

Updated secondary anti-helium cosmic ray fluxes

A first estimate for ${}^4\overline{\text{He}}$

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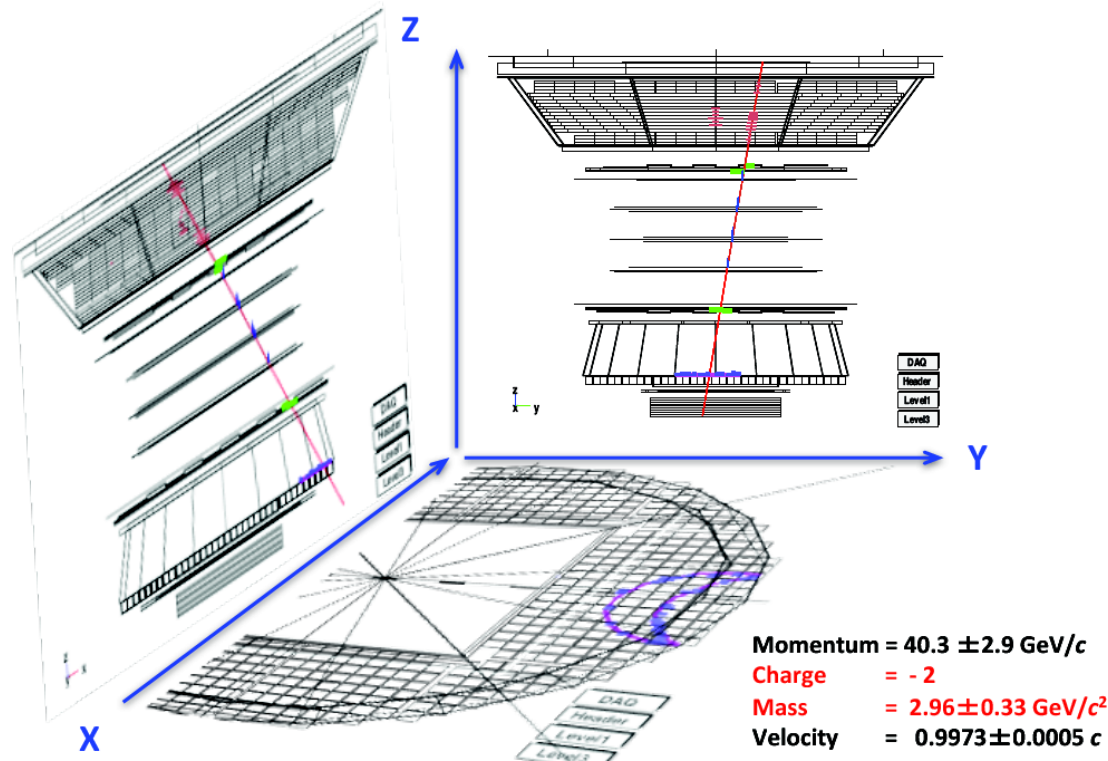
Outline

- 1) Anti-helium et advocati diaboli
- 2) Anti-helium production and the coalescence factor
- 3) Cosmic ray anti-helium Galactic propagation
- 4) A word on dark matter

Based on Phys. Rev. **D99** (2019) 023016

V. Poulin, **P.S.**, I. Cholis, M. Kamionkowski & J. Silk

1) Anti-helium et advocati diaboli

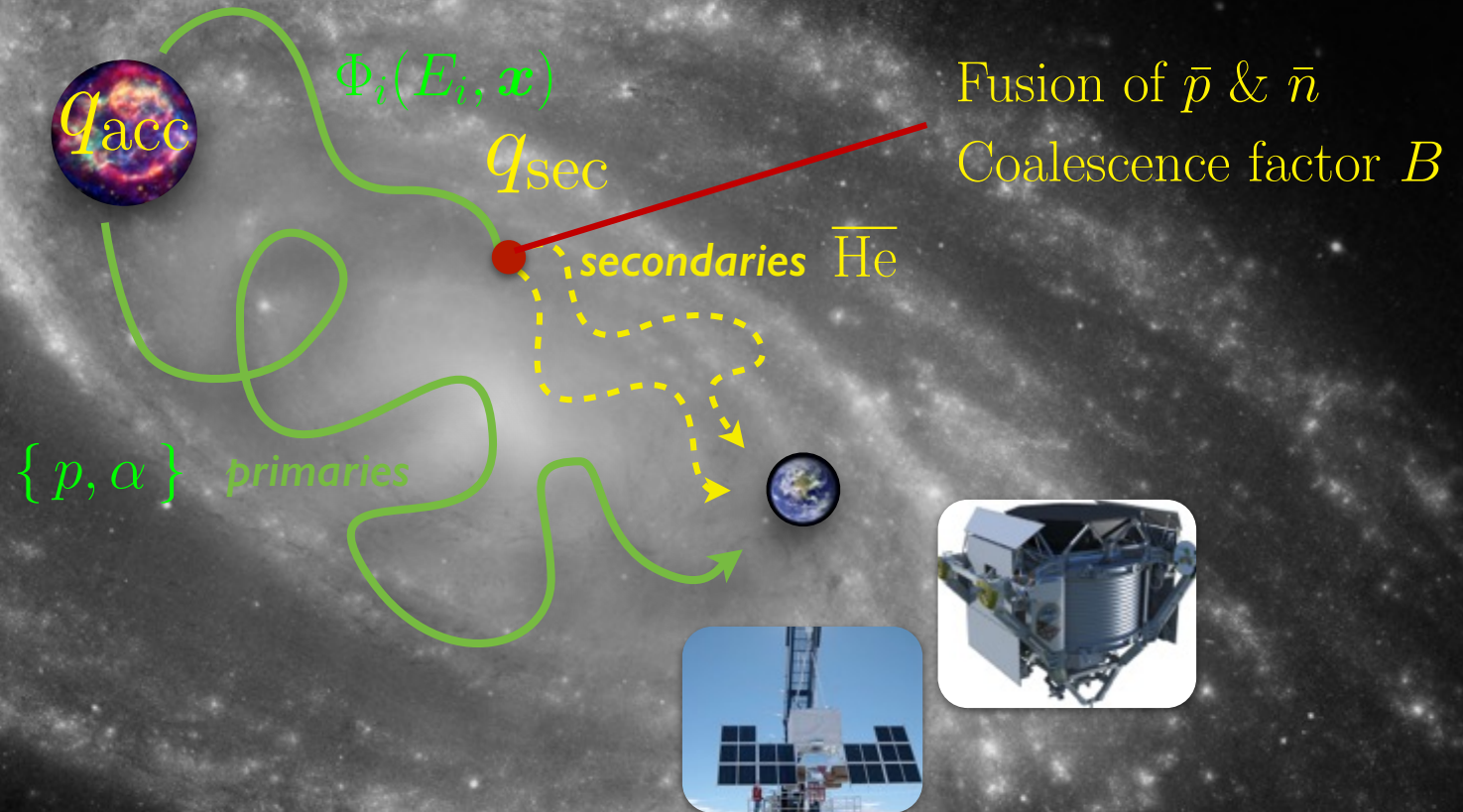


P. von Doetinchem, DSU June 18, 2018

- ${}^3\overline{\text{He}}$ (6) and ${}^4\overline{\text{He}}$ (2) candidates have been identified by AMS-02.
The event rate is ~ 1 anti-helium in 100 million helium.
- Massive background simulations are carried out to evaluate significance.
The probability of a background origin for $\overline{\text{He}}$ events is very small.
- More data are needed. Number of collected $\overline{\text{He}}$ events should increase,
while probability of background origin should decrease.

Secondary cosmic-ray anti-helium

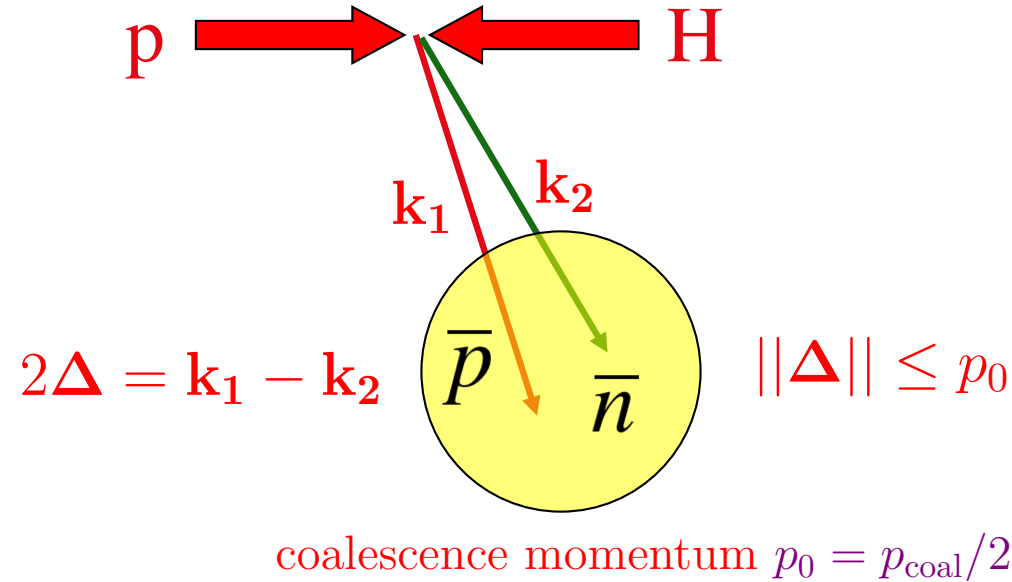
$$q_{\text{sec}}(\overline{\text{He}} | E_{\overline{\text{He}}}, \mathbf{x}) = \sum_{i \in p, \alpha} \sum_{j \in \text{H, He}} 4\pi \int dE_i \Phi_i(E_i, \mathbf{x}) n_j(\mathbf{x}) \frac{d\sigma_{ij \rightarrow \overline{\text{He}}}}{dE_{\overline{\text{He}}}}(E_i, E_{\overline{\text{He}}})$$



Solar modulation with $\phi_p^{\text{F}} \neq \phi_{\bar{p}}^{\text{F}}$

2) Anti-helium production and the coalescence factor

coalescence $\equiv$ fusion of $\bar{p}$ & $\bar{n}$ into $\bar{d}$, ${}^3\overline{\text{He}}$ or ${}^4\overline{\text{He}}$



$$d^3\mathcal{N}_{\bar{d}}(\mathbf{K}) = \int d^6\mathcal{N}_{\bar{p},\bar{n}}\{\mathbf{k}_1, \mathbf{k}_2\} \times \mathcal{C}(\Delta) \times \delta^3(\mathbf{K} - \mathbf{k}_1 - \mathbf{k}_2)$$

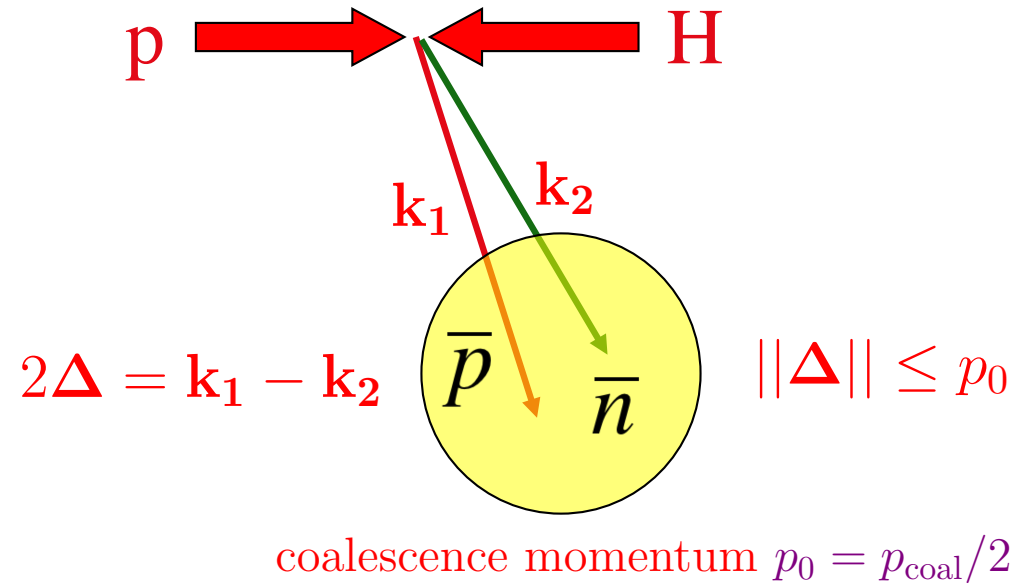
$$B_2 = \frac{E_{\bar{d}}}{E_{\bar{p}}E_{\bar{n}}} \int d^3\Delta \mathcal{C}(\Delta) \simeq \frac{m_{\bar{d}}}{m_{\bar{p}}m_{\bar{n}}} \left\{ \frac{4}{3}\pi p_0^3 \equiv \frac{\pi}{6} p_{\text{coal}}^3 \right\}$$

Coalescence factor B_2

$$\frac{E_{\bar{d}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{d}}}{d^3\mathbf{K}} = B_2 \left\{ \frac{E_{\bar{p}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{p}}}{d^3\mathbf{k}_1} \right\} \left\{ \frac{E_{\bar{n}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{n}}}{d^3\mathbf{k}_2} \right\}$$

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Production on anti-nuclei with mass A

$$\frac{E_{\bar{A}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{A}}}{d^3\mathbf{k}_{\bar{A}}} = B_A \left\{ \frac{E_{\bar{p}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{p}}}{d^3\mathbf{k}_{\bar{p}}} \right\}^Z \left\{ \frac{E_{\bar{n}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{n}}}{d^3\mathbf{k}_{\bar{n}}} \right\}^{A-Z} \quad \text{with} \quad \mathbf{k}_{\bar{p}} = \mathbf{k}_{\bar{n}} = \mathbf{k}_{\bar{A}}/A$$

Coalescence factor B_A

$$B_A = \frac{m_A}{m_p^Z m_n^{A-Z}} \left\{ \frac{\pi}{6} p_{\text{coal}}^3 \right\}^{A-1}$$

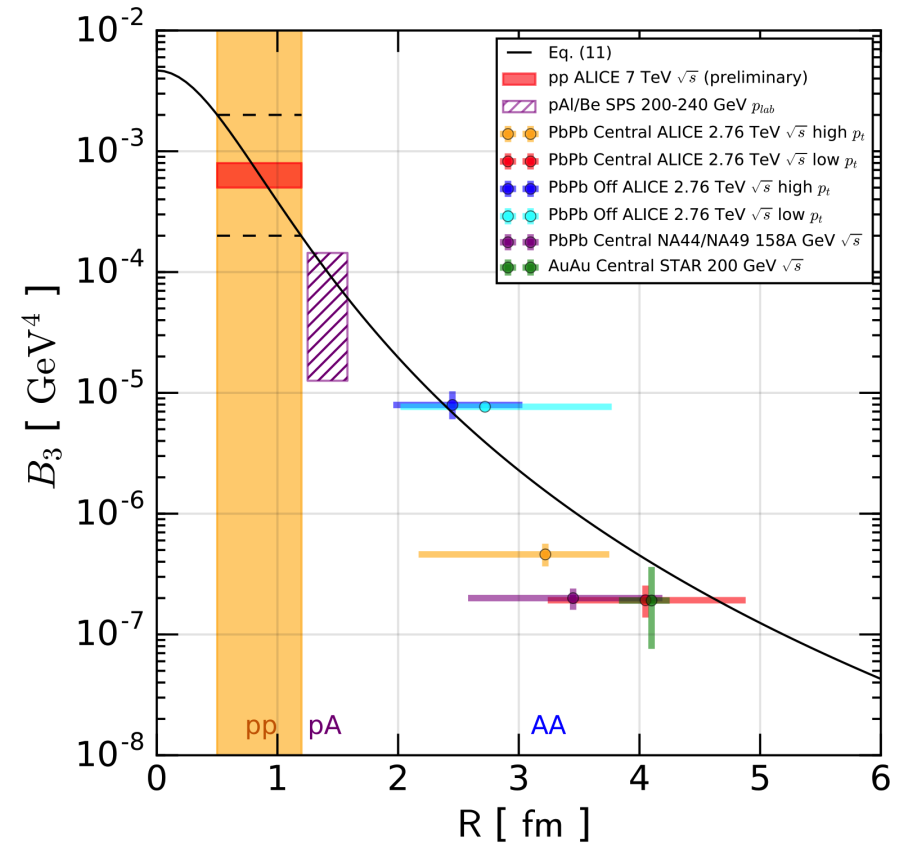
Determination of the coalescence momentum

- No ab initio determination of p_0 which needs to be fitted to data.
To do so, a model is required.

In Blum et al., $B_A \propto V^{1-A}$. The hadronic volume V is probed by the HBT two-pion correlation measurements.

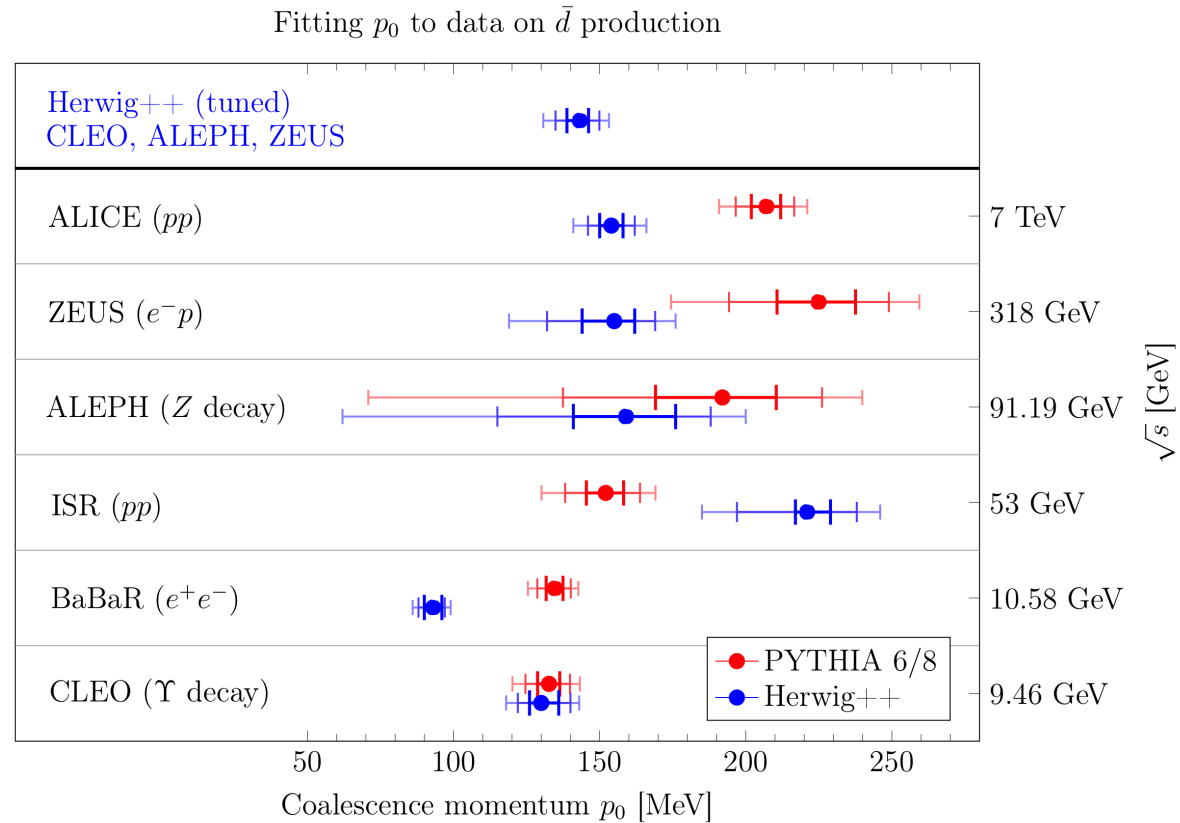
$$\frac{B_2}{\text{GeV}^2} \approx 0.068 \left(\left(\frac{R(p_t)}{1 \text{ fm}} \right)^2 + 2.6 \left(\frac{b_2}{3.2 \text{ fm}} \right)^2 \right)^{-\frac{3}{2}} \quad (12)$$

$$\frac{B_3}{\text{GeV}^4} \approx 0.0024 \left(\left(\frac{R(p_t)}{1 \text{ fm}} \right)^2 + 0.8 \left(\frac{b_3}{1.75 \text{ fm}} \right)^2 \right)^{-3} \quad (11)$$



Determination of the coalescence momentum

- Monte-Carlo event-generators are not devoid of problems.
They are tuned to specific processes $\neq$ antinucleon production.
They yield different p_0 when adjusted to different data sets.
 p_0 depends on $\sqrt{s}$.

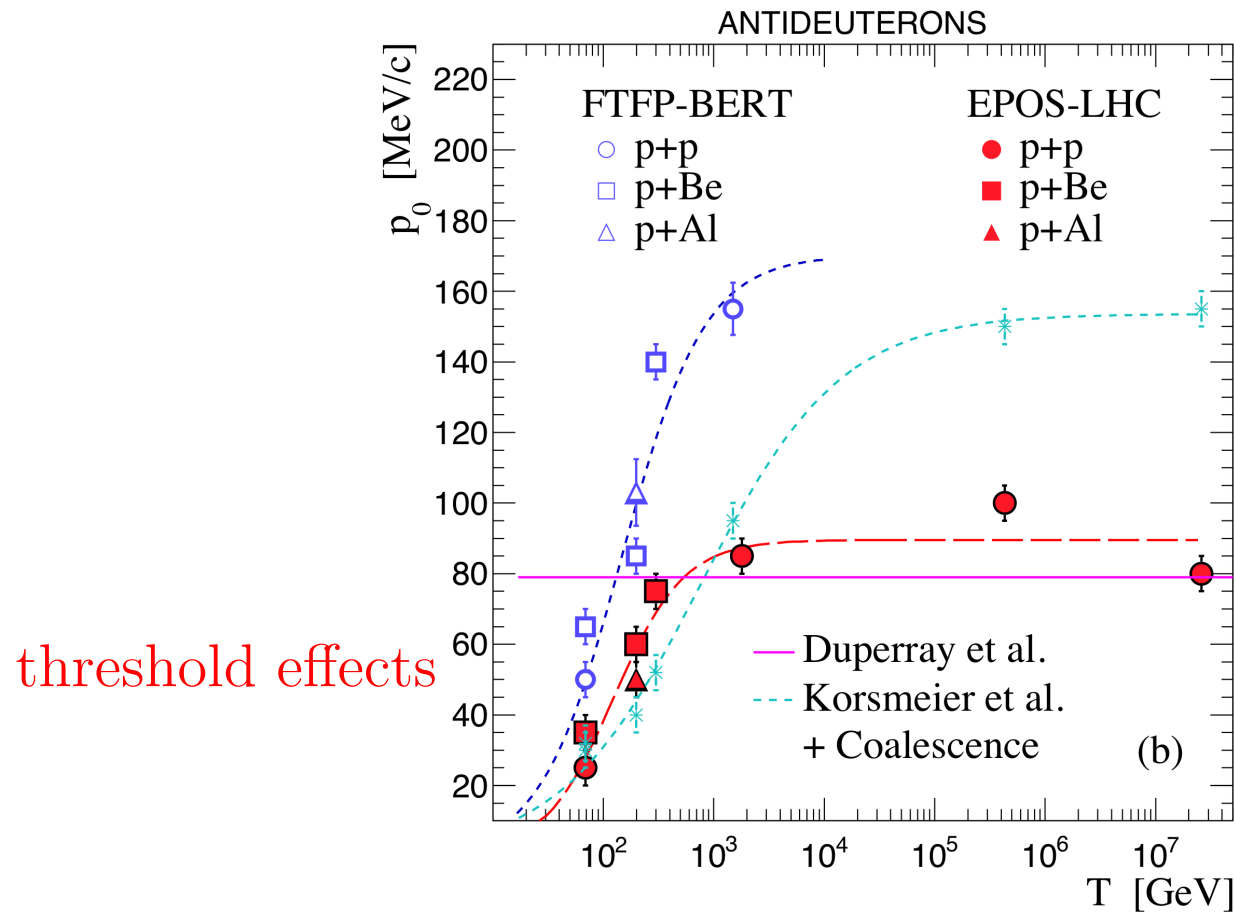


A. Ibarra & S. Wild, JCAP **1302** (2013) 021

L.A. Dal & A.R. Raklev, Phys. Rev. **D89** (2014) 103504

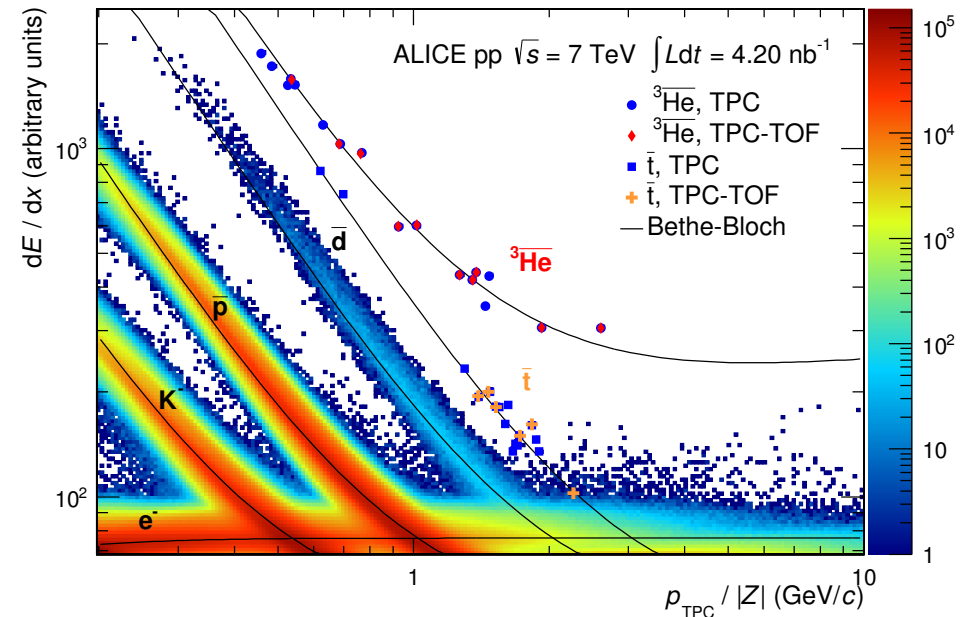
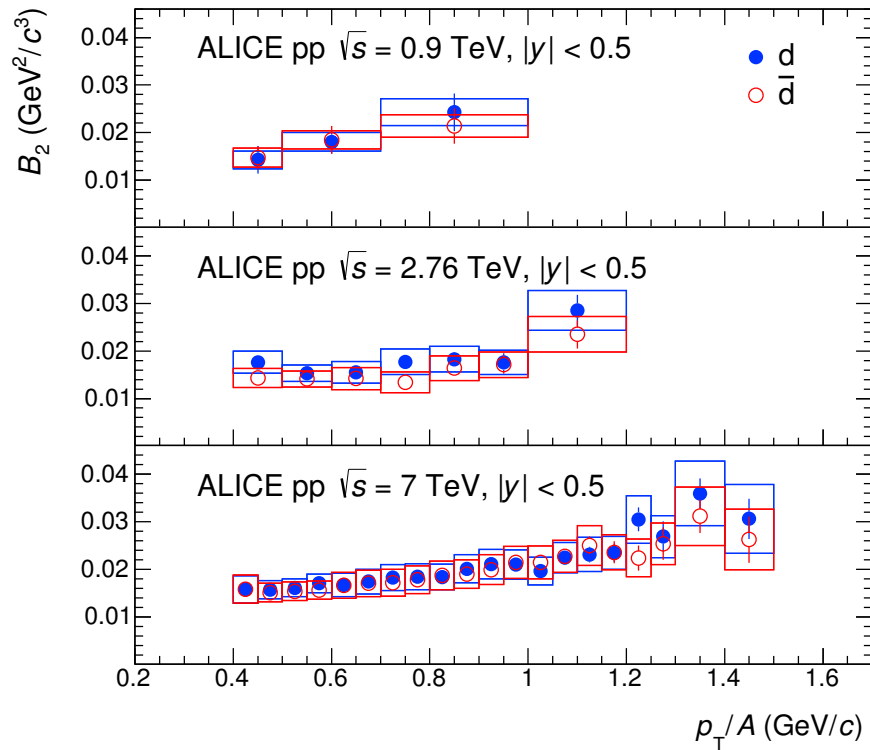
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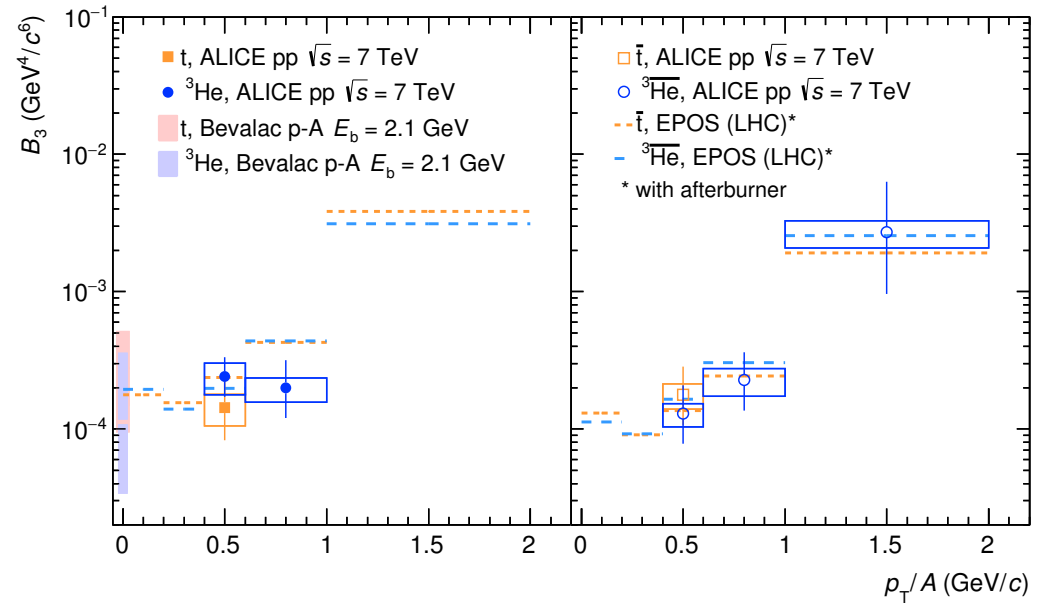
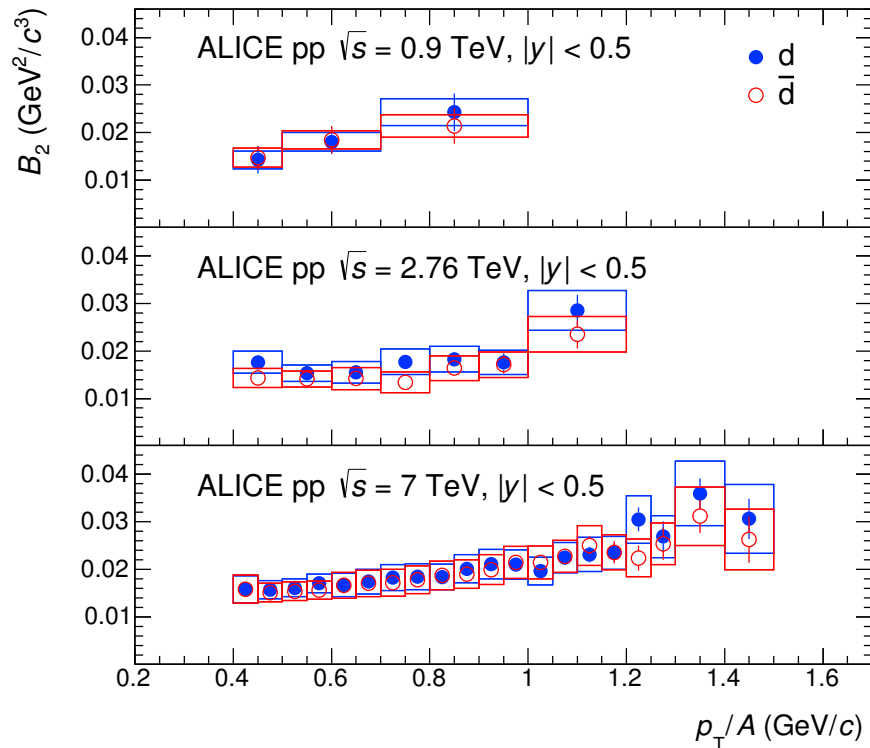
Determination of the coalescence momentum

- ALICE provides an **experimental** determination of B_2 and B_3 .
 $\bar{p}$ production cross-section is **measured**.
 Approximately the same value for p_0 from $\bar{d}$, $\bar{t}$ and ${}^3\overline{\text{He}}$.



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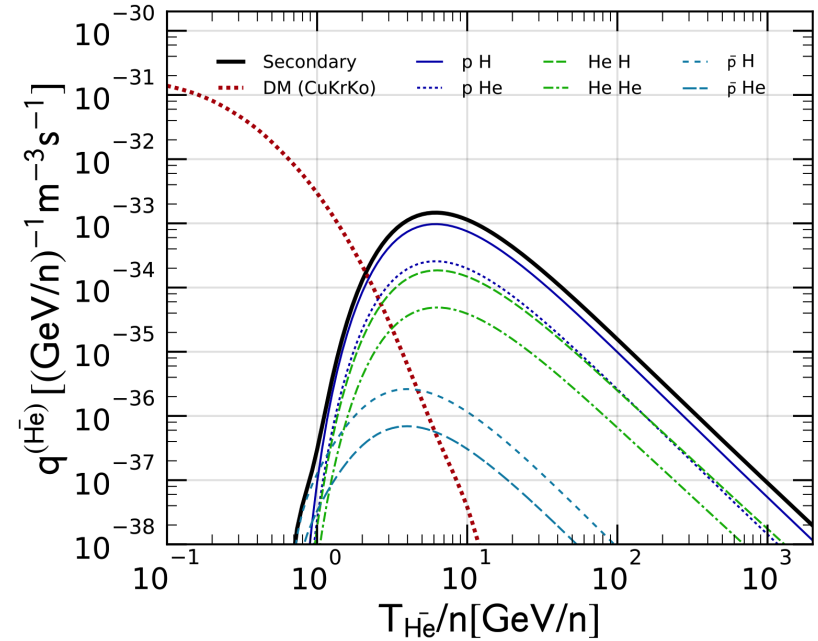
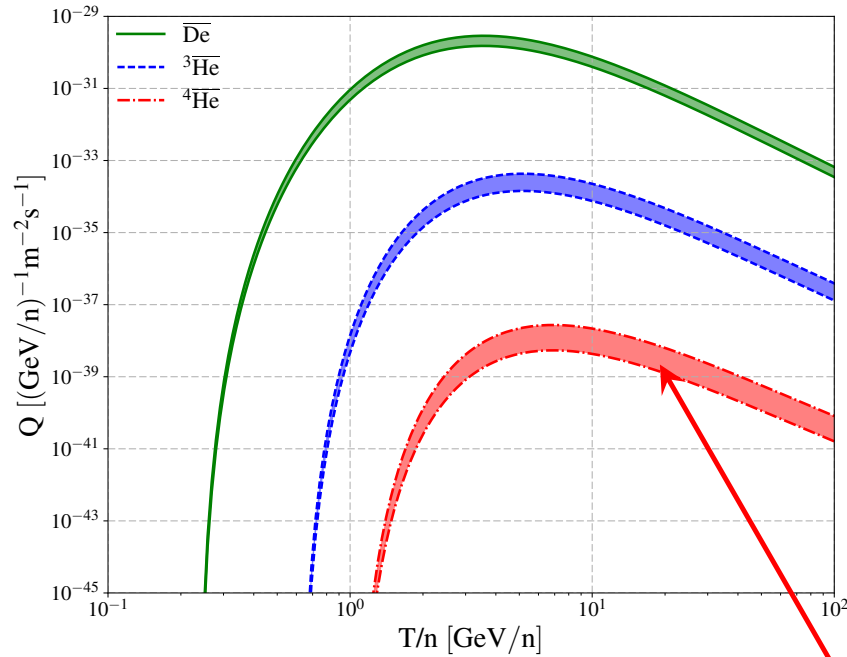


$$208 \text{ MeV} \leq p_{\text{coal}} \leq 262 \text{ MeV}$$

$$218 \text{ MeV} \leq p_{\text{coal}} \leq 262 \text{ MeV}$$

Local source term for anti-nuclei production in cosmic-rays

$$q_{\text{sec}}(\overline{\text{He}} | E_{\overline{\text{He}}}, \mathbf{x}) = \sum_{i \in \text{p}, \alpha} \sum_{j \in \text{H}, \text{He}} 4\pi \int dE_i \Phi_i(E_i, \mathbf{x}) n_j(\mathbf{x}) \frac{d\sigma_{ij \rightarrow \overline{\text{He}}}}{dE_{\overline{\text{He}}}}(E_i, E_{\overline{\text{He}}})$$



V. Poulin et al., Phys. Rev. **D99** (2019) 023016

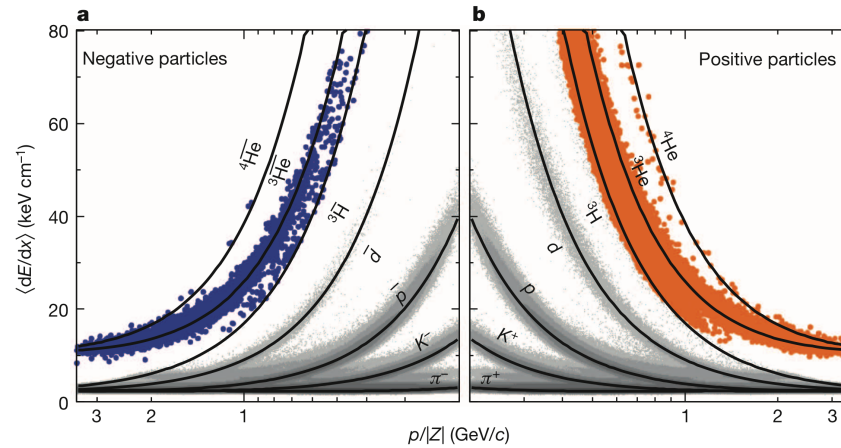
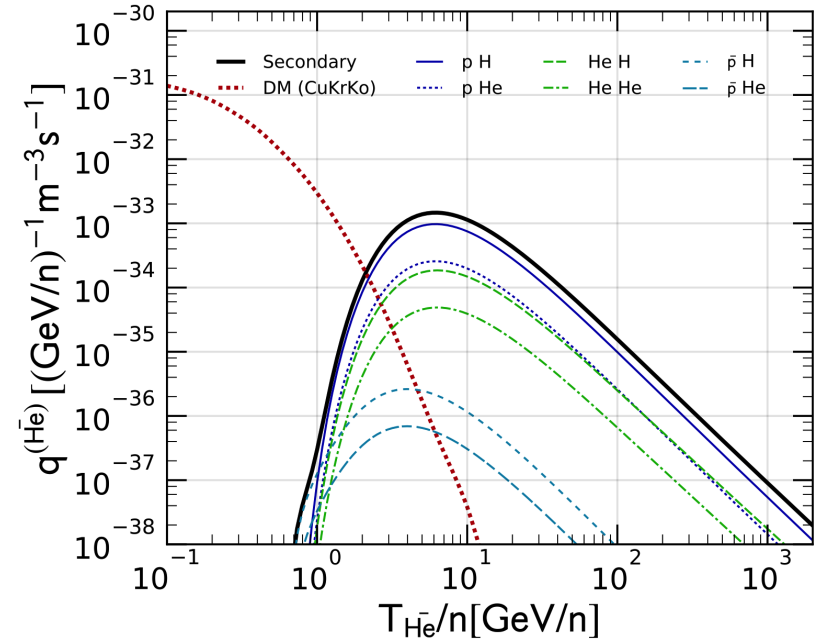
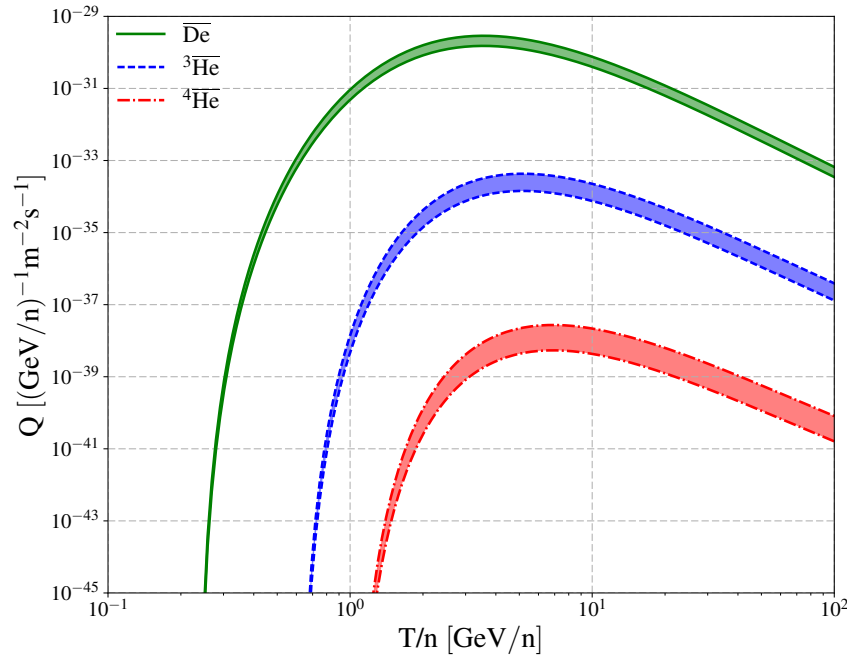
M. Korsmeier et al., Phys. Rev. **D97** (2018) 103011

$$7.7 \times 10^{-7} \leq \frac{B_4}{\text{GeV}^6} \leq 3.9 \times 10^{-6}$$

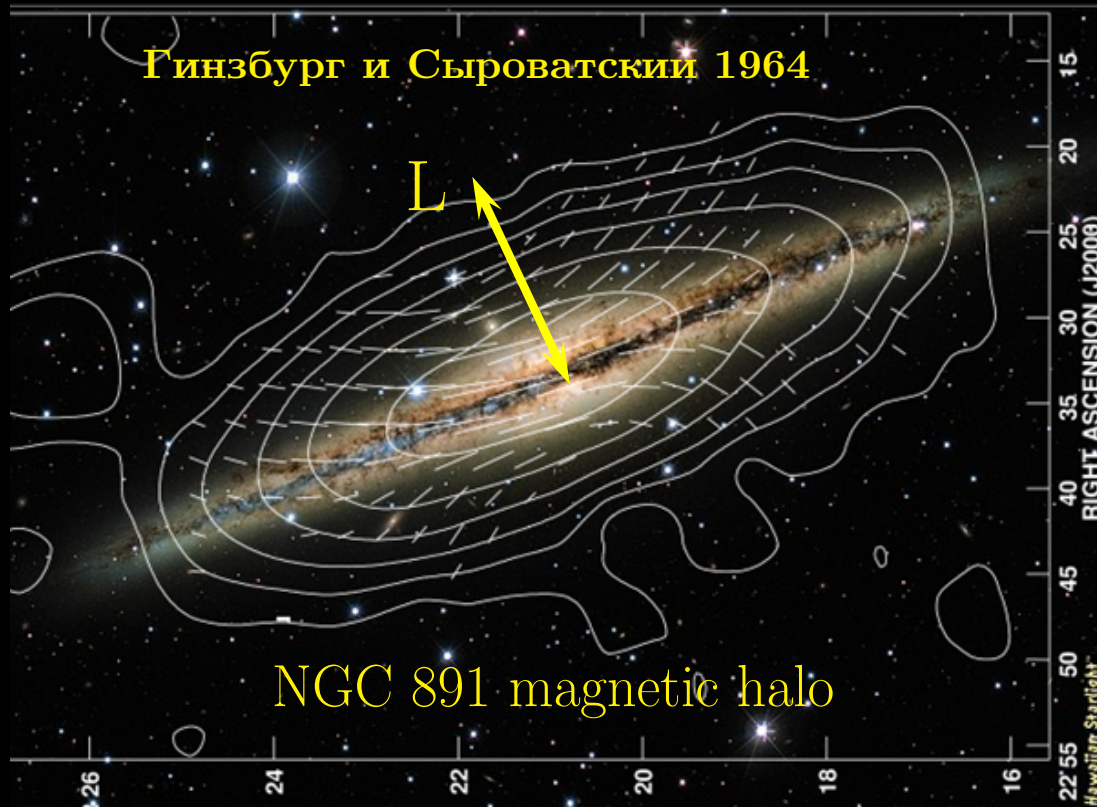
$\bar{p}$ production modeled as in
M. di Mauro et al., Phys. Rev. **D90** (2014) 085017

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3) Cosmic ray anti-helium Galactic propagation



$$\psi = \frac{dn}{dE} = \frac{d^4N}{d^3\mathbf{x}dE}$$

$$\Phi = \frac{1}{4\pi} v \psi$$

$$(\text{GeV}/\text{nuc})^{-1} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

ISM spallation

$$\underbrace{\dot{\psi} + \nabla \cdot \{-K \nabla \psi + \psi \mathbf{V}_C\}}_{\text{convection}} + \underbrace{\frac{\partial}{\partial E} \left\{ b \psi - D_{EE} \frac{\partial \psi}{\partial E} \right\}}_{\text{E losses}} = q - (\sigma v n_H) \psi$$

$q = q_{\text{acc}}, q_{\text{sec}}, q_{\text{DM}}$

x diffusion

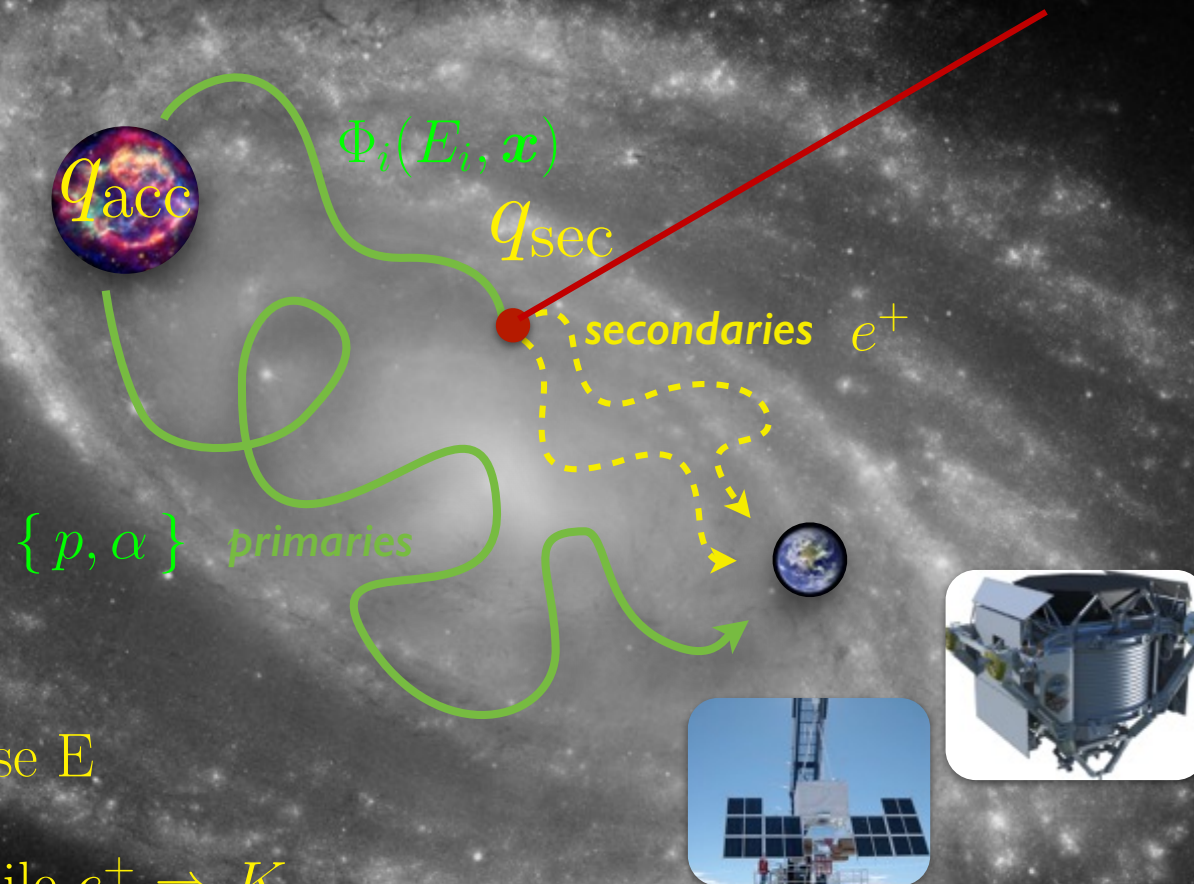
$$K = K_0 \beta \mathcal{R}^\delta \left\{ 1 + \left(\frac{\mathcal{R}}{\mathcal{R}_b} \right)^{\Delta\delta/s} \right\}^{-s}$$

E diffusion

$$D_{EE} = \frac{2}{9} \frac{V_A^2 \beta^4 E^2}{K}$$

Positron constraints on cosmic ray propagation

$$q_{\text{sec}}(e^+ | E_e, \mathbf{x}) = \sum_{i \in p, \alpha} \sum_{j \in H, He} 4\pi \int dE_i \Phi_i(E_i, \mathbf{x}) n_j(\mathbf{x}) \frac{d\sigma_{ij \rightarrow e^+}}{dE_e}(E_i, E_e)$$

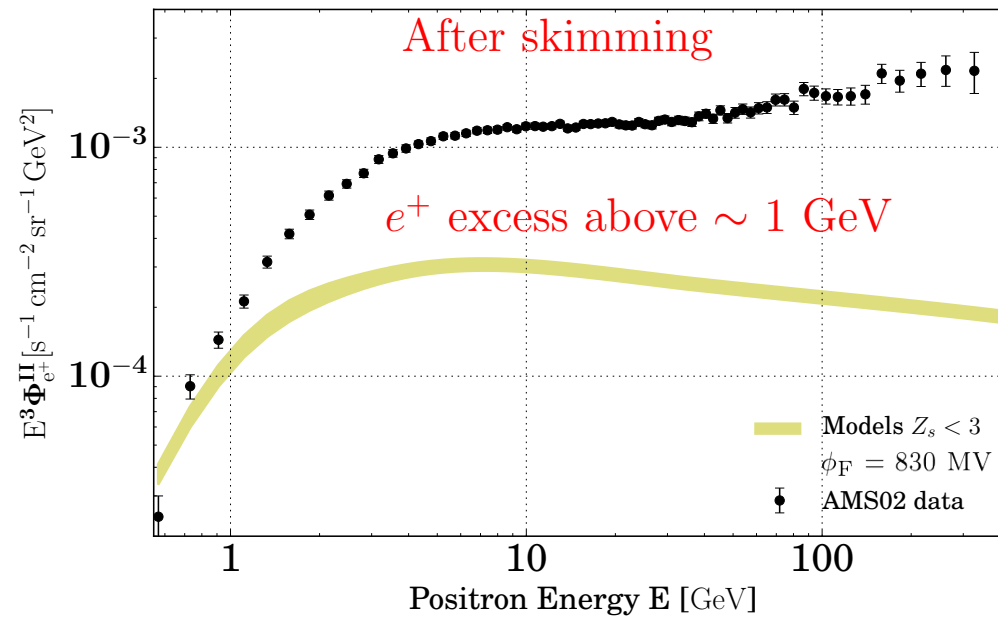
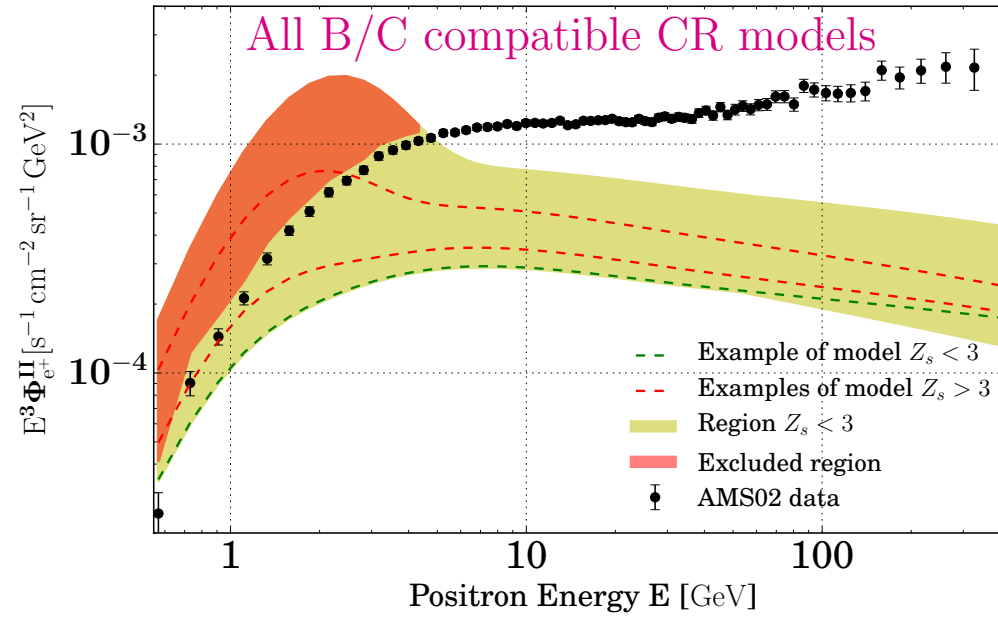


e^+ diffuse and lose E

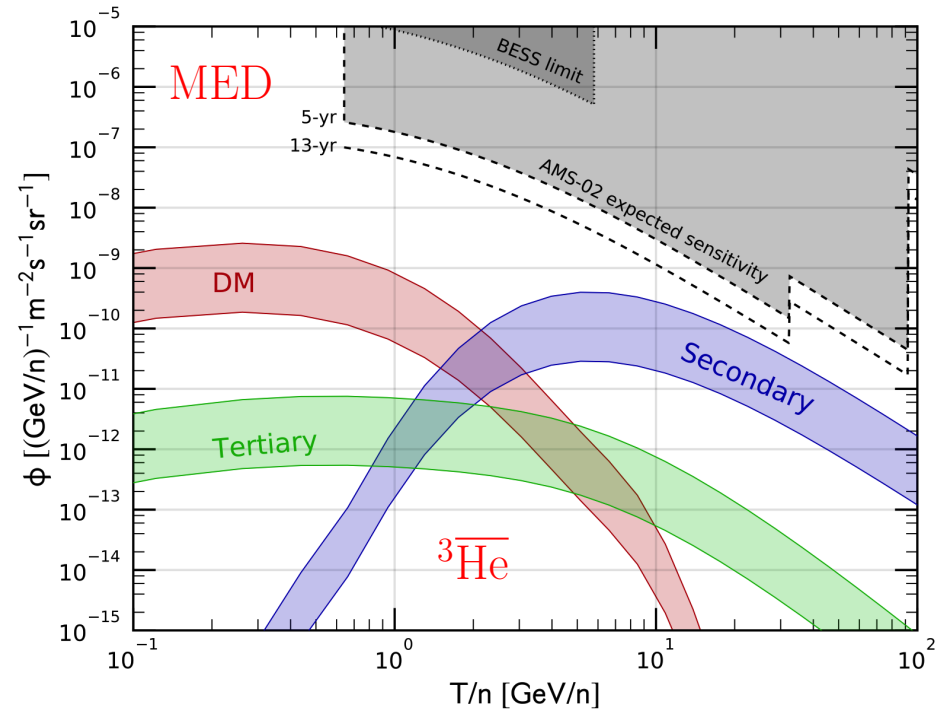
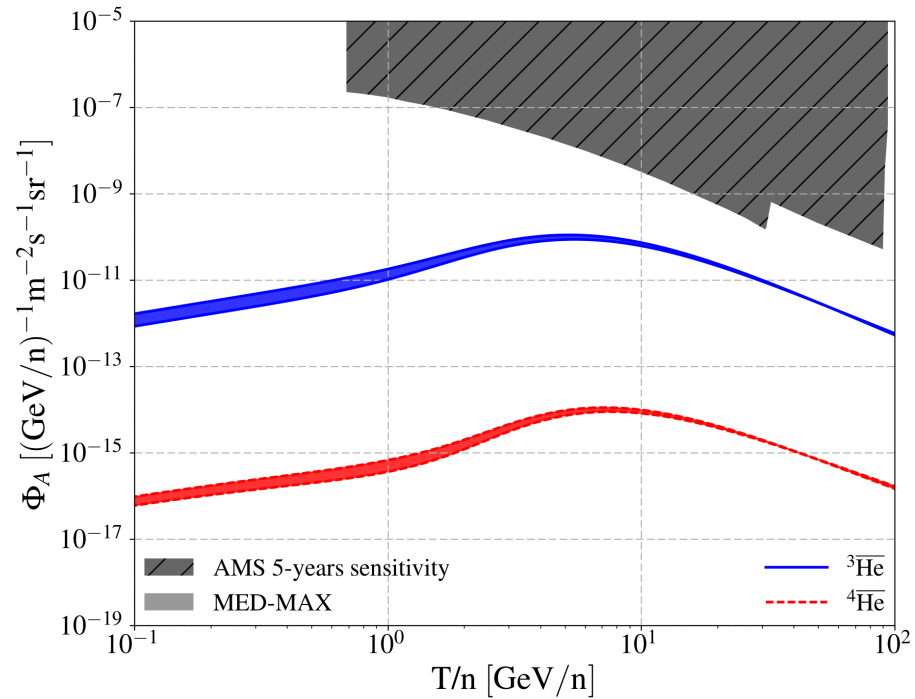
B/C $\Rightarrow$ K/L while $e^+ \Rightarrow K$

Degeneracy between K and L can be lifted

$\Phi_{e^+}^{\text{sec}}$ forces K to be large $\Rightarrow L \geq L_{\text{min}} \sim 8.5$ kpc



Secondary anti-helium fluxes



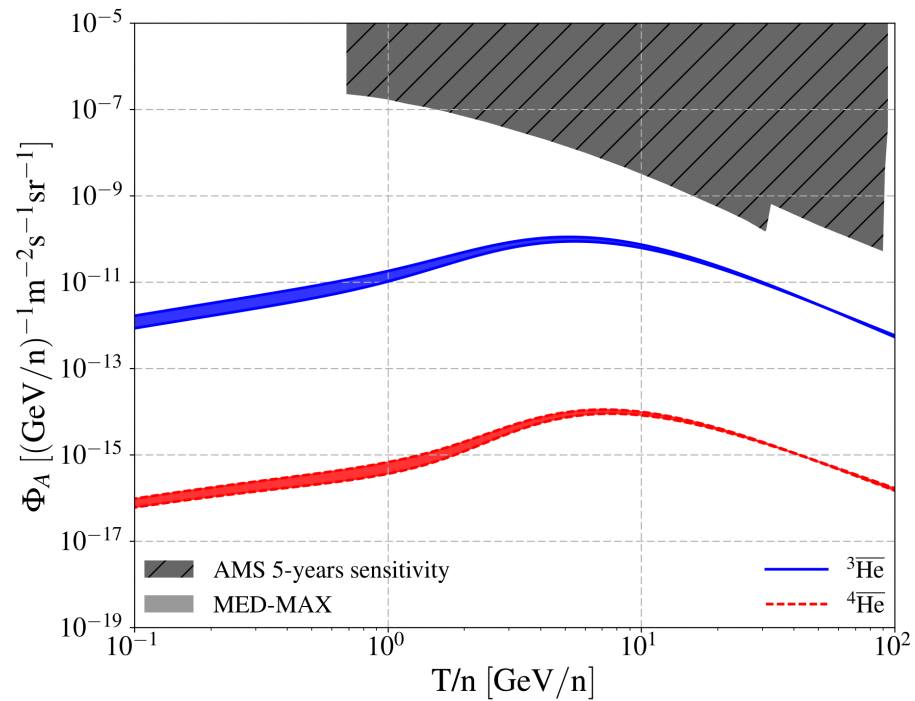
V. Poulin et al., Phys. Rev. **D99** (2019) 023016

M. Korsmeier et al., Phys. Rev. **D97** (2018) 103011

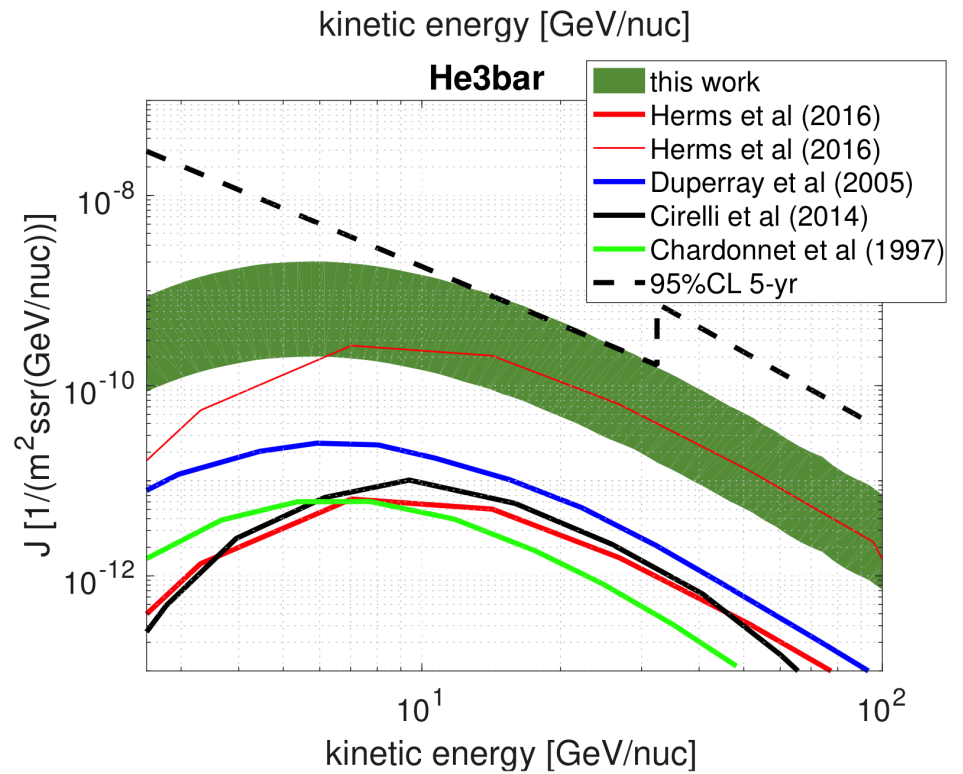
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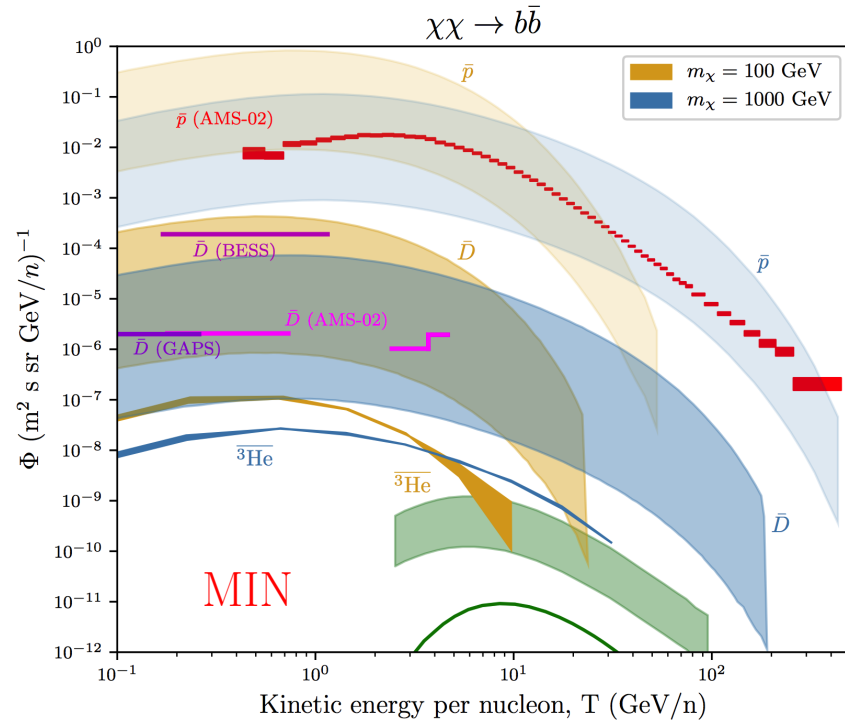
K. Blum et al., Phys. Rev. **D96** (2017) 103021

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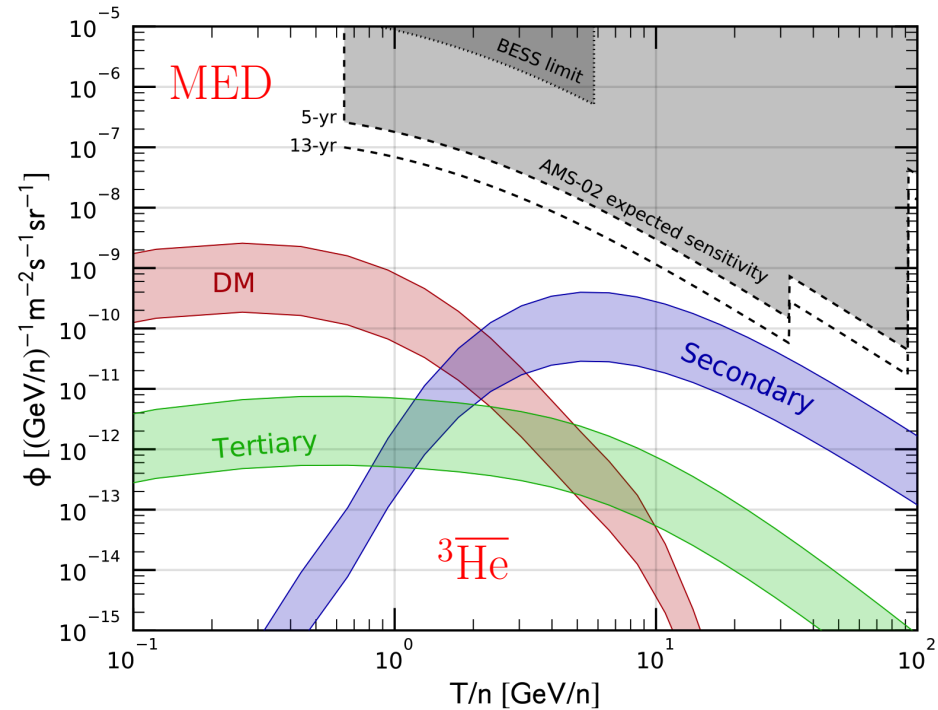
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4) A word on dark matter

Could anti-helium (${}^3\overline{\text{He}}$) events be produced by DM?



A. Coogan & S. Profumo,
Phys. Rev. **D96** (2017) 083020

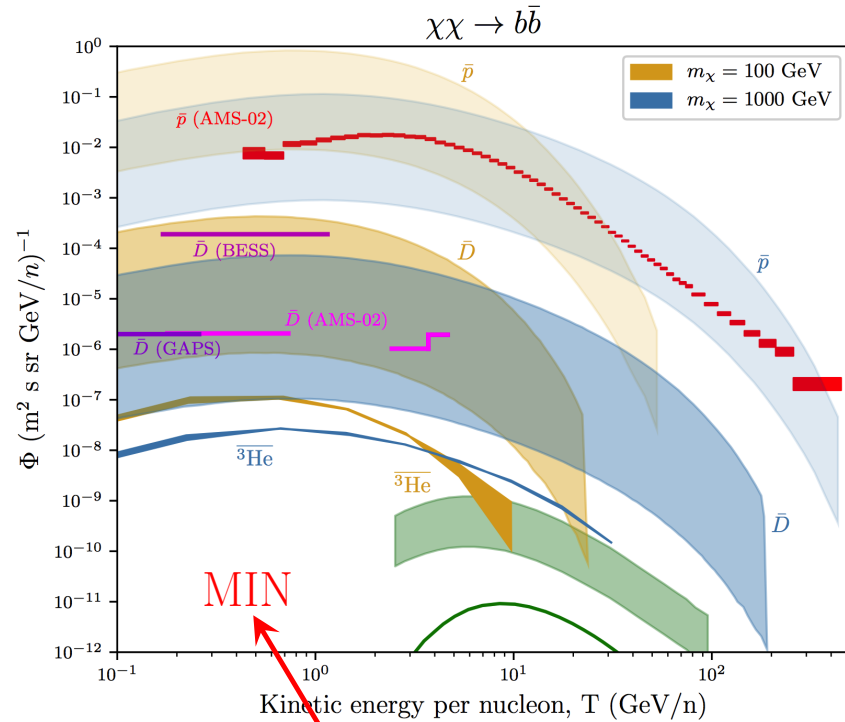


M. Korsmeier et al., Phys. Rev. **D97** (2018) 103011

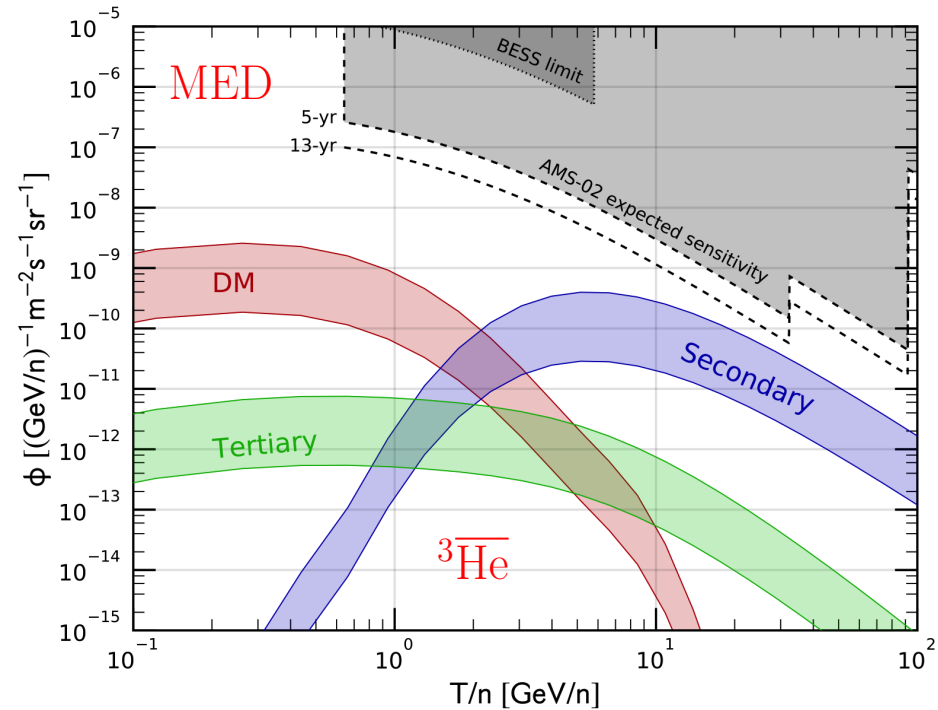
- If AMS-02 $\overline{\text{He}}$ events are from DM, beware of $\bar{p}$ flux.
- To evade the $\bar{p}$ constraint, p_{coal} exceedingly large.

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MIN is excluded by positrons

Concluding remarks

- Anti-helium-3 and anti-helium-4 candidates have been identified by AMS-02. Massive background simulations are carried out to evaluate significance. More data are needed.
- ${}^3\overline{\text{He}}$ events
AMS-02 should **not** see secondary CR ${}^3\overline{\text{He}}$.
If $\overline{\text{He}}$ events are produced by DM, a large $\bar{p}$ excess is expected.
Apart from a possible anomaly, no such excess is seen.
- ${}^4\overline{\text{He}}$ events
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If confirmed, a single ${}^4\overline{\text{He}}$ would be a major discovery

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Thanks for your attention