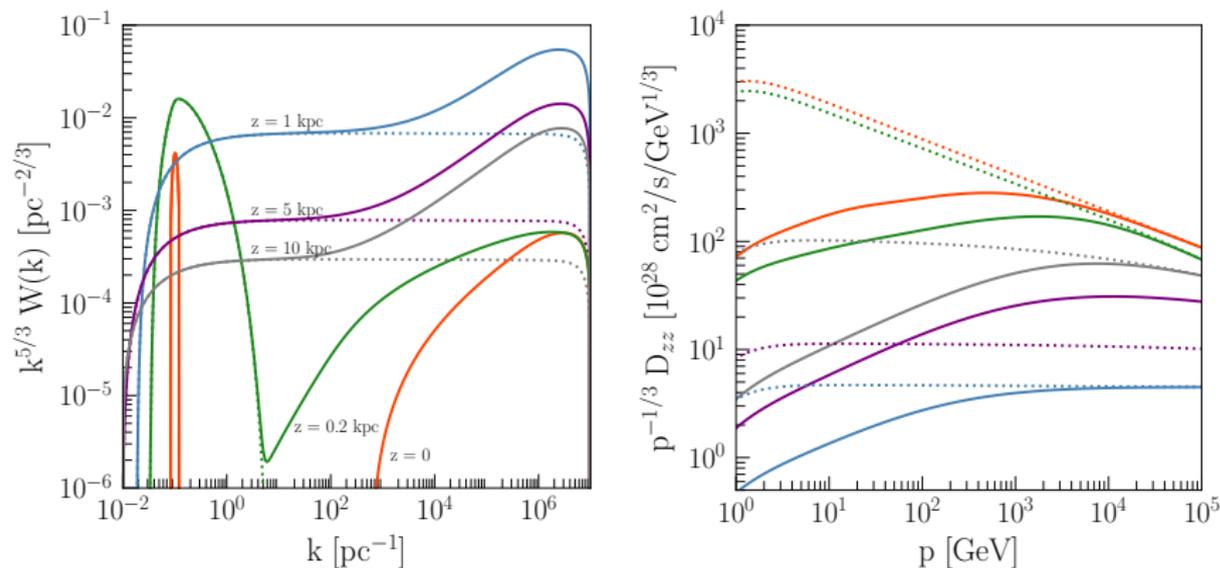
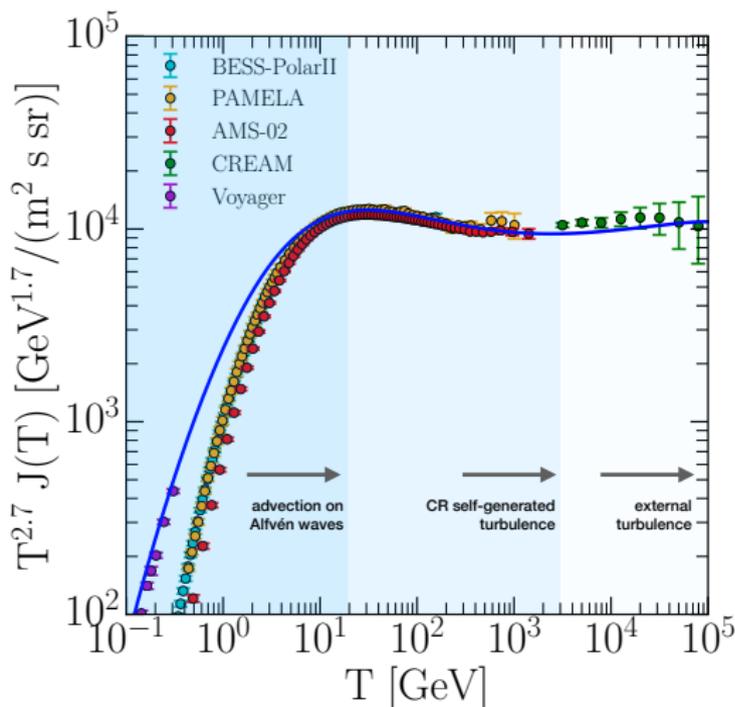


Figure: Turbulence spectrum without (dotted) and with (solid) CR self-generated waves at different distance from the galactic plane.

# Non-linear cosmic ray transport: a global picture

Evoli, Blasi, Morlino & Aloisio, 2018, PRL



- ▶ Pre-existing waves (Kolmogorov) dominates above the break
- ▶ Self-generated turbulence between 1-100 GeV
- ▶ Voyager data are reproduced with no additional breaks, but due to advection with self-generated waves (single injection slope)
- ▶  $H$  is not predetermined here.
- ▶ None of these effects were included in the numerical simulations of CR transport before.

# Conclusions

- ▶ Recent findings by PAMELA and AMS-02 (breaks in the spectra of primaries, B/C à la Kolmogorov, flat anti-protons, rising positron fraction) are challenging the standard scenario of CR propagation.
- ▶ Non-linearities might play an essential role for propagation (as they do for acceleration). They allow to reproduce local observables (primary spectra) without ad hoc breaks.
- ▶ We present a non-linear model in which SNRs inject: a) turbulence at a given scale with efficiency  $\epsilon_w \sim 10^{-4}$  and b) cosmic-rays with a single power-law and  $\epsilon_{CR} \sim 10^{-1}$ . The turbulent halo and the change of slope at  $\sim 300$  GV are obtained self-consistently.
- ▶ As a bonus, these models enable us a deeper understanding of the interplay between CR, magnetic turbulence and ISM in our Galaxy.

# The road ahead..

